

Observation of spinon Anderson localization in a spin-1/2 antiferromagnetic Heisenberg chain

Shiyan Li

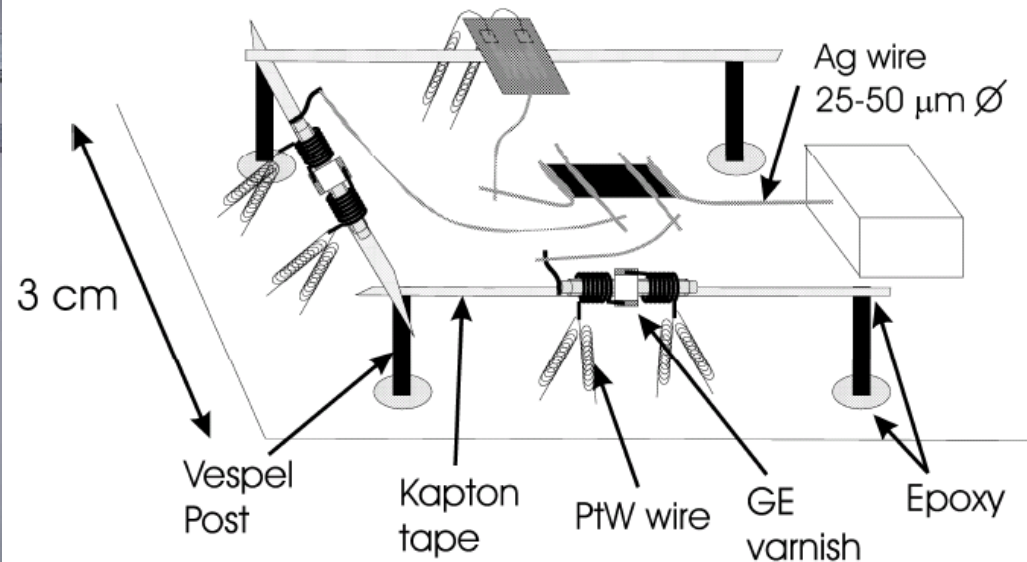
Fudan University

April 2013, IAS, Tsinghua

Outline:

- 1. Ultra-low-temperature heat transport measurement**
- 2. Some examples of heat transport by magnetic excitations**
- 3. Anderson localization of spinons in a spin-1/2 antiferromagnetic Heisenberg chain**
- 4. Unveiling the quantum critical point of an Ising chain in a transverse field**
- 5. Summary**

1. Ultra-low-temperature heat transport measurement



^3He - ^4He dilution fridge
 $T \rightarrow 7 \text{ mK}$; $H \rightarrow 17 \text{ T}$

$$\kappa = \alpha \frac{\dot{Q}}{\Delta T}$$

Heat transport: A tool to probe low-lying quasiparticles

$$\kappa = \kappa_{\text{electrons}} + \kappa_{\text{phonons}} + \kappa_{\text{magnons}} + \kappa_{\text{spinons}} \dots$$

$$\kappa = 1/3 C v l$$

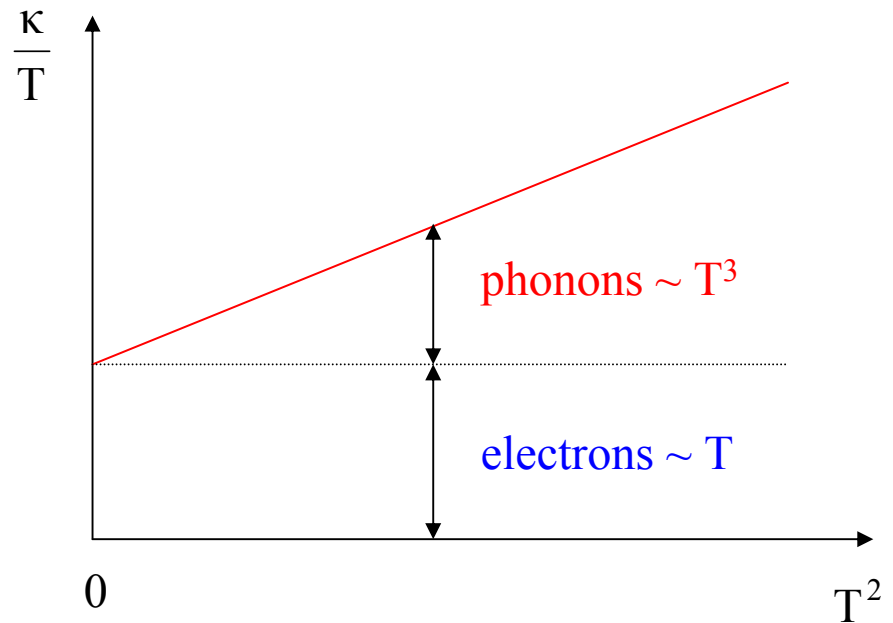
FERMIONS (Electrons)

$$\kappa \propto C_e \propto T$$

BOSONS (Phonons)

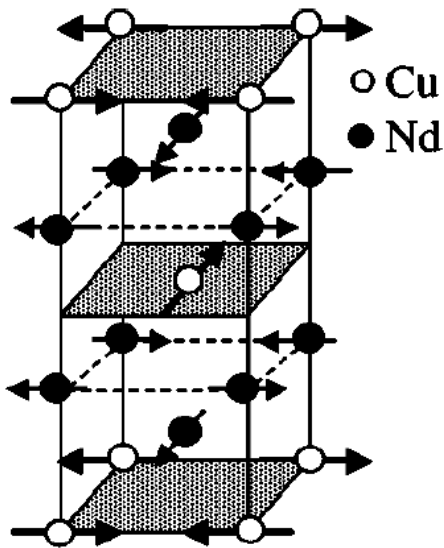
$$\kappa \propto C_{\text{ph}} \propto T^3$$

$$\kappa/T = A + BT^2$$



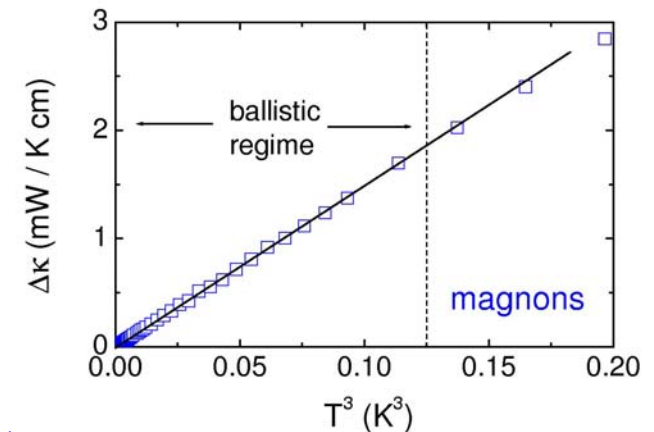
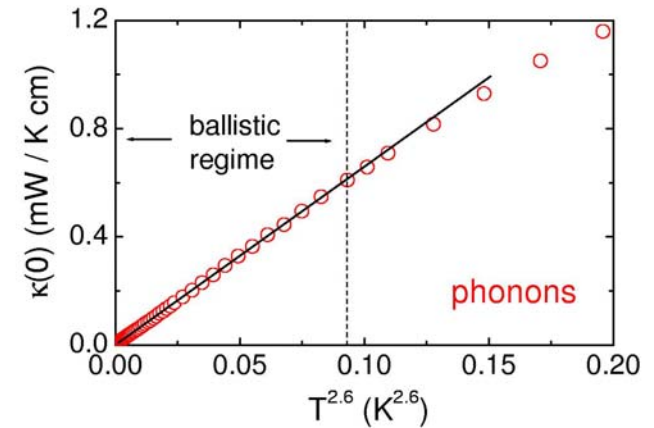
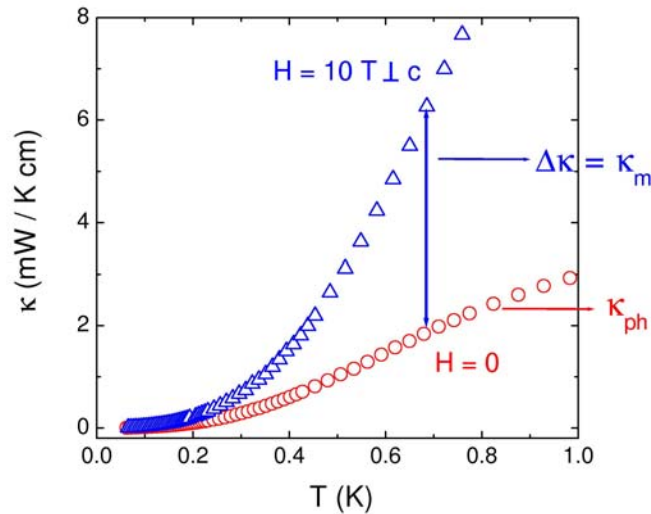
2. Some examples of heat transport by magnetic excitations

Example 1: AF magnon heat transport in Nd_2CuO_4



Spin-flop transition
in $H \perp c$

Switch on acoustic magnons



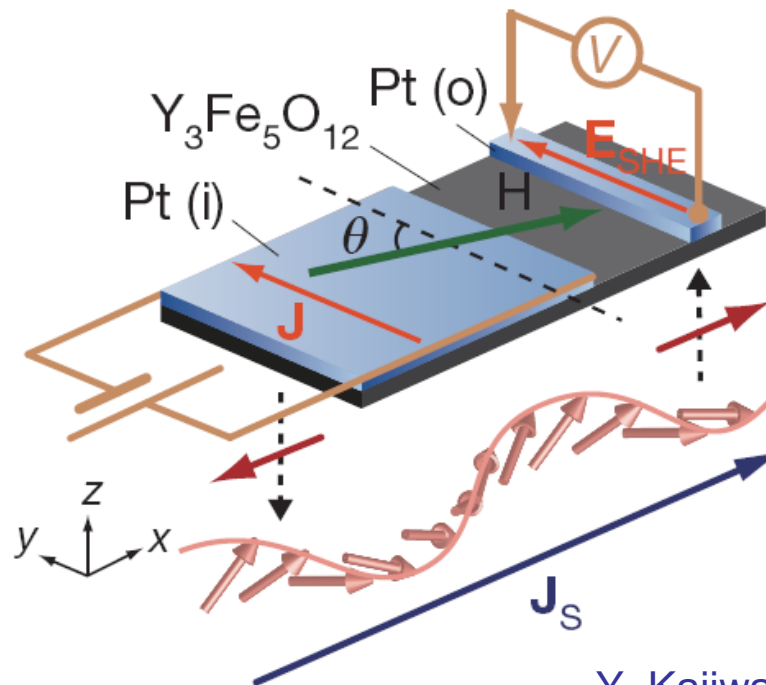
First observation of $\kappa \sim T^3$ AF magnon heat transport



Example 2: FM magnon heat transport in YIG

Transmission of electrical signals by spin-wave interconversion in a magnetic insulator

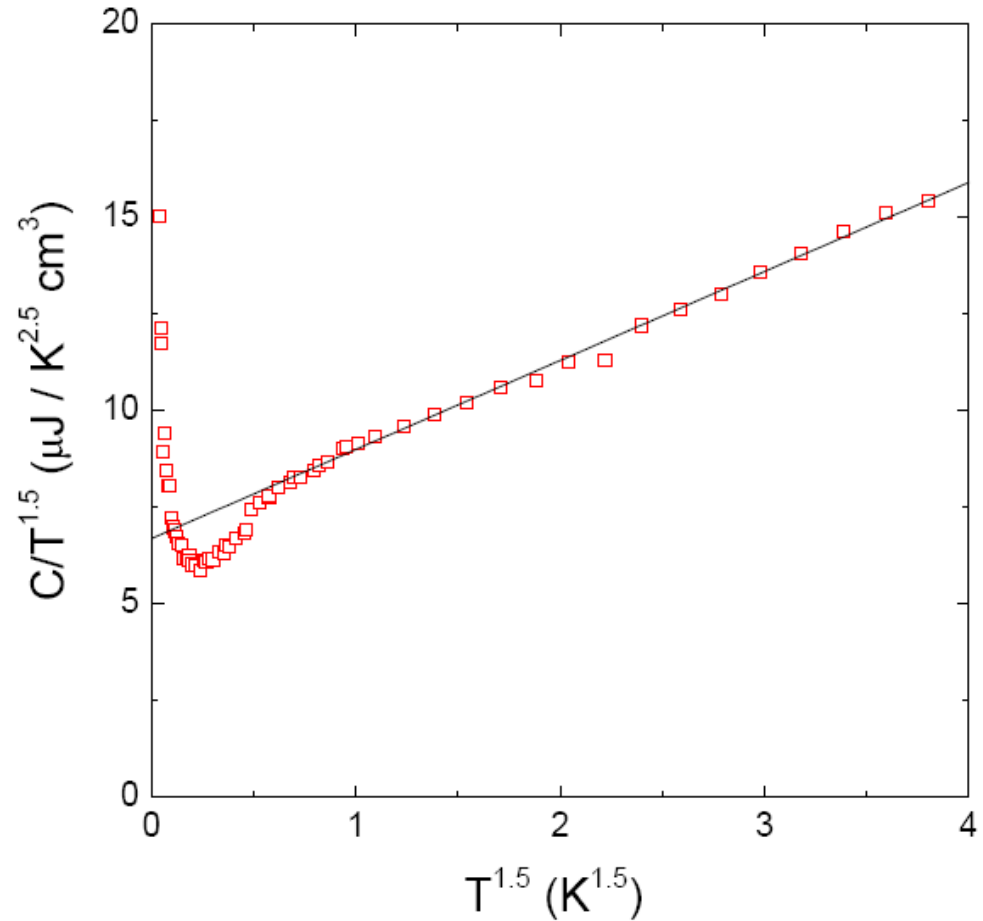
Y. Kajiwara^{1,2}, K. Harii¹, S. Takahashi^{1,3}, J. Ohe^{1,3}, K. Uchida¹, M. Mizuguchi¹, H. Umezawa⁵, H. Kawai⁵, K. Ando^{1,2}, K. Takanashi¹, S. Maekawa^{1,3} & E. Saitoh^{1,2,4}



$\text{Y}_3\text{Fe}_5\text{O}_{12}$ (YIG)
typical ferrimagnet

Specific heat:

FM magnon in YIG single crystal



At not very low T : $E = Dk^2$

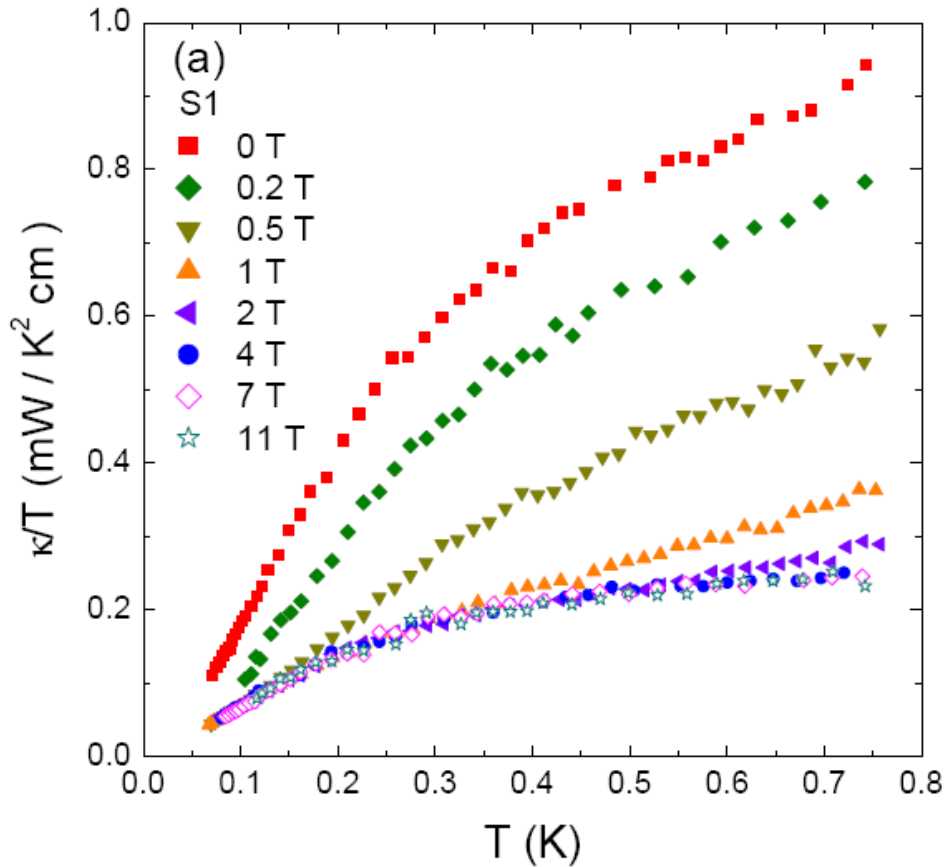
$$C_m(T) = \frac{15\zeta(5/2)k_B^{2.5}T^{1.5}}{32\pi^{1.5}D^{1.5}}$$

$$\kappa_m(T) = \frac{\zeta(3)k_B^3LT^2}{\pi^2\hbar D}$$

0.77 K < T < 2.5 K: $C = 6.7T^{1.5} + 2.3T^3$
T < 0.77 K: dipole-dipole correction

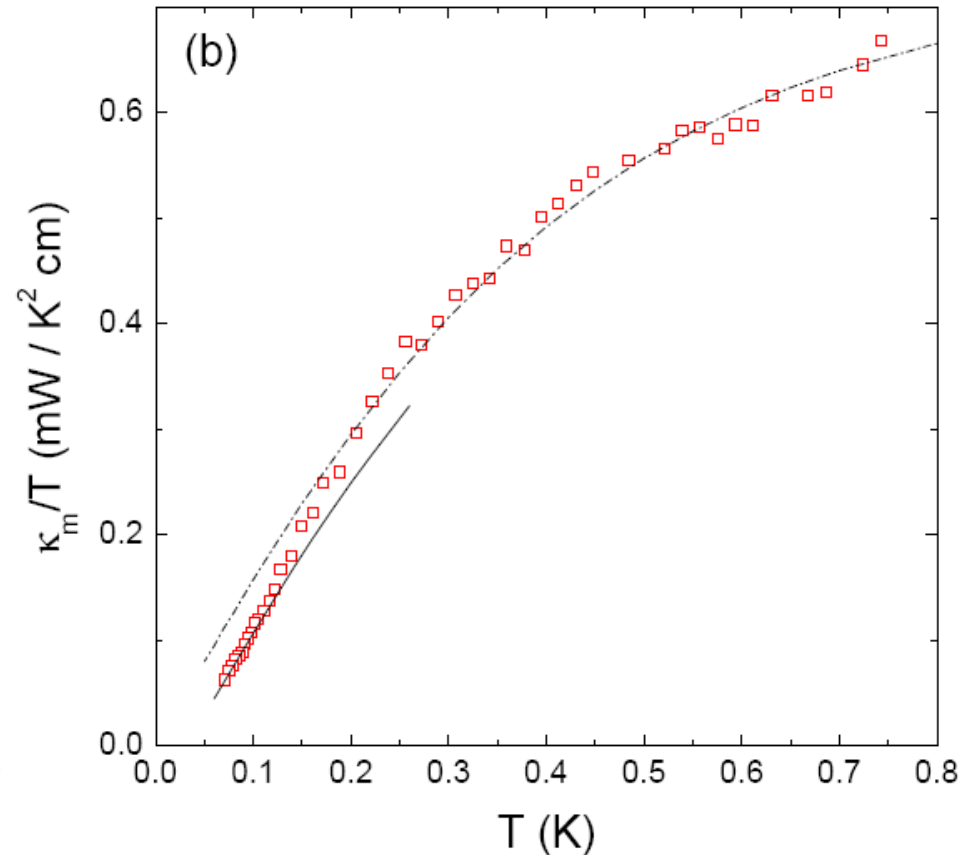
B. Y. Pan, S. Y. Li *et al.*, arXiv:1302.6739

Thermal conductivity: FM magnon in YIG single crystal



Magnon gap in field: $\Delta = g\mu_B H$

$$\kappa_m = \kappa(0T) - \kappa(4T)$$



If no corrections: $\kappa_m \sim T^2$

Our result suggests the corrections of defects and dipole-dipole interaction are needed.

Example 3: Spinon heat transport in spin liquids



EXOTIC MATTER

Leon Balents, Nature **464**, 199 (2010)

Spin liquids in frustrated magnets

Leon Balents¹

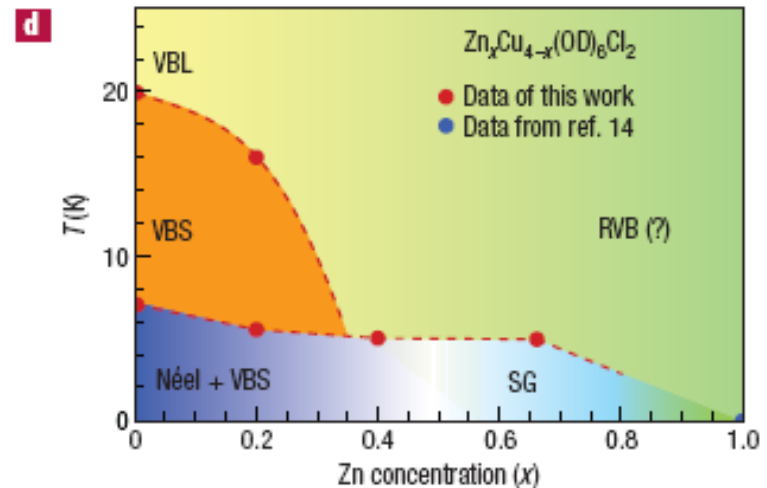
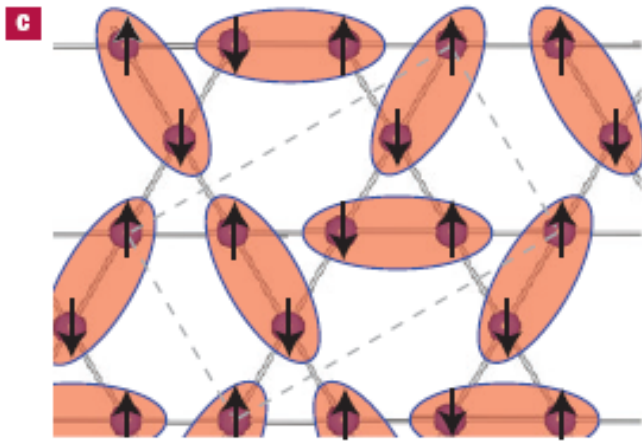
New magnetic ground state!

Table 1 | Some experimental materials studied in the search for QSLs

Material	Lattice	S	Θ_{CW} (K)	R^*	Status or explanation
κ -(BEDT-TTF) ₂ Cu ₂ (CN) ₃	Triangular†	½	-375‡	1.8	Possible QSL
EtMe ₃ Sb[Pd(dmit) ₂] ₂	Triangular†	½	-(375-325)‡	?	Possible QSL
Cu ₃ V ₂ O ₇ (OH) ₂ •2H ₂ O (volborthite)	Kagomé†	½	-115	6	Magnetic
ZnCu ₃ (OH) ₆ Cl ₂ (herbertsmithite)	Kagomé	½	-241	?	Possible QSL
BaCu ₃ V ₂ O ₈ (OH) ₂ (vesignieite)	Kagomé†	½	-77	4	Possible QSL
Na ₄ Ir ₃ O ₈	Hyperkagomé	½	-650	70	Possible QSL
Cs ₂ CuCl ₄	Triangular†	½	-4	0	Dimensional reduction
FeSc ₂ S ₄	Diamond	2	-45	230	Quantum criticality

BEDT-TTF, bis(ethylenedithio)-tetrathiafulvalene; dmit, 1,3-dithiole-2-thione-4,5-dithiolate; Et, ethyl; Me, methyl. * R is the Wilson ratio, which is defined in equation (1) in the main text. For EtMe₃Sb[Pd(dmit)₂]₂ and ZnCu₃(OH)₆Cl₂, experimental data for the intrinsic low-temperature specific heat are not available, hence R is not determined. †Some degree of spatial anisotropy is present, implying that $J' \neq J$ in Fig. 1a. ‡ Θ_{CW} a theoretical Curie-Weiss temperature (Θ_{CW}) calculated from the high-temperature expansion for an $S = \frac{1}{2}$ triangular lattice; $\Theta_{\text{CW}} = 3J/2k_B$, using the J fitted to experiment.

Quantum-spin-liquid states in the two-dimensional kagome antiferromagnets $\text{Zn}_x\text{Cu}_{4-x}(\text{OD})_6\text{Cl}_2$

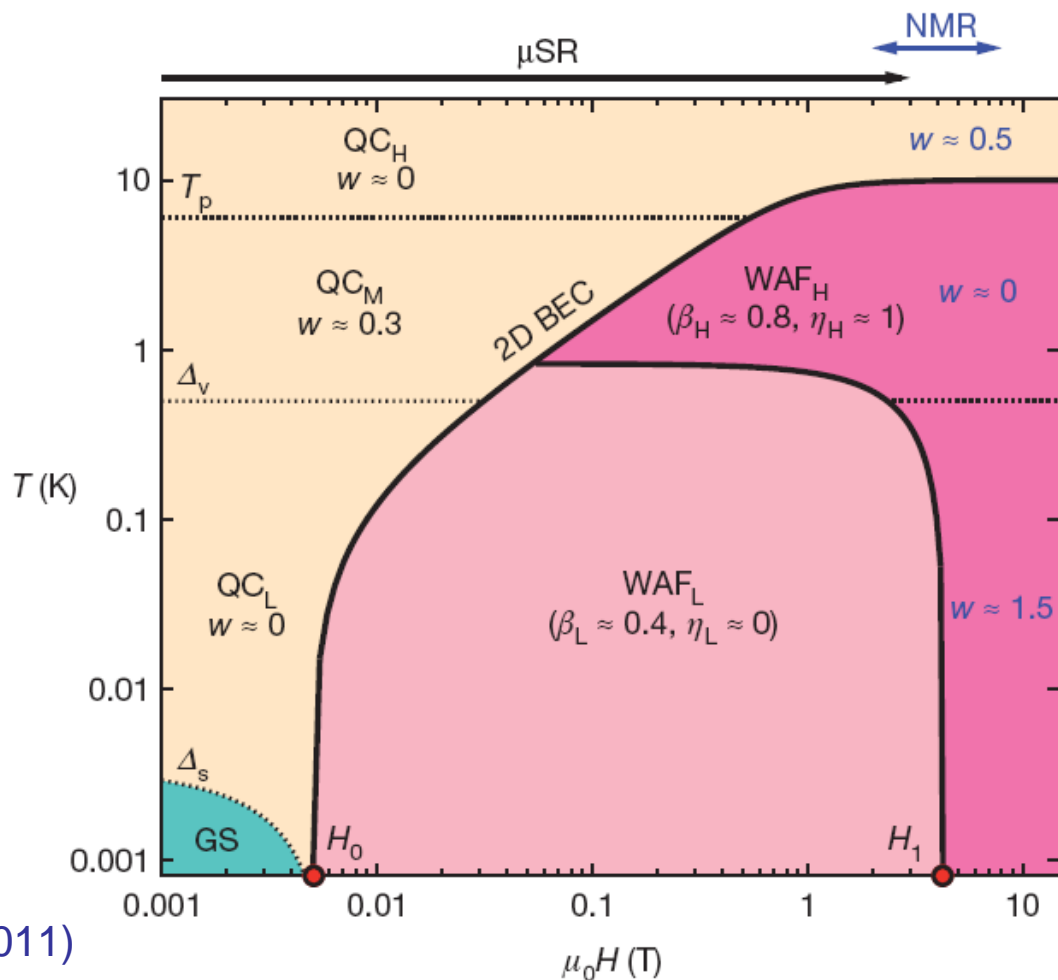


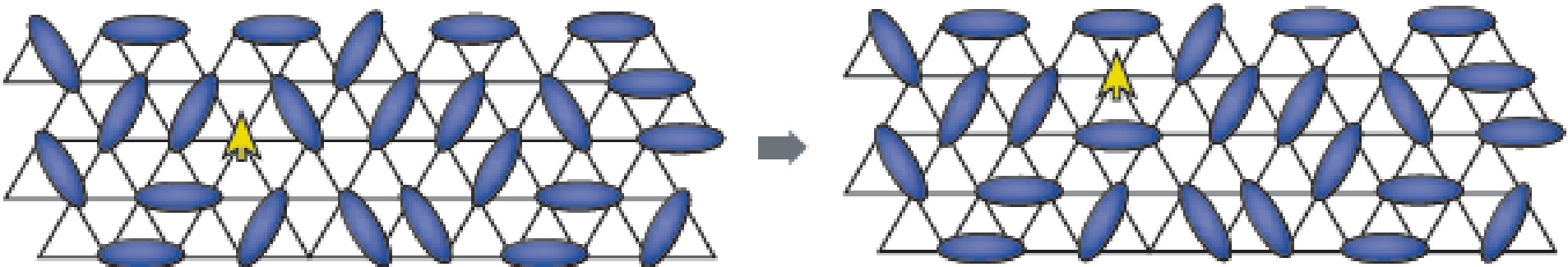
Neutron scattering

Magnetic and non-magnetic phases of a quantum spin liquid

μ SR, NMR

κ -(BEDT-TTF)₂Cu₂(CN)₃





Spinon excitation in a 2D QSL detected by heat transport

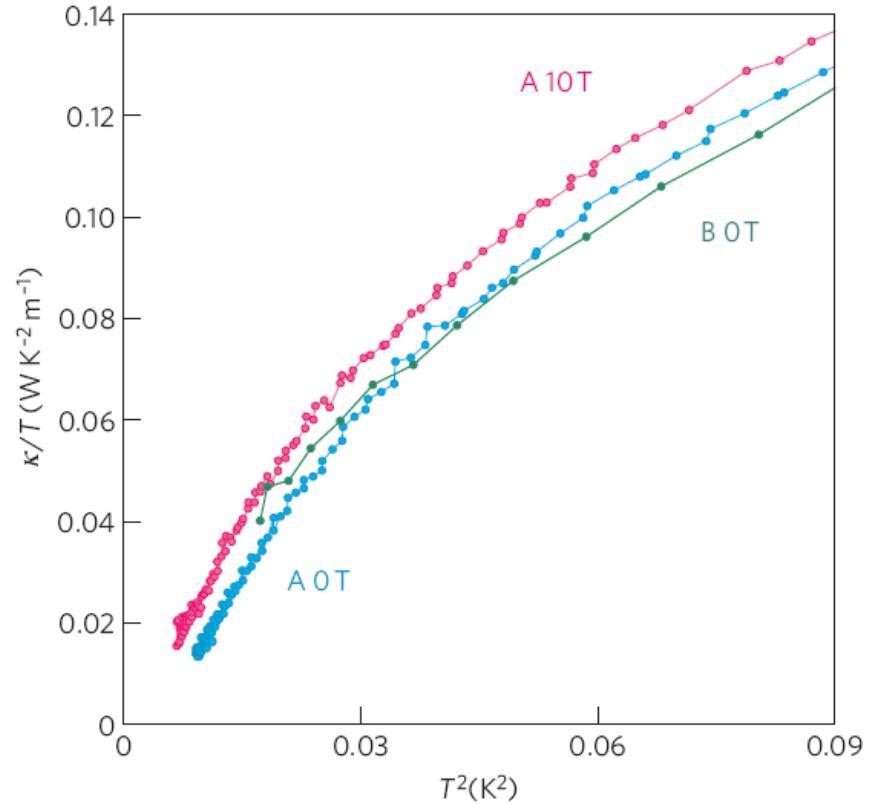
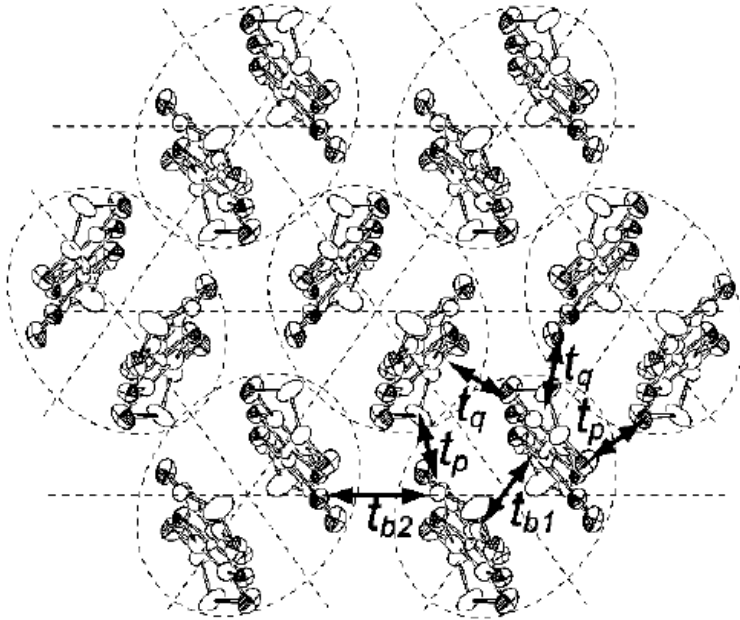
Amperean Pairing Instability in the U(1) Spin Liquid State with Fermi Surface and Application to κ -(BEDT-TTF)₂Cu₂(CN)₃ ARPES

Sung-Sik Lee, Patrick Lee and T. Senthil, PRL **98**, 067006 (2007)

Prediction: $\kappa \sim T$, like electrons in a metal

Heat transport:

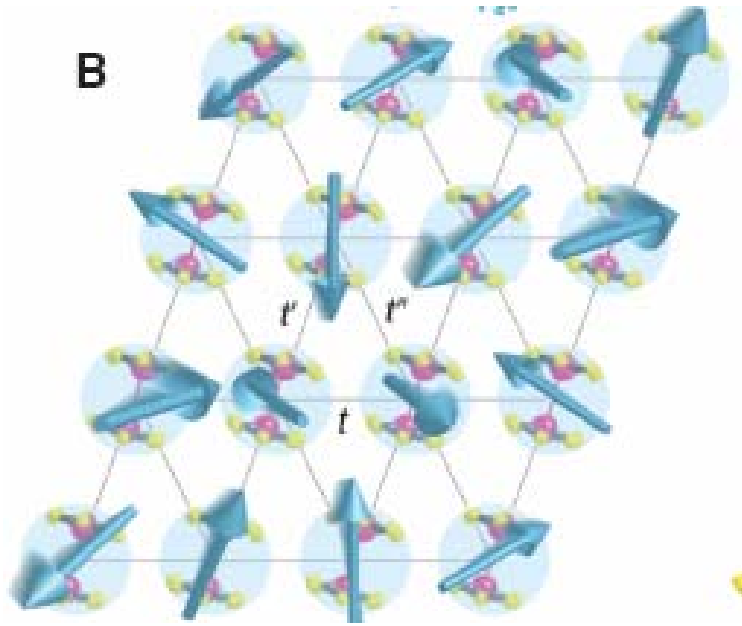
A tool to probe spinons



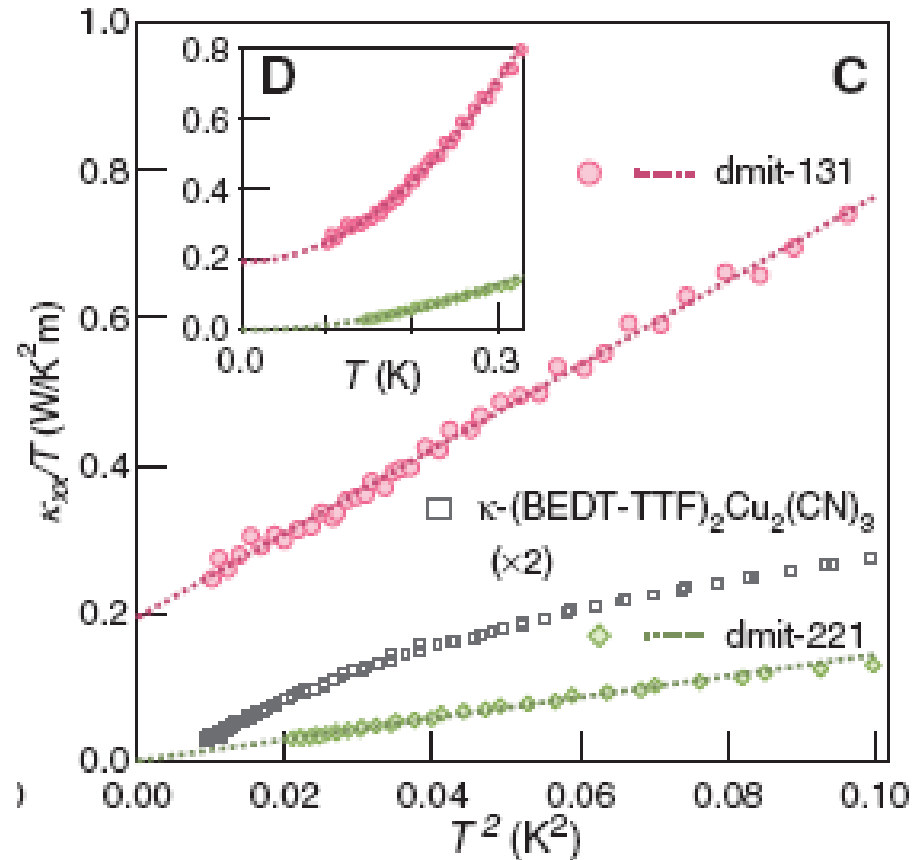
No κ_0/T : are spinons gapped?

Heat transport:

A tool to probe spinons

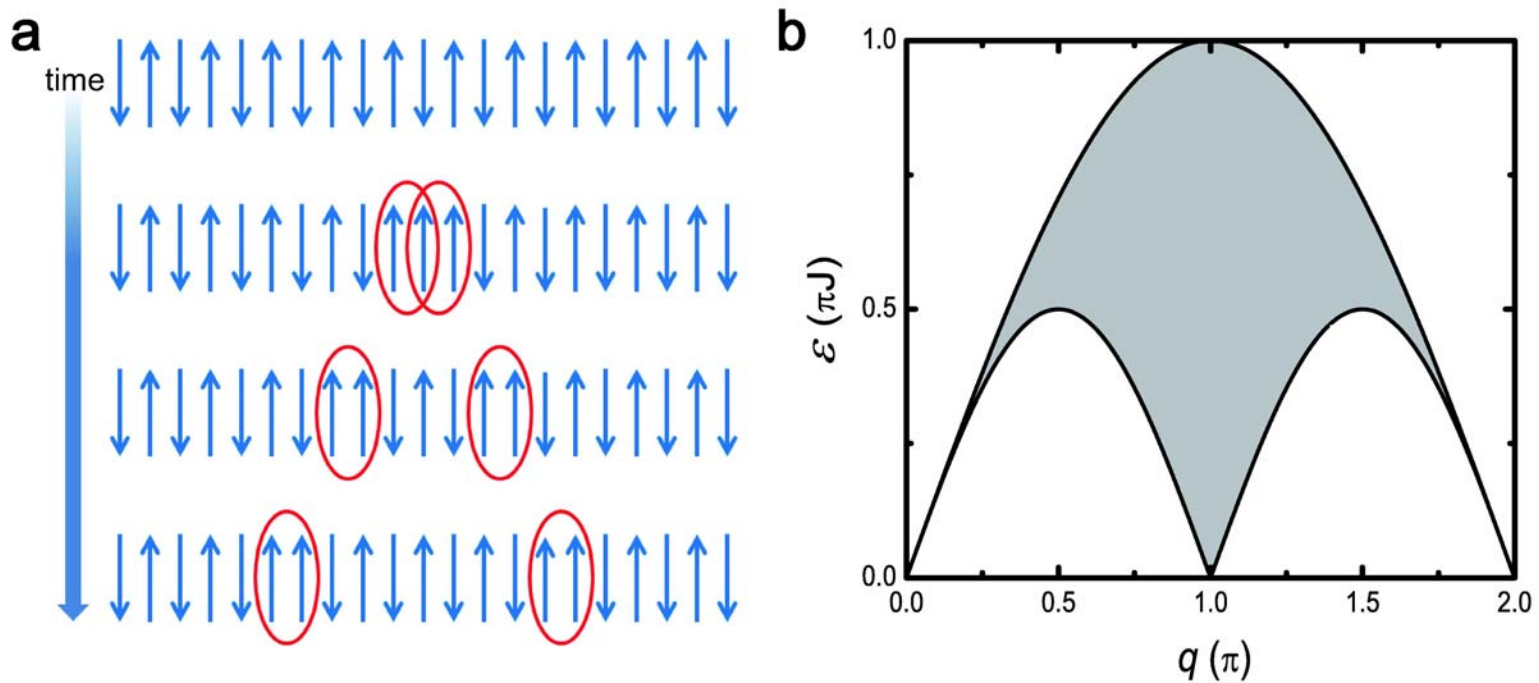


EtMe₃Sb[Pd(dmit)₂]₂: dmit-131



Significant κ_0/T : evidence for spinons in a spin-liquid candidate.

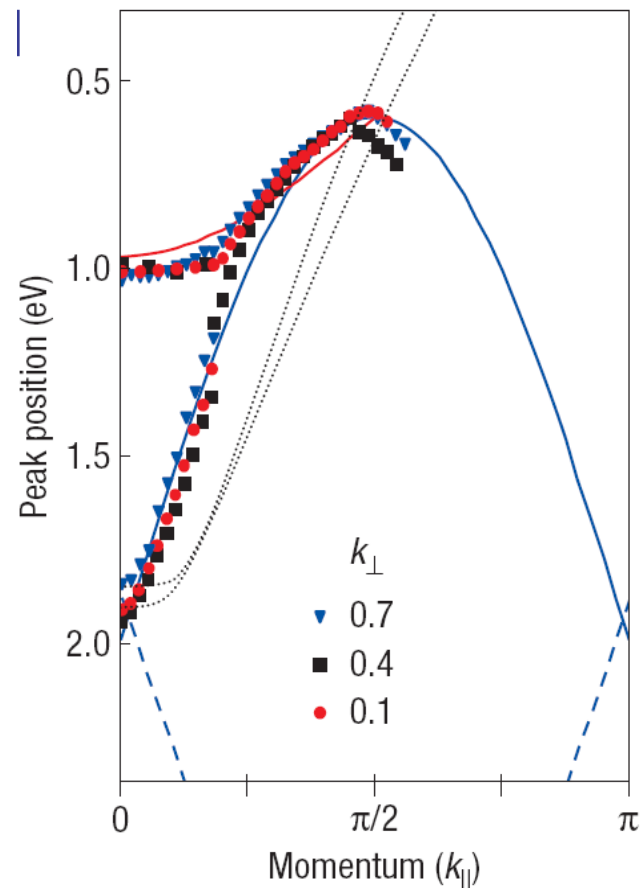
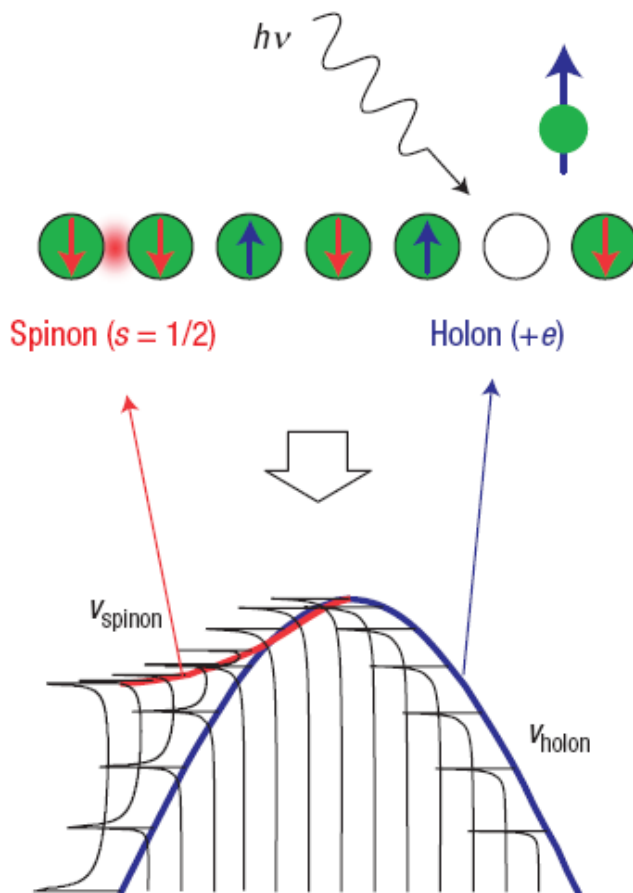
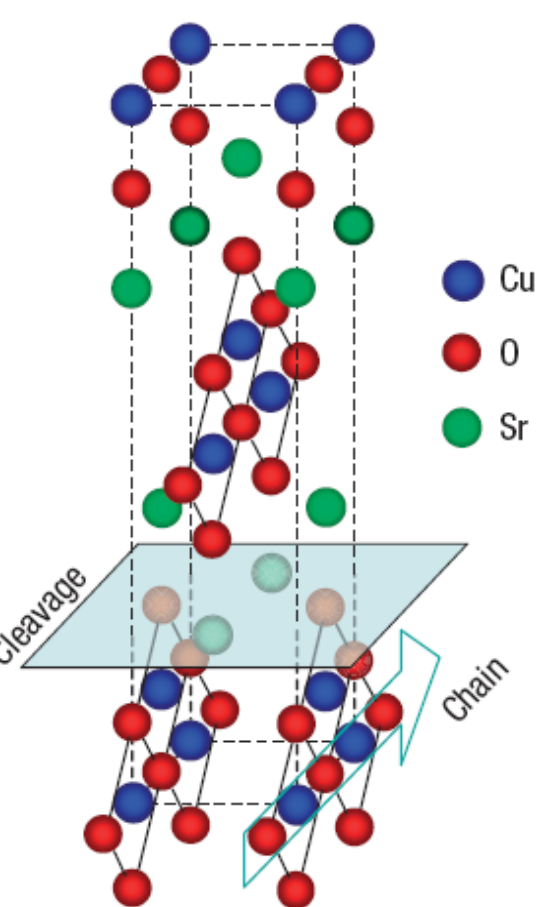
3、 Anderson localization of spinons in a spin-1/2 antiferromagnetic Heisenberg chain



The model of spin-1/2 AF Heisenberg chain can be exactly solved, and the excitations are called spinon.

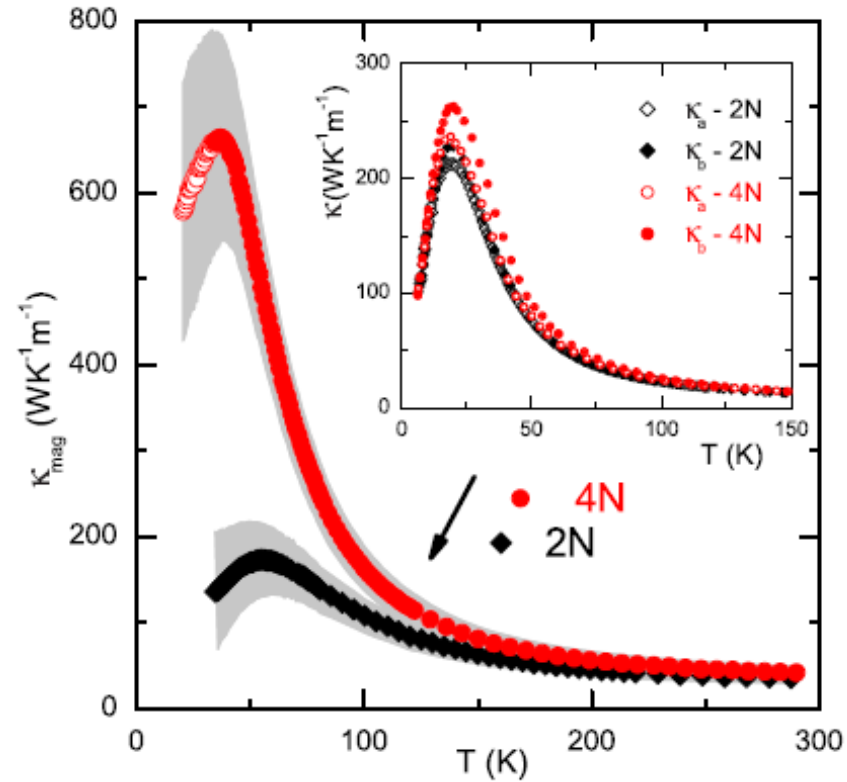
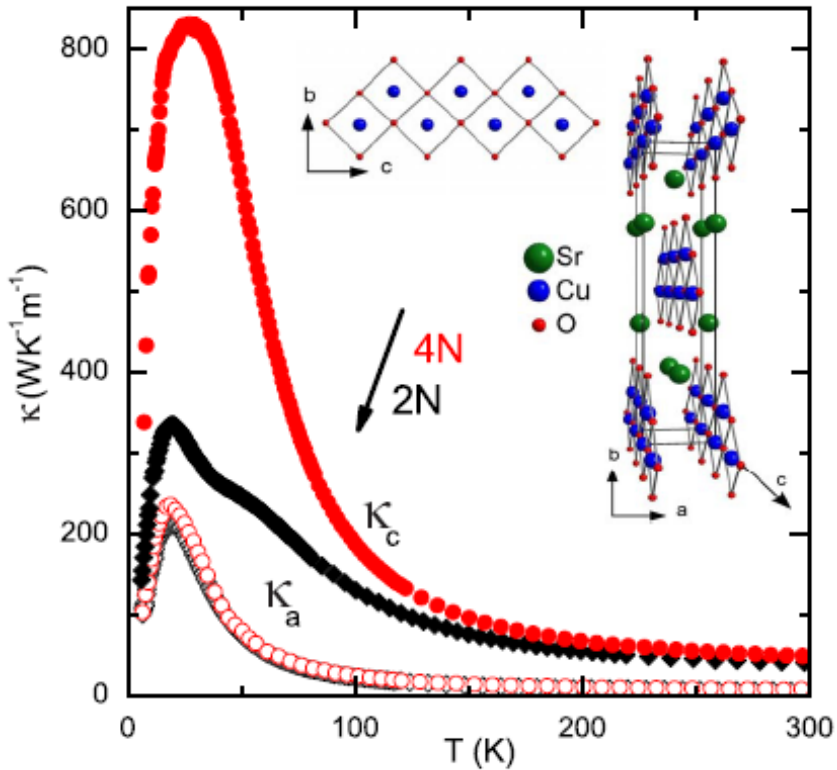
SrCuO₂:

Spin-charge separation by ARPES



SrCuO₂:

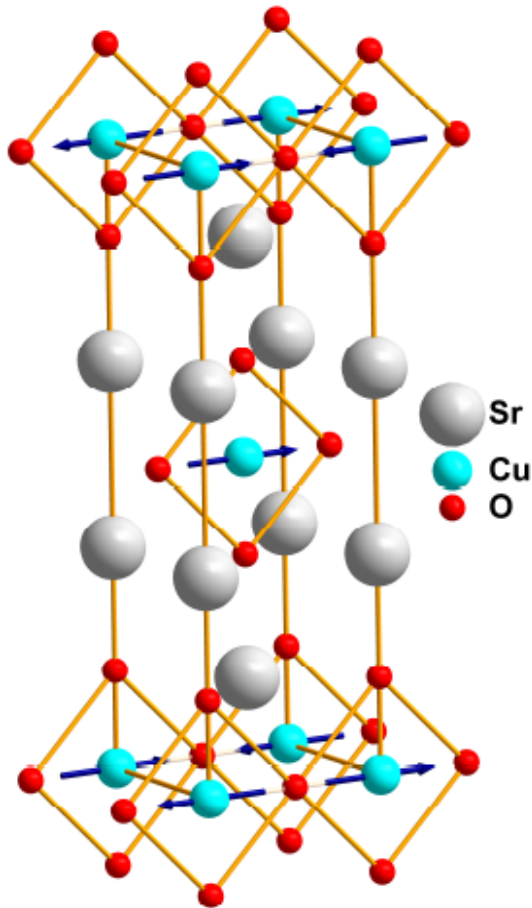
extra heat conduction along the chain



$$\kappa_{\text{spinon}} = \kappa_c - \kappa_a$$

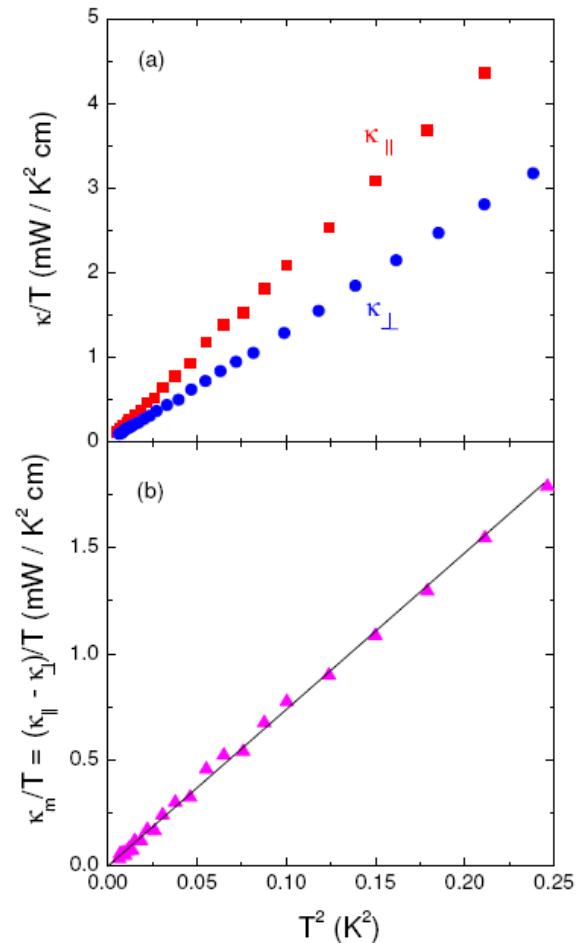
Sr_2CuO_3 :

extra heat conduction along the chain



$$J \sim 2000 \text{ K}$$

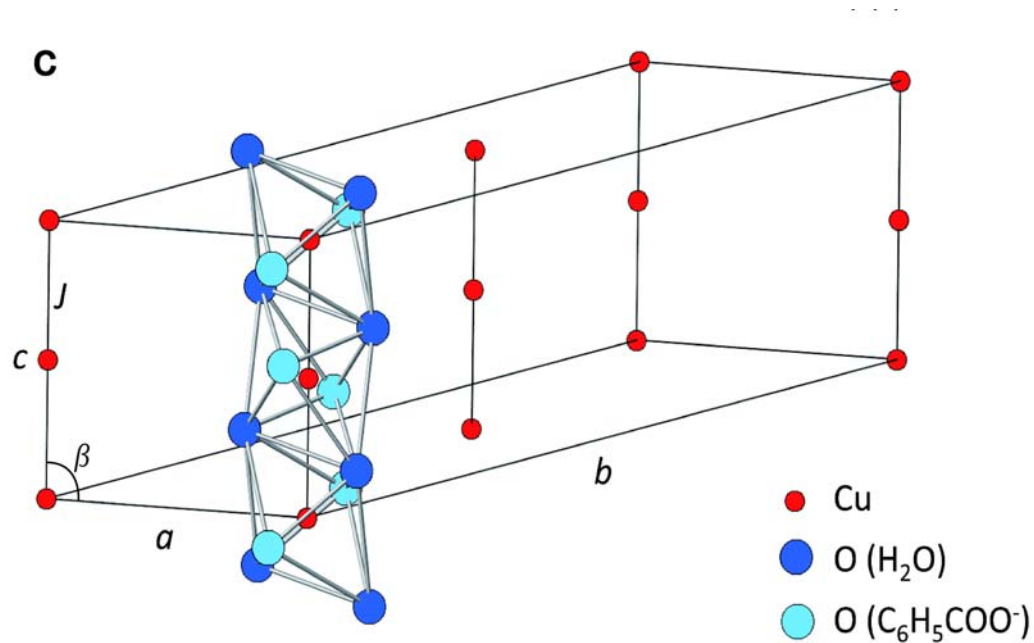
$$J' \sim T_N = 5.4 \text{ K}$$



$$\kappa_{\text{magnon}} = \kappa_{\parallel} - \kappa_{\perp}$$

Cu Benzoate:

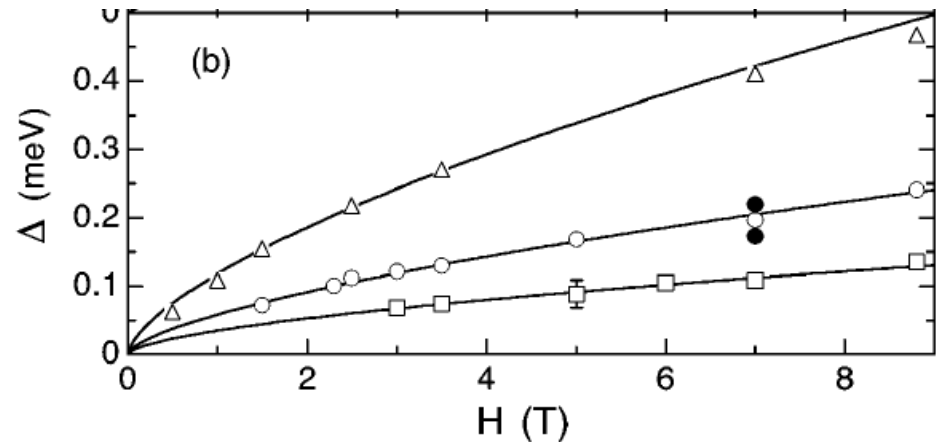
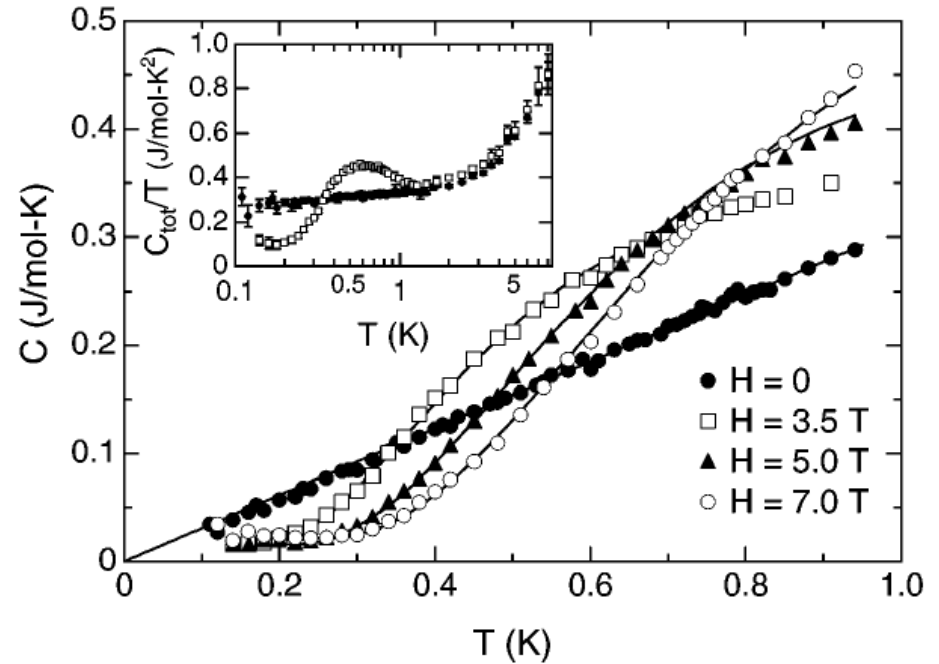
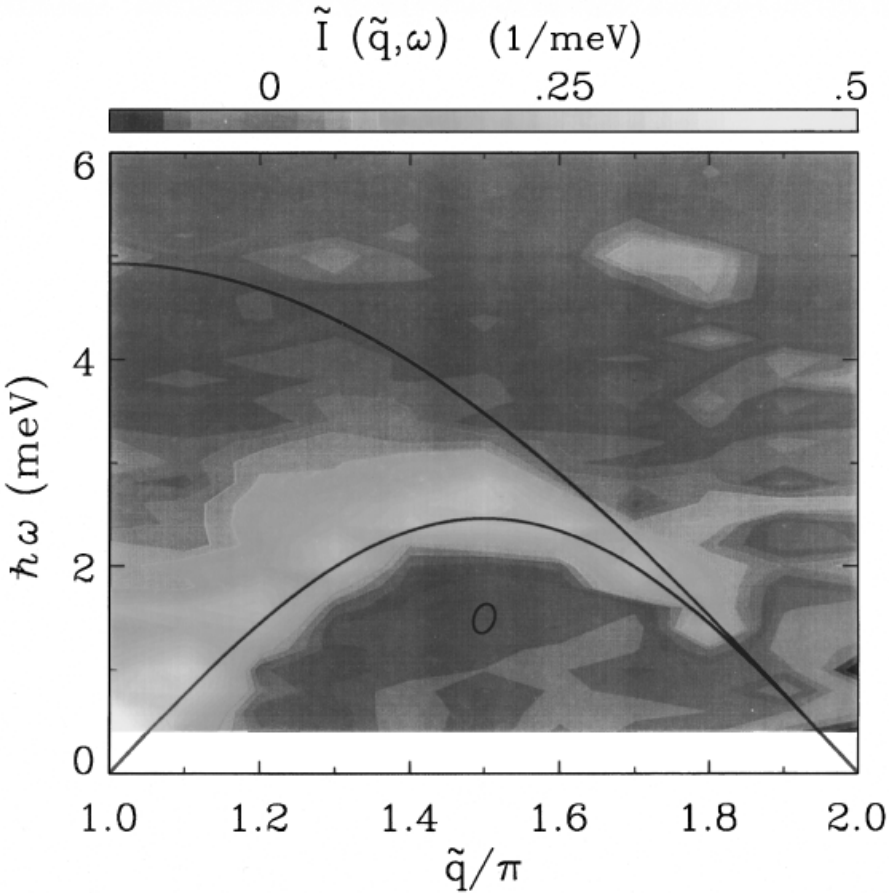
an ideal spin-1/2 Heisenberg chain



$\text{Cu}(\text{C}_6\text{H}_5\text{COO})_2 \cdot 3\text{H}_2\text{O}$: $J \sim 18.6 \text{ K}$, $J' < 50 \text{ mK}$
no order down to 50 mK

Cu Benzoate:

an ideal spin-1/2 Heisenberg chain



D. C. Dender *et al.*, PRB **53**, 2583 (1996)

D. C. Dender *et al.*, PRL **79**, 1750 (1997)

Field-Induced Gap in $S = 1/2$ Antiferromagnetic ChainsMasaki Oshikawa¹ and Ian Affleck^{1,2}

PHYSICAL REVIEW B

VOLUME 60, NUMBER 2

1 JULY 1999-II

Field-induced gap in Cu benzoate and other $S = \frac{1}{2}$ antiferromagnetic chains

Ian Affleck

PHYSICAL REVIEW B

VOLUME 59, NUMBER 22

1 JUNE 1999-II

Sine-Gordon low-energy effective theory for copper benzoate

Fabian H. L. Eßler

VOLUME 90, NUMBER 20

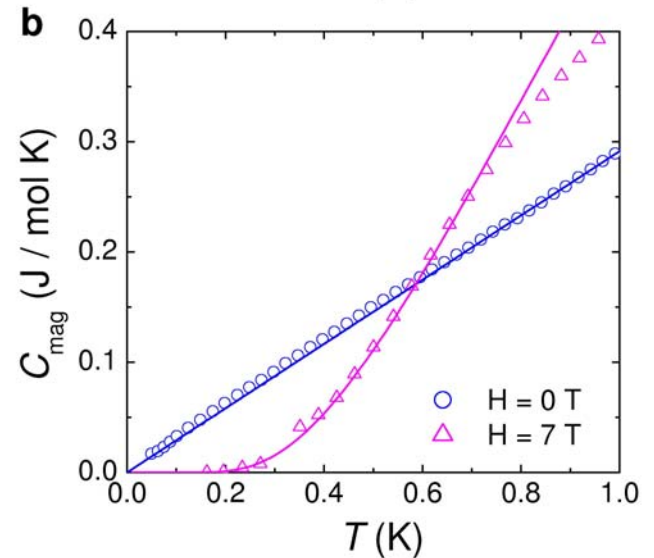
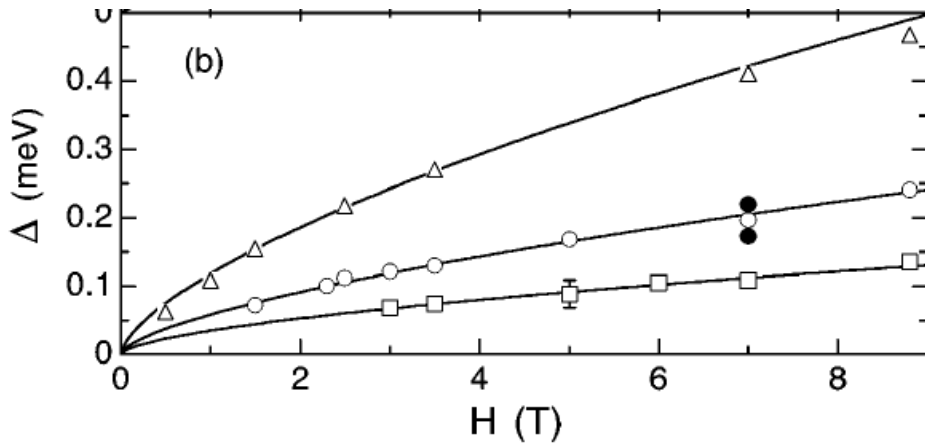
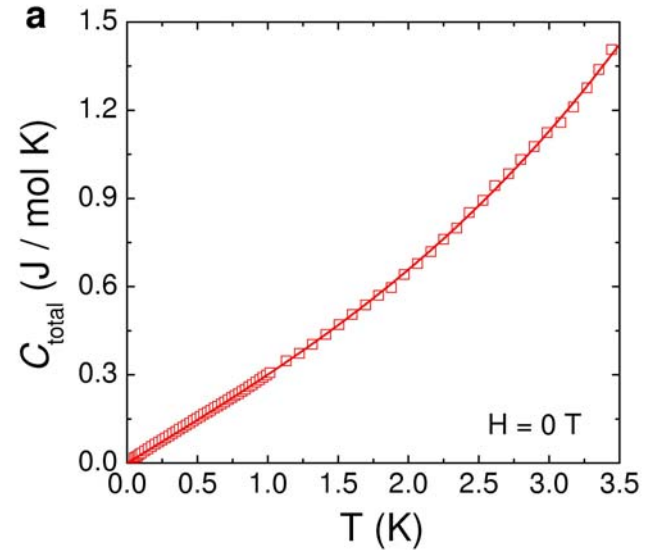
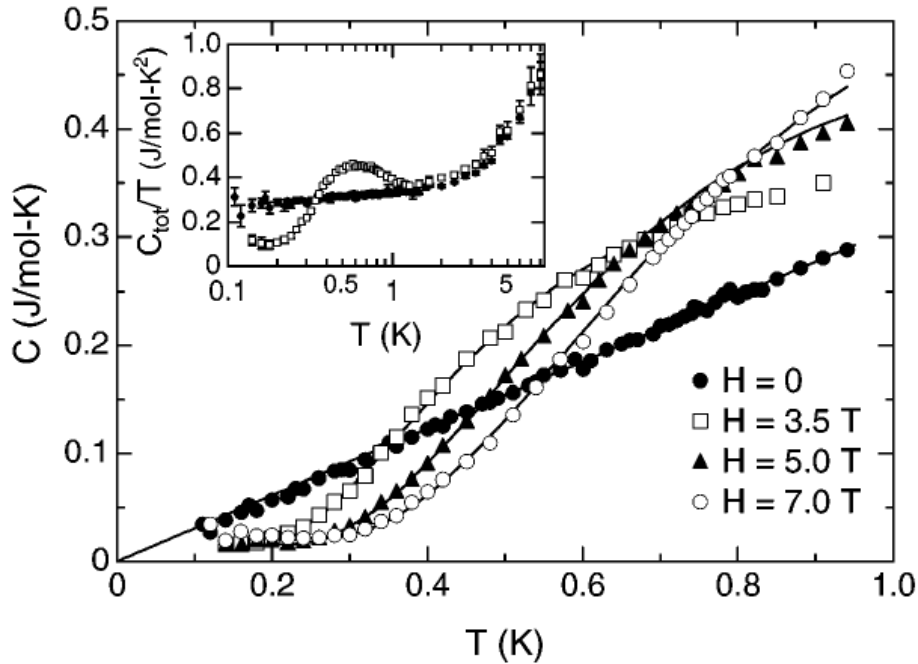
PHYSICAL REVIEW LETTERS

week ending
23 MAY 2003

Effects of the Dzyaloshinskii-Moriya Interaction on Low-Energy Magnetic Excitations in Copper BenzoateJ. Z. Zhao,¹ X. Q. Wang,^{1,2,3} T. Xiang,^{1,2} Z. B. Su,^{1,2} and L. Yu^{1,2}

Cu Benzoate:

spinon specific heat $C_s \sim T$

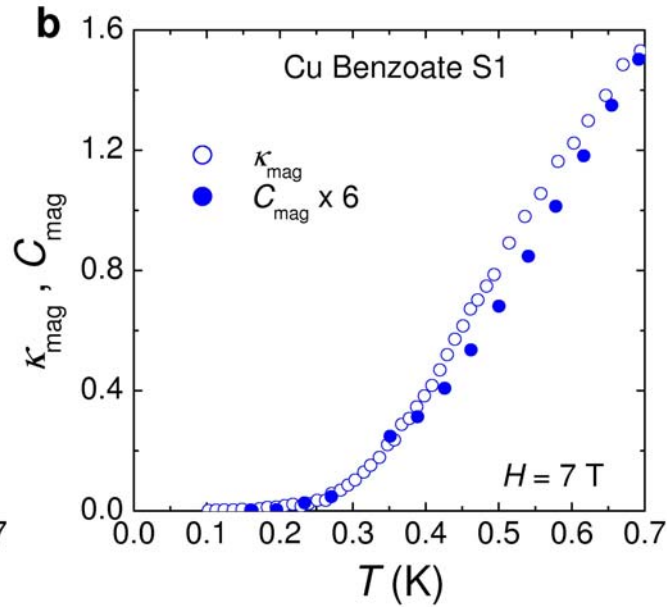
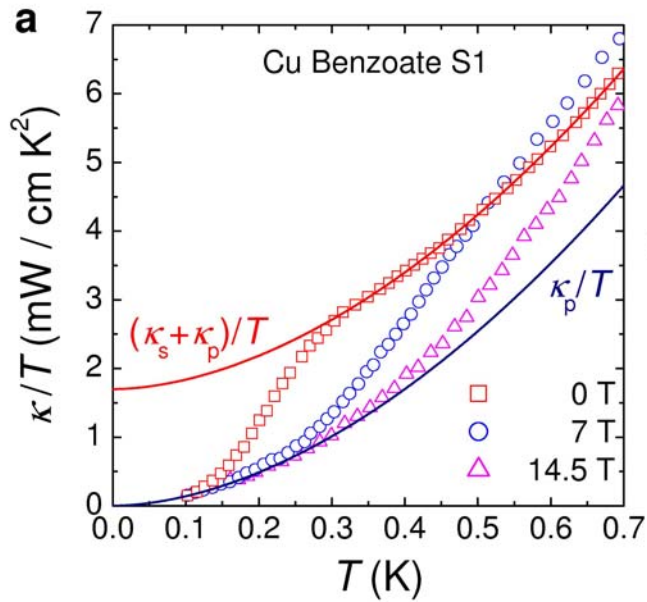


D. C. Dender *et al.*, PRL **79**, 1750 (1997)

B. Y. Pan, S. Y. Li *et al.*, arXiv:1208.3803

Cu Benzoate:

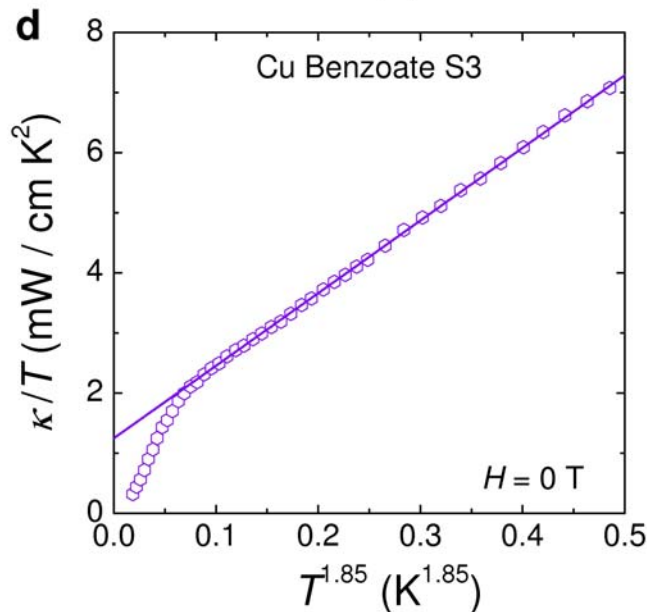
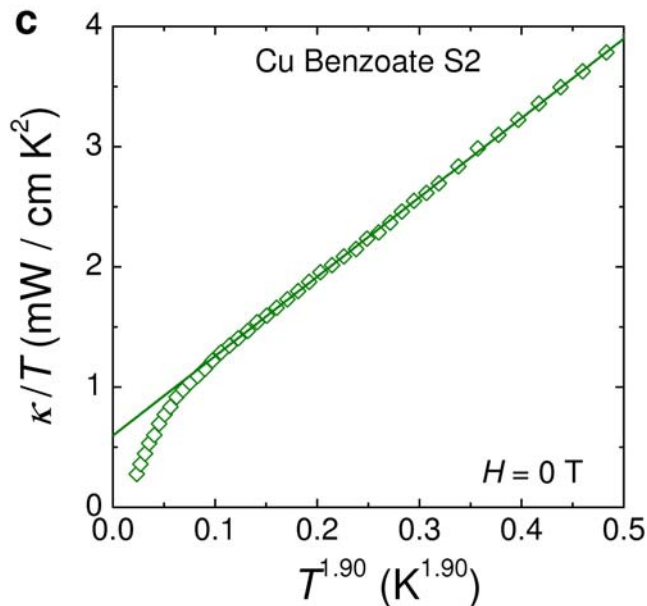
thermal conductivity



$$\kappa = C v l$$

$$H = 0 \text{ T}$$

$$\kappa_s \sim C_s \sim T$$



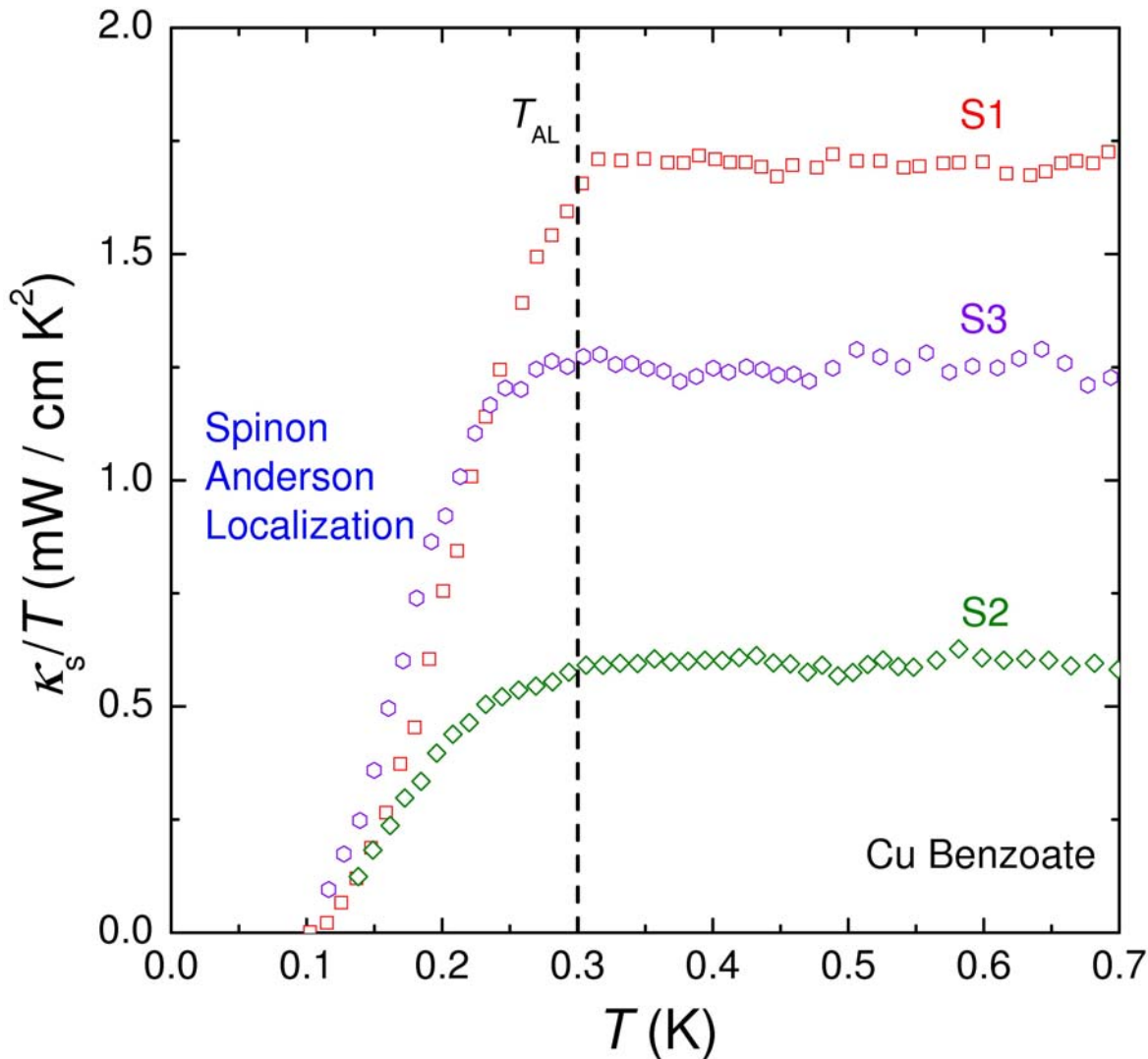
$$H = 7 \text{ T}$$

$$\kappa_{\text{mag}} \sim C_{\text{mag}}$$

B. Y. Pan, S. Y. Li *et al.*,
arXiv:1208.3803

Cu Benzoate:

spinon thermal conductivity



Linear term
Compare to electrons:

$$\kappa_0/T = L_0/\rho_0$$

$$1.7 \text{ mW/cmK}^2 \\ \sim 15 \mu\Omega \text{ cm}$$

Compare to spinons
in the spin liquid:

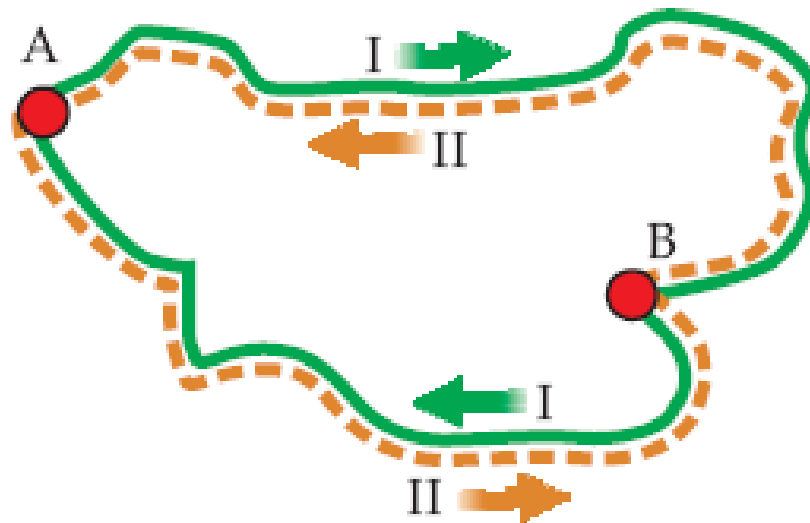
$$\kappa_s/T \approx 2 \text{ mW/cmK}^2$$

Absence of Diffusion in Certain Random Lattices

P. W. ANDERSON

Bell Telephone Laboratories, Murray Hill, New Jersey

(Received October 10, 1957)



Anderson localization of waves in disordered systems originates from interference in multiple elastic scattering.

Localization of light in a disordered medium

Nature **390**, 671 (1997)

Statistical signatures of photon localization

Nature **404**, 850 (2000)

Transport and Anderson localization in disordered two-dimensional photonic lattices

Nature **446**, 52 (2007)

Localization of ultrasound in a three-dimensional elastic network

HEFEI HU^{1*}, A. STRYBULEVYCH¹, J. H. PAGE^{1†}, S. E. SKIPETROV² AND B. A. VAN TIGGELEN²

¹Department of Physics and Astronomy, University of Manitoba, Winnipeg, Manitoba, R3T 2N2, Canada

²Université Joseph Fourier, Laboratoire de Physique et Modélisation des Milieux Condensés, CNRS, 25 Rue des Martyrs, BP 166, 38042 Grenoble, France

*Present address: Department of Physics, University of Illinois at Urbana-Champaign, Urbana, Illinois 61801-3080, USA

†e-mail: jhpage@cc.umanitoba.ca

Nature Physics 4,945 (2008)

Direct observation of Anderson localization of matter waves in a controlled disorder

Nature 453, 891 (2008)

Anderson localization of a non-interacting Bose–Einstein condensate

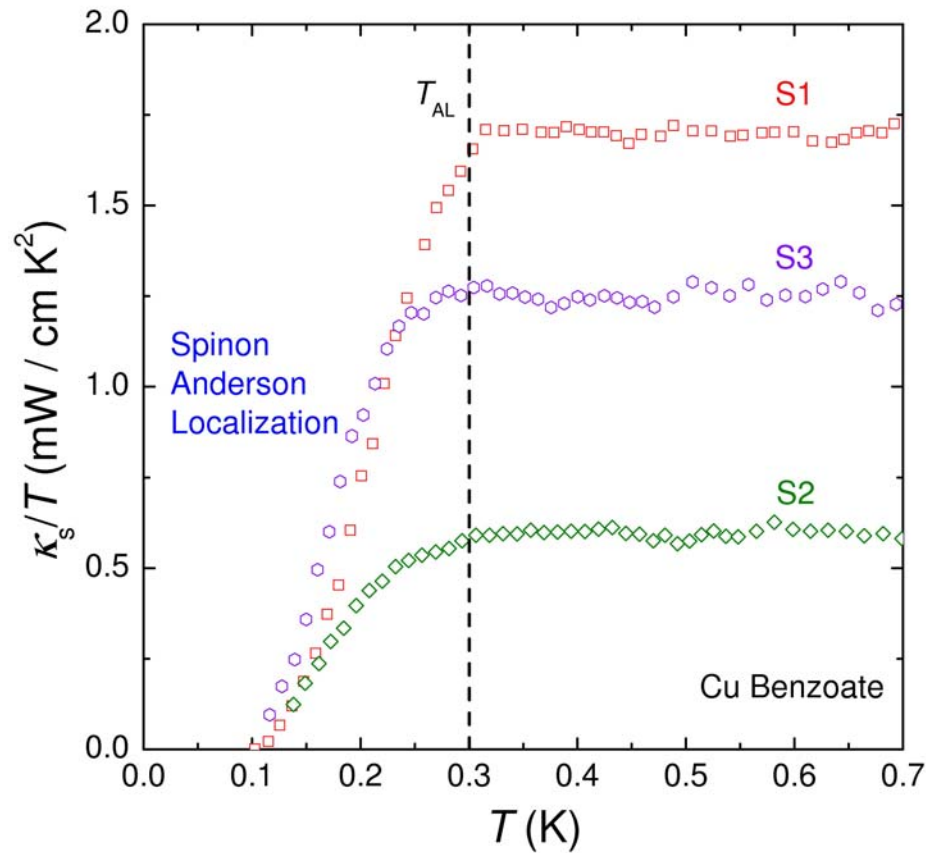
Nature 453, 895 (2008)

Three-Dimensional Anderson Localization of Ultracold Matter

Science 333, 66 (2011)

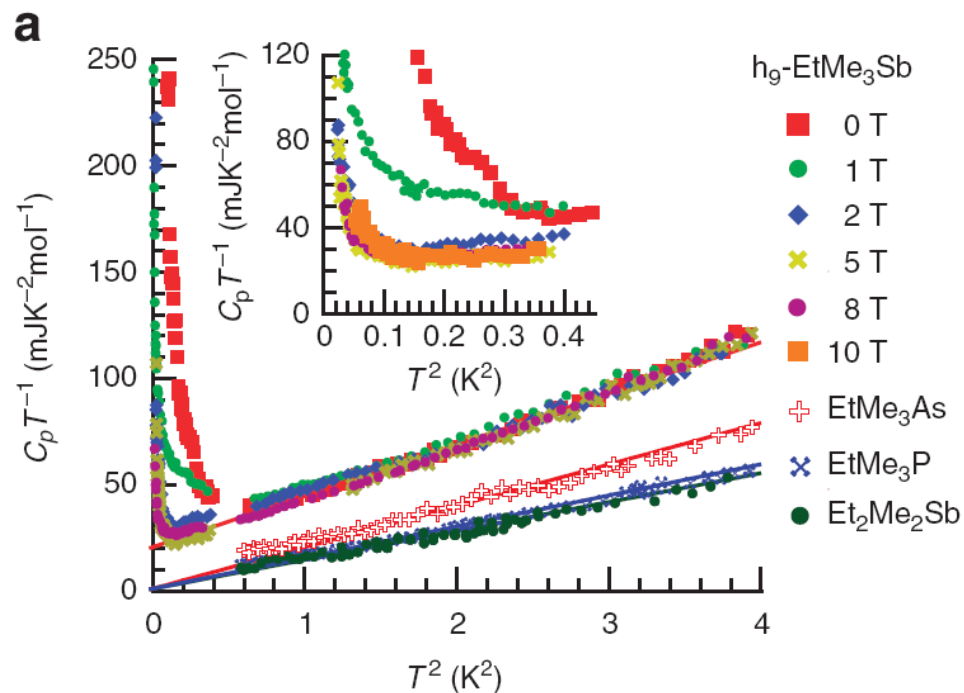
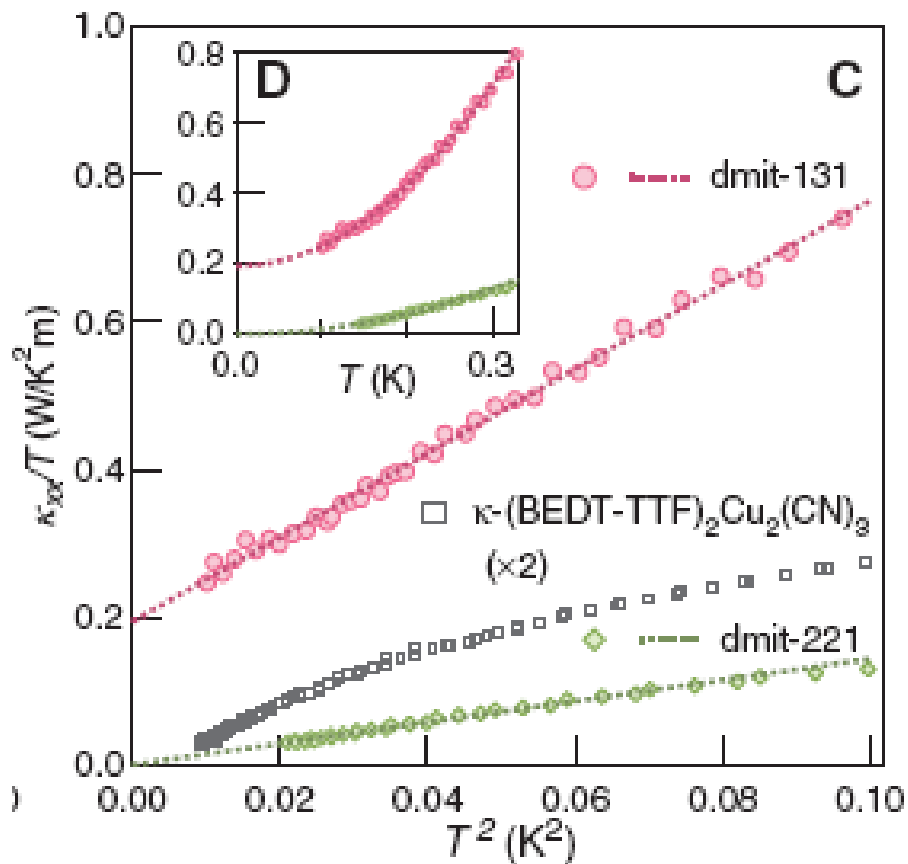
Three-dimensional localization of ultracold atoms in an optical disordered potential

Nature Physics 8, 398 (2012)



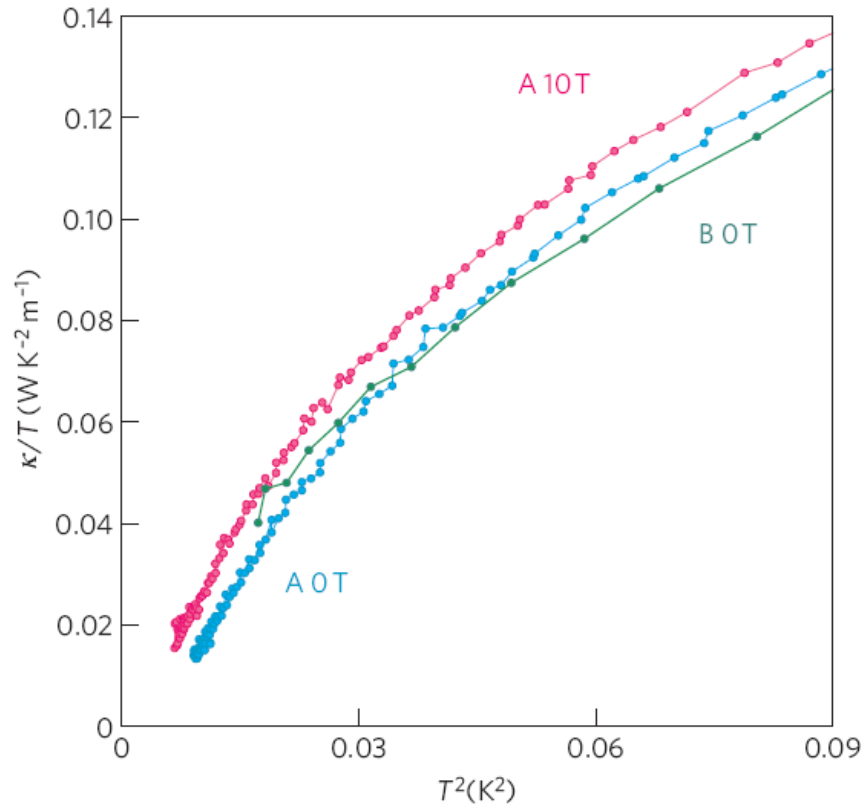
First observation of Anderson localization of magnetic excitations.
1D system is the best place for Anderson localization to occur.

EtMe₃Sb[Pd(dmit)₂]₂: dmit-131



Spin liquid:

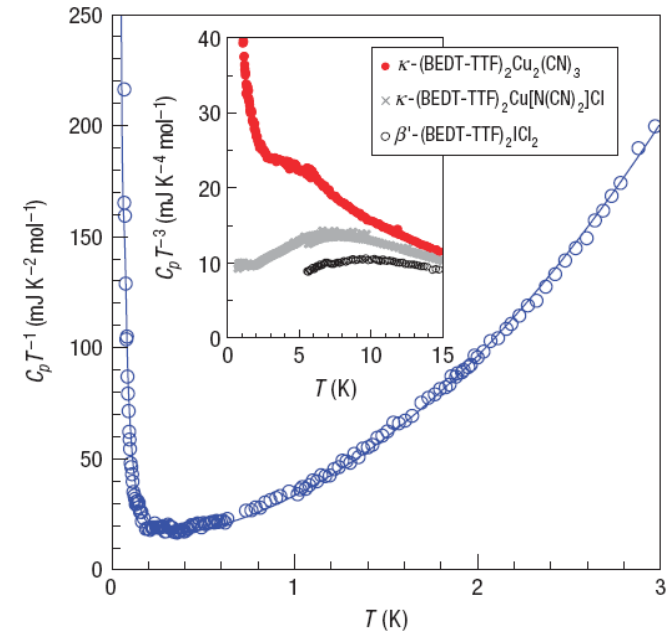
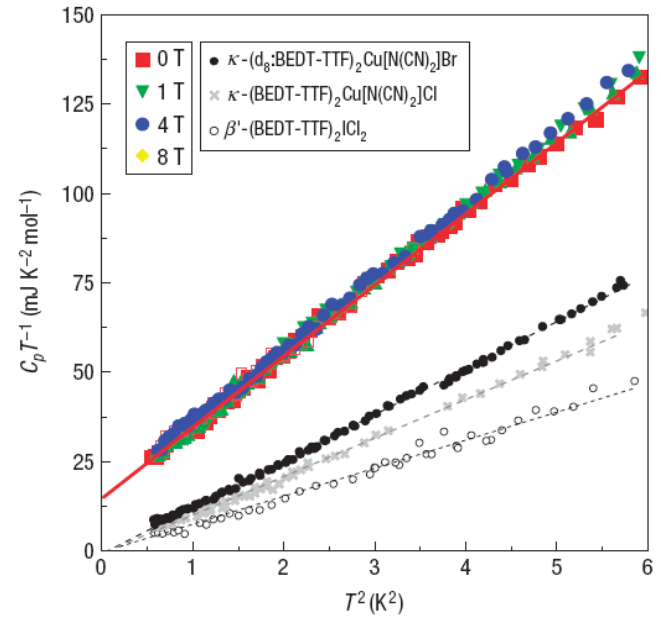
Spinons



No κ_0/T : Gapped? Localized?

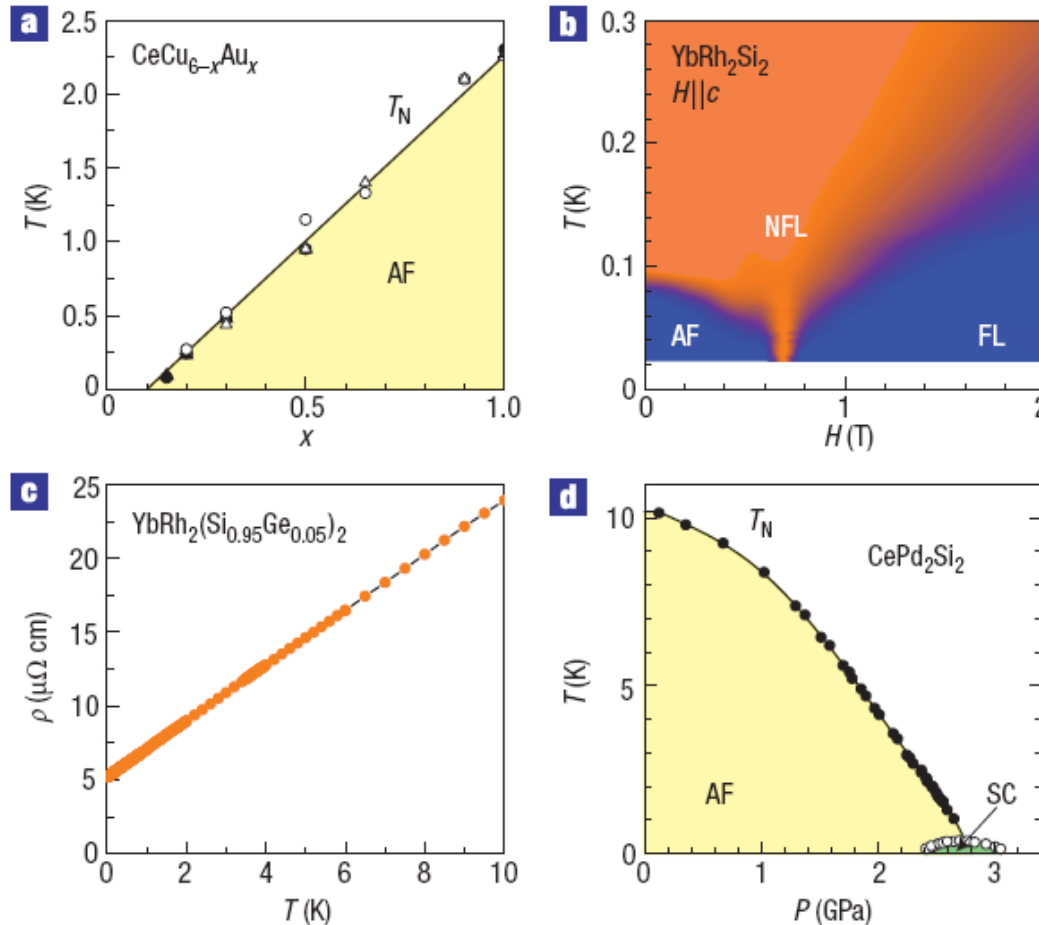
M. Yamashita *et al.*, Nature Physics **5**, 44 (2008)

S. Yamashita *et al.*, Nature Physics **4**, 459 (2008)



4. Unveiling the quantum critical point of an Ising chain in a transverse field

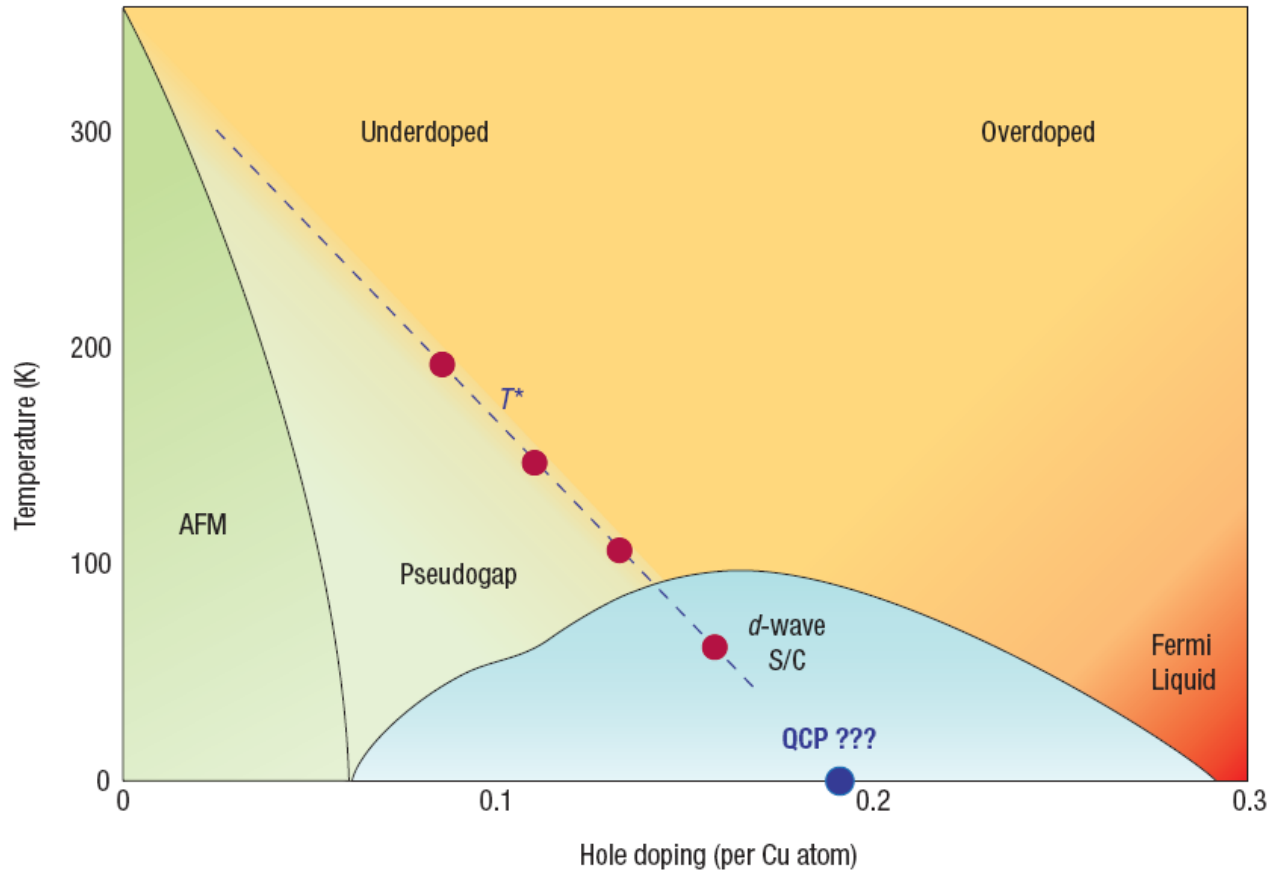
Quantum Phase Transition: big issue in condensed matter physics



Heavy-fermion systems

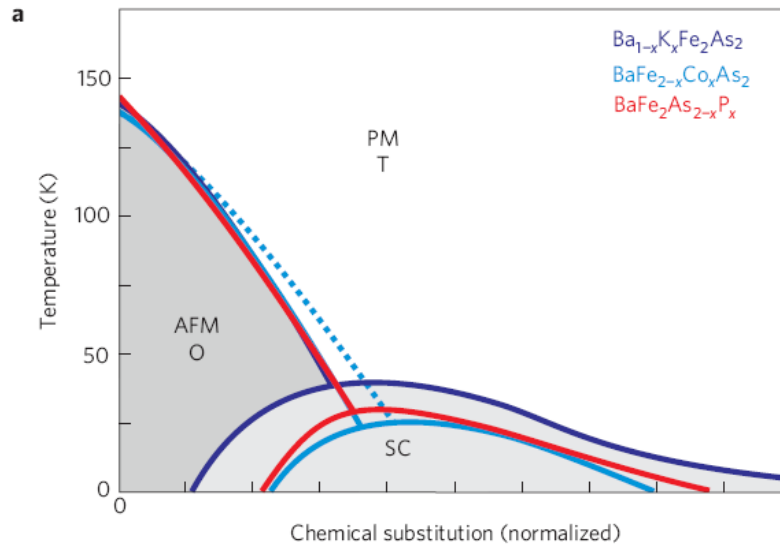
QPT occurs at zero temperature, tuned by nonthermal parameters: chemical doping, magnetic field, pressure ...

Quantum Phase Transition: big issue in condensed matter physics

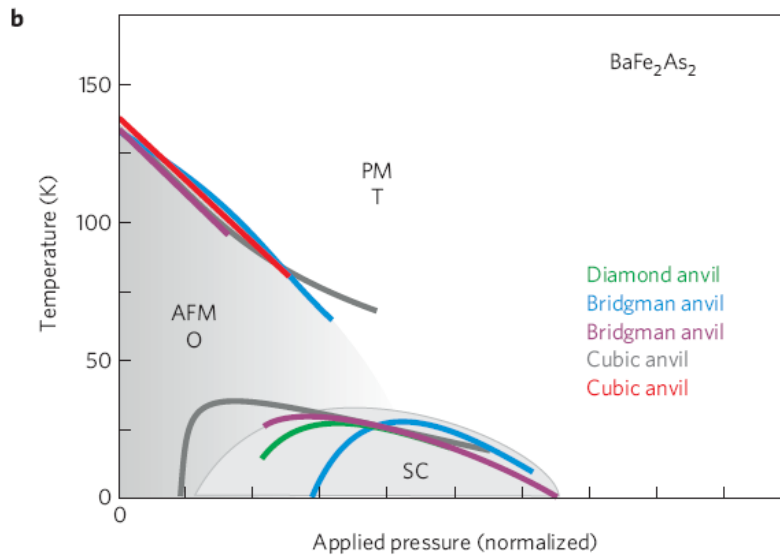


D. M. Broun, Nature Phys. **4**, 170 (2008)

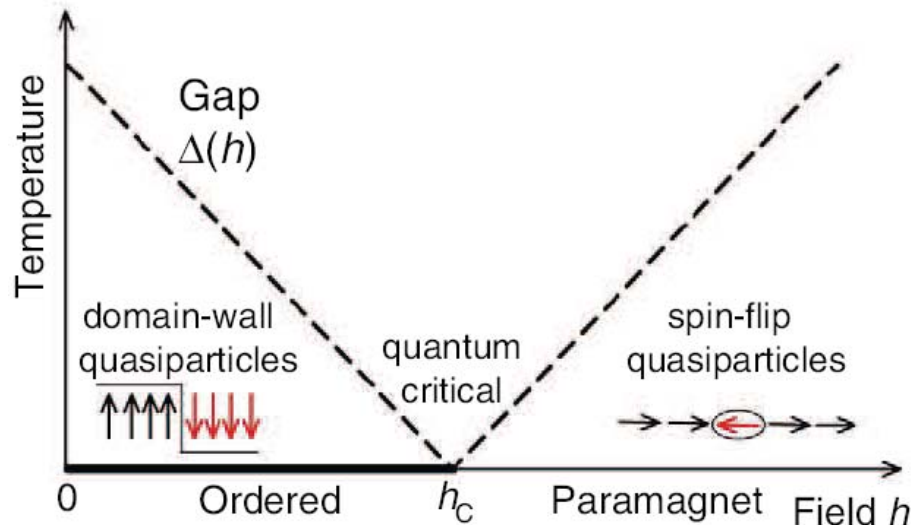
Quantum Phase Transition: big issue in condensed matter physics



Iron pnictides



TFIC: a relatively simple model undergoing QPT



Hamiltonian:

$$H = -J \sum_i (\hat{\sigma}_i^z \hat{\sigma}_{i+1}^z - h \hat{\sigma}_i^x)$$

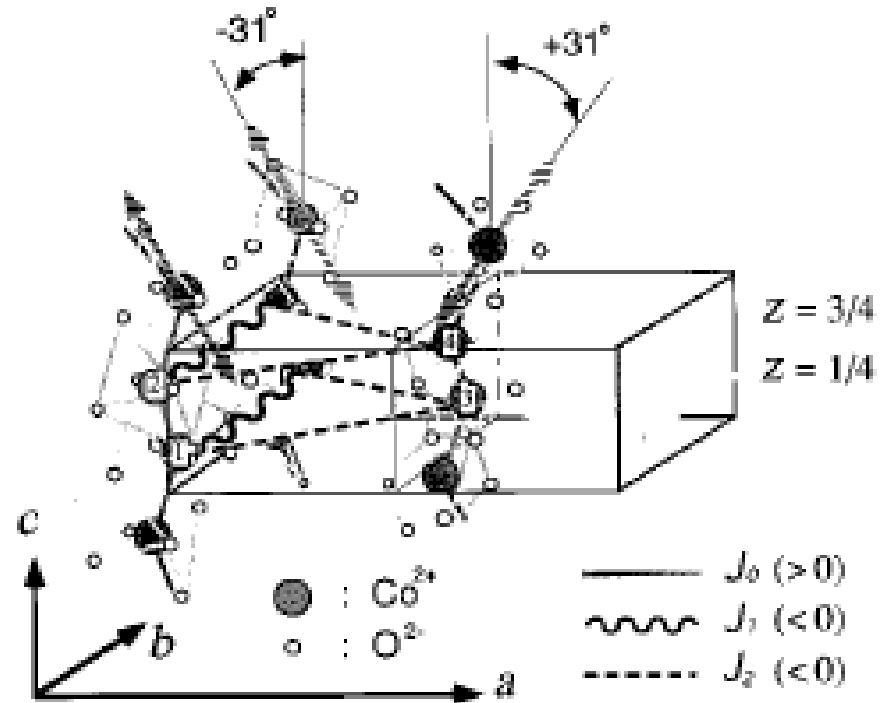
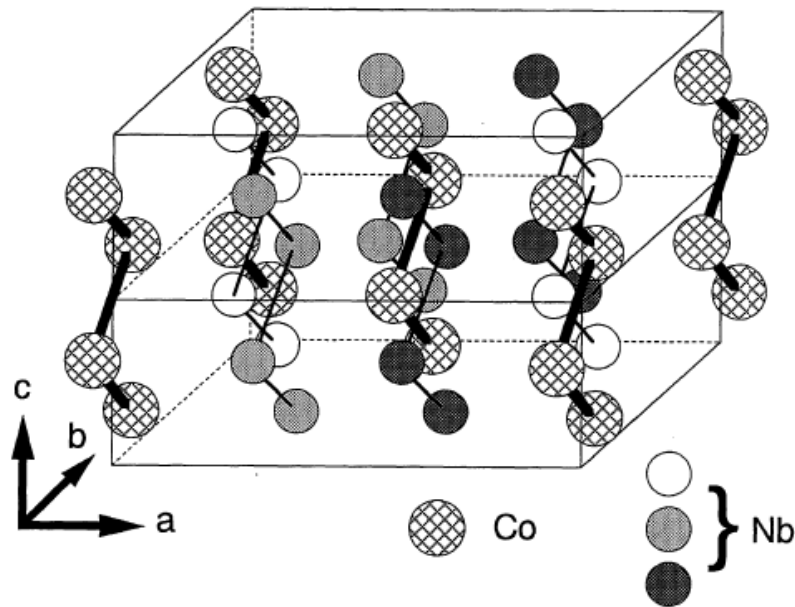
The Ising chain in a transverse field (TFIC):
one of the most-studied model in condensed matter physics.

By using the Jordan-Wigner transformation, the spins can be transformed to noninteracting spinless fermions, and this model can be exactly solved.

The minimum single-particle excitation energy, or the energy gap: $\Delta = 2J|1-h|$

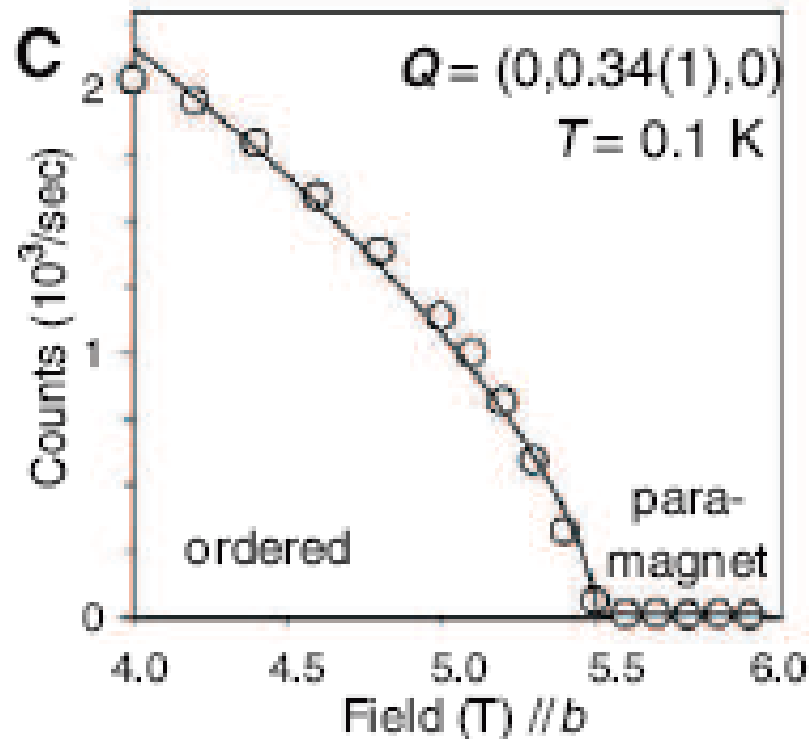
Quantum critical point: $h = 1, \Delta = 0$

CoNb₂O₆: a rare experimental realization of the TFIC model

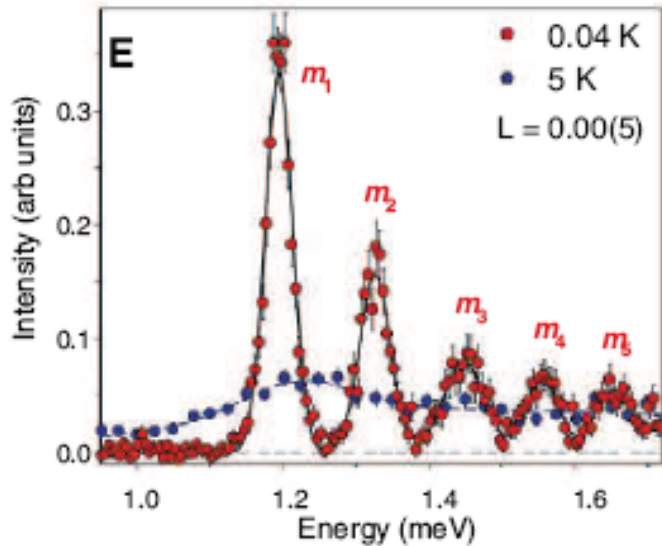
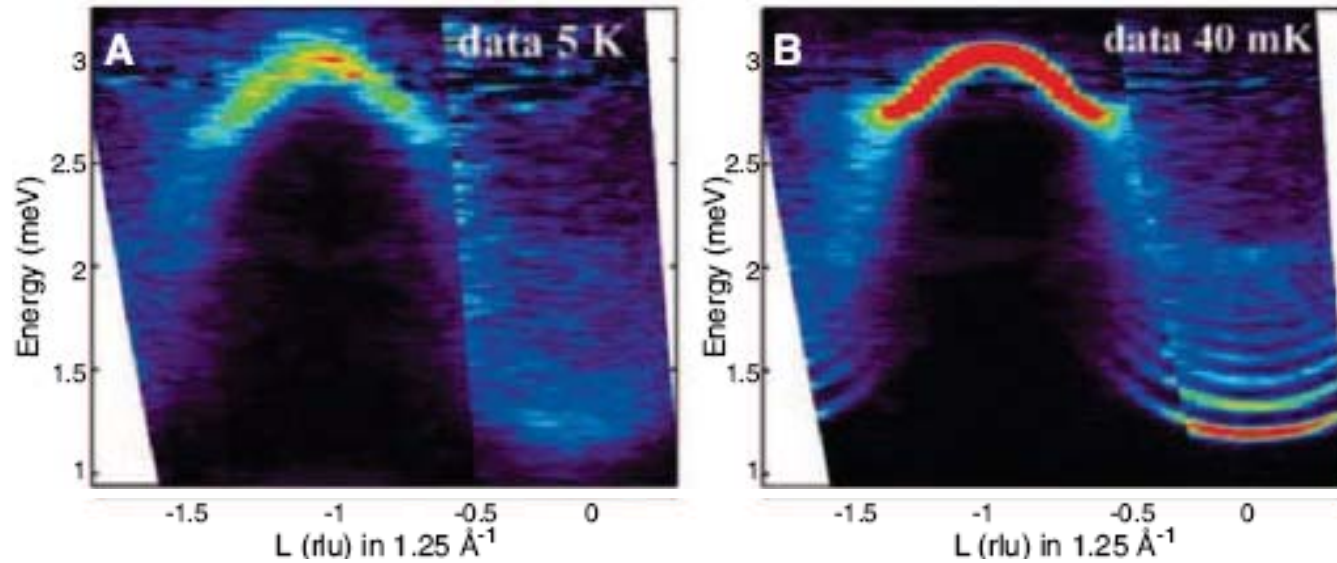


Strong easy-axis anisotropy due to CFEs: easy-axis in *ac* plane, $\pm 31^\circ$ to *c*-axis
Intrachain coupling $J > 0$: favors FM ordering along *c*-axis
Interchain coupling $J_1, J_2 < 0, J_1, J_2 \ll J$: favors AF ordering between chains

CoNb₂O₆: neutron scattering experiments in a transverse field



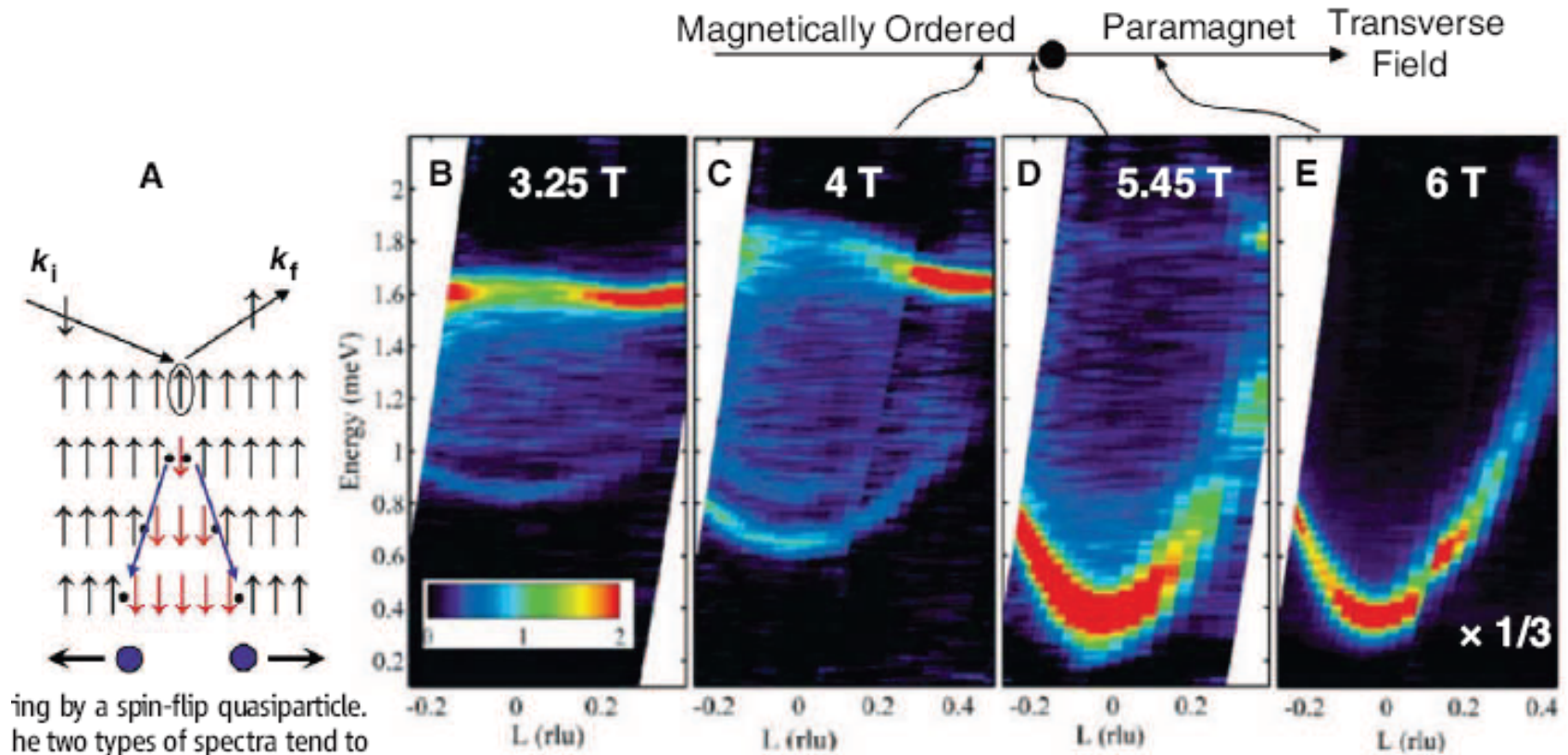
Elastic scattering in H || b: QPT at H = 5.5 T.



Inelastic scattering in H = 0 and at 40 mK:
a few bound states m₁, m₂, m₃, ...
(domain-wall quasiparticles)

CoNb₂O₆:

neutron experiments in a transverse field



ing by a spin-flip quasiparticle.
he two types of spectra tend to

domain-wall
quasiparticles



Inelastic scattering at 0.1 K:
domain-wall quasiparticles for $H < 5.5$ T
spin-flip quasiparticles for $H > 5.5$ T.

spin-flip
quasiparticles



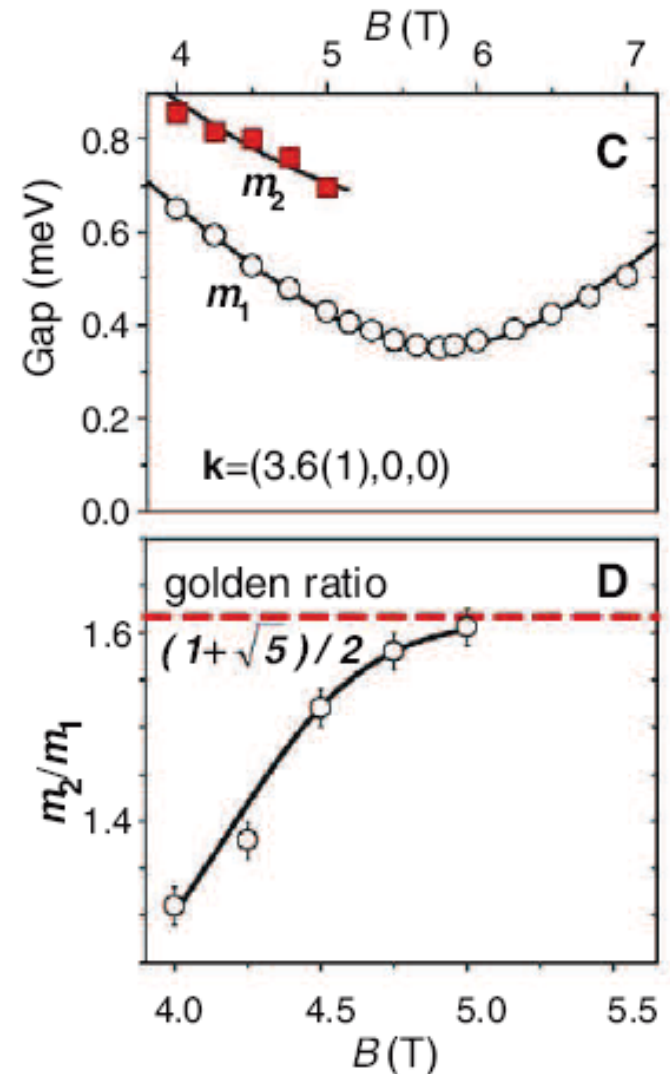
Our idea: Probe the low-lying magnetic excitation in CoNb_2O_6

In a quasi-1D system such as CoNb_2O_6 with finite interchain couplings, a complete gap softening is only expected (23) at the location of the 3D magnetic long-range order Bragg peaks, which occur at a finite interchain wave vector \mathbf{q}_\perp that minimizes the Fourier transform of the antiferromagnetic interchain couplings; the measurements shown in Fig. 4C were in a scattering plane where no magnetic Bragg peaks occur, so an incomplete gap softening would be expected here, as indeed was observed.

$$\text{QCP: } \Delta = 0$$

Technical difficulties for neutron scattering to probe the QCP with $\Delta = 0$.

Heat transport should be able to detect the low-energy quasiparticles near the QCP.

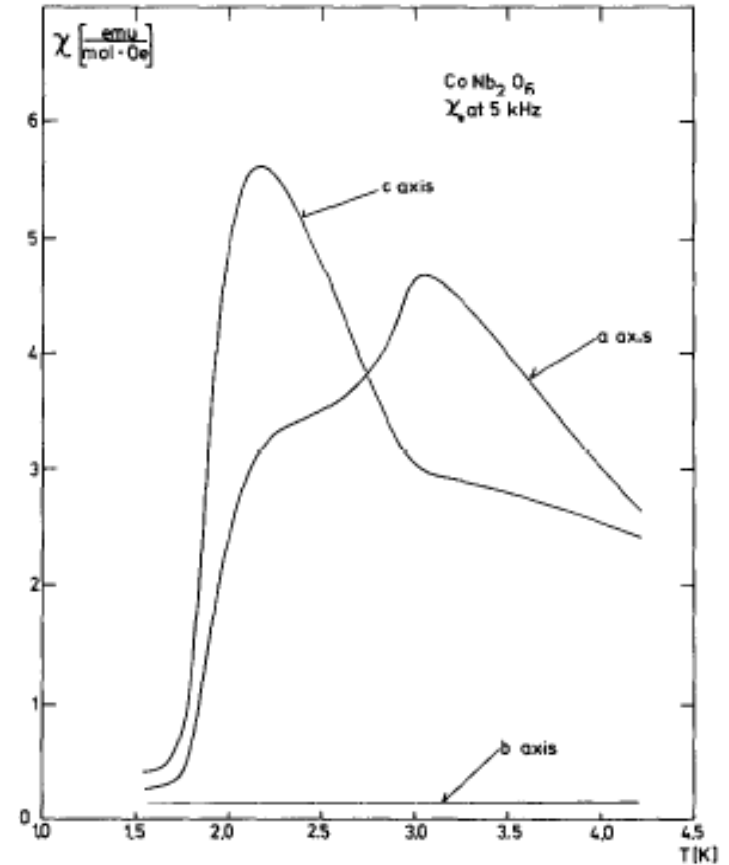
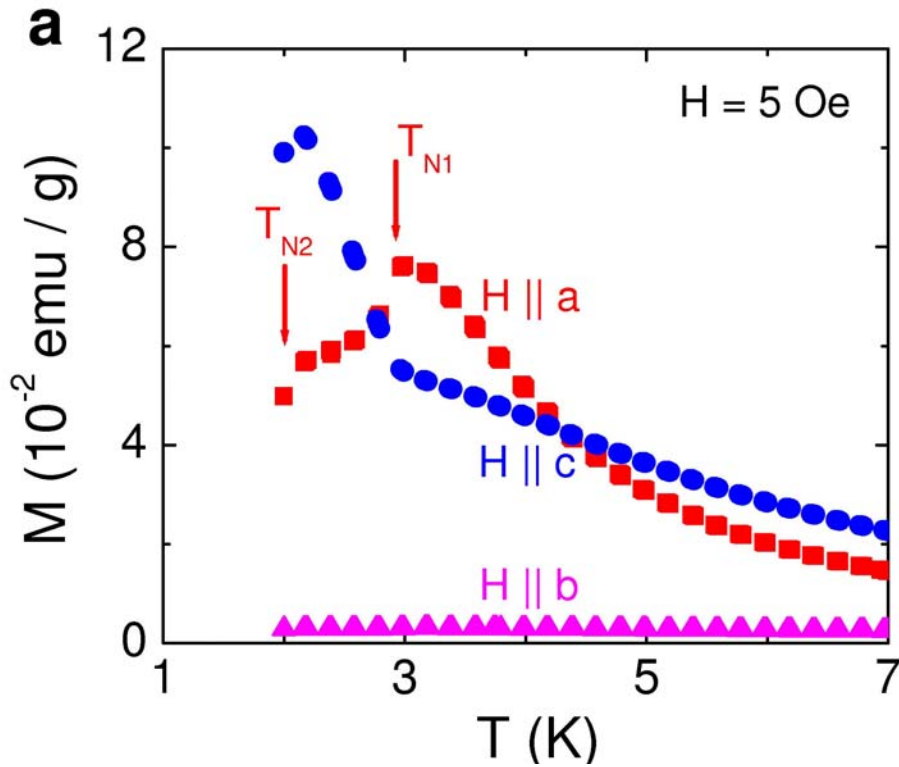


CoNb₂O₆:

Single crystal growth



Floating-zone
optical furnace

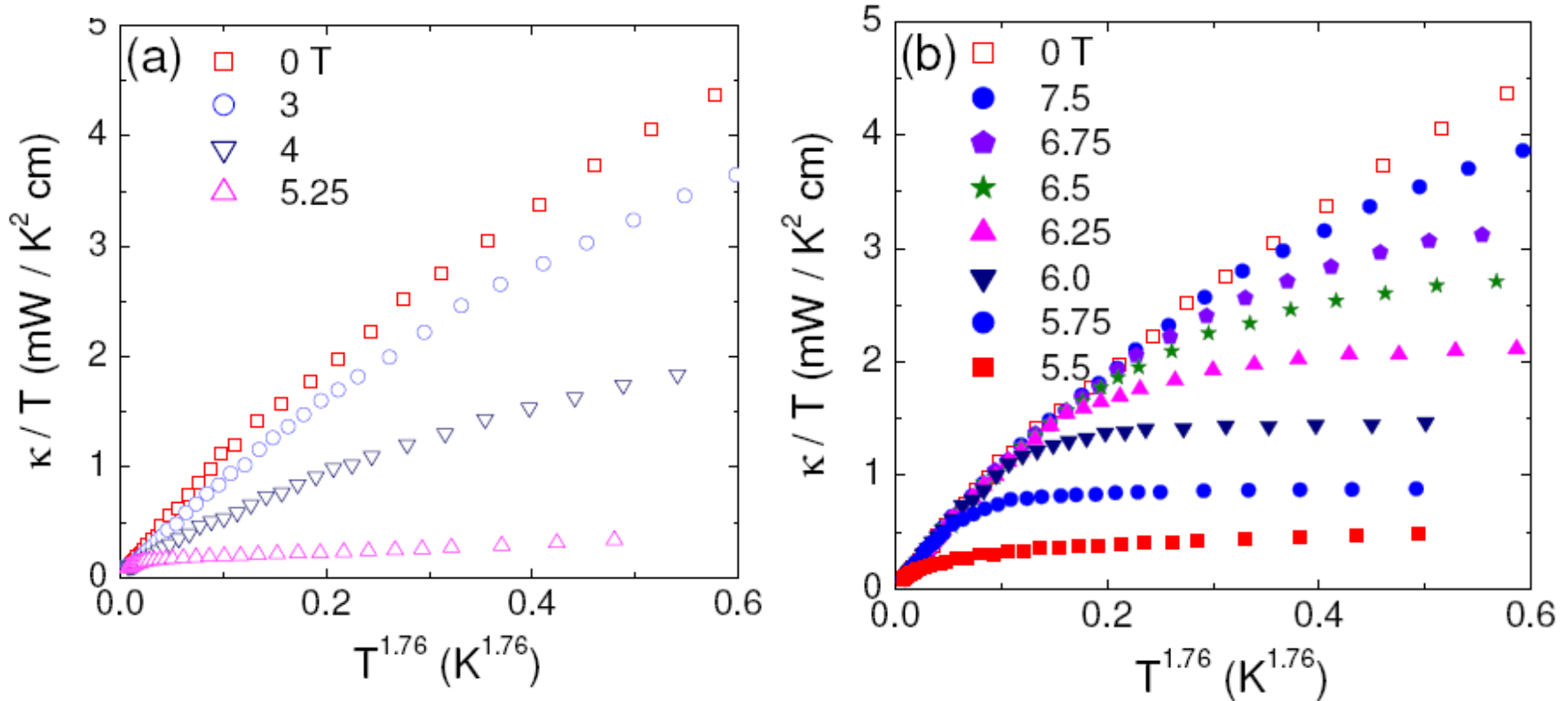


The interchain couplings: two 3D transitions

$T_{N1} = 2.95 \text{ K}$: incommensurate SDW transition

$T_{N2} = 1.97 \text{ K}$: commensurate AF transition

W. Scharf *et al.*, JMMM 13, 121 (1979)



- 1) No significant positive contribution to κ/T by magnetic excitations, likely due to low J . The suppression of κ/T is due to the scattering of phonons by these magnetic excitations.
- 2) At the left of QCP, there are some gapless excitations (AF magnons?).
- 3) At the QCP, some gapless excitations strongly scatter phonons.
- 4) At the right of QCP, the gap develops with increasing magnetic field.

Summary

Low-T thermal conductivity and specific heat are nice tools to probe low-lying magnetic excitations in quantum magnets:

- 1、 AFM magnons in 3D Nd_2CuO_4 : $\kappa_m \sim T^3$
- 2、 FM magnons in 3D YIG: $C_m \sim T^{1.5} + \text{correction}$; $\kappa_m \sim T^2 + \text{corrections}$
- 3、 Spinons in 2D spin liquid: $C_s \sim T$; $\kappa_s \sim T$
- 4、 Spinons in 1D spin-1/2 Heisenberg chain Cu Benzoate:
 $C_s \sim T$ down to 50 mK, $\kappa_s \sim T$ down to 300 mK, observing Anderson localization of spinons at lower temperature.
- 5、 Quasi-1D Ising chain CoNb_2O_6 in a transverse field: the magnetic excitations strongly scatter phonons, which unveils the QCP.