

Quantum Impurities

A Challenge for Quantum Simulation with Ultracold Atoms

Richard Schmidt

- *Harvard University / ITAMP* -

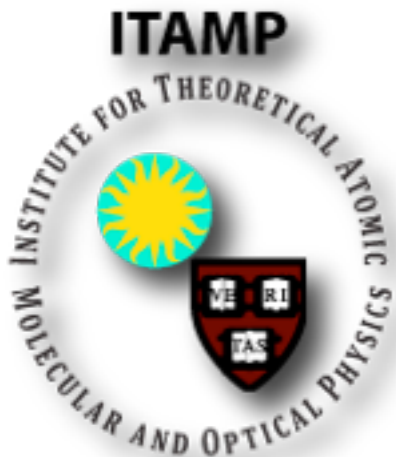
Theoretical Physics Seminar
Tsinghua University, Institute for Advanced Study
Beijing, 10/29/14

Collaborators

Harvard/ITAMP: David Benjamin, Eugene Demler,
Mikhail Lukin, Hossein Sadeghpour

Munich: Steffen P. Rath, Wilhelm Zwerger

Vienna: Mikhail Lemeshko



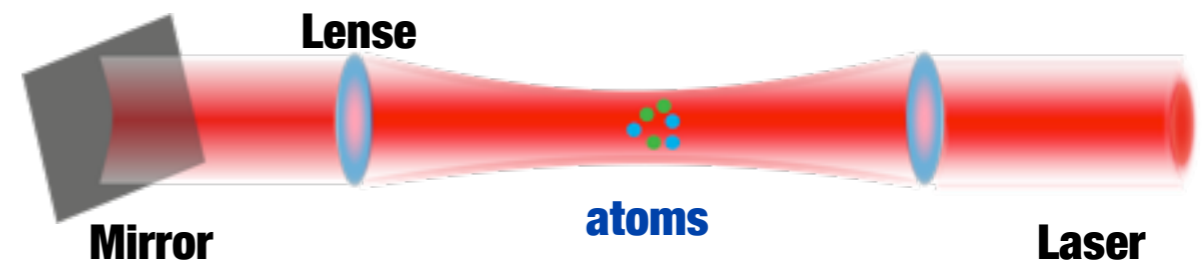
Harvard University



ultracold quantum gases

atoms trapped by laser in harmonic confinement

→ cooled to very low temperatures ~ 100 nK



condensed matter system with **well controlled Hamiltonian**

e.g.

$$H = \sum_{\mathbf{p}} \epsilon_{\mathbf{p}} \hat{c}_{\mathbf{p}}^{\dagger} \hat{c}_{\mathbf{p}} + g \sum_{\mathbf{k}, \mathbf{k}', \mathbf{q}} \hat{c}_{\mathbf{k}' - \mathbf{q}}^{\dagger} \hat{c}_{\mathbf{k} + \mathbf{q}}^{\dagger} \hat{c}_{\mathbf{k}} \hat{c}_{\mathbf{k}'}$$

low scattering energies: *contact* interactions

→ interaction strength tunable

→ experimental accessible:

- density
- transport coefficients (viscosity, spin diffusion...)
- responses, correlations

**quantum simulator of Hamiltonians
&
testground for few and many-body theories**

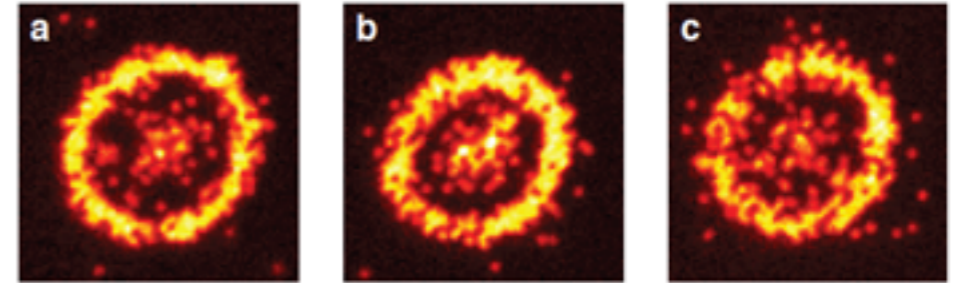
ultracold quantum gases

cold atoms as quantum simulators

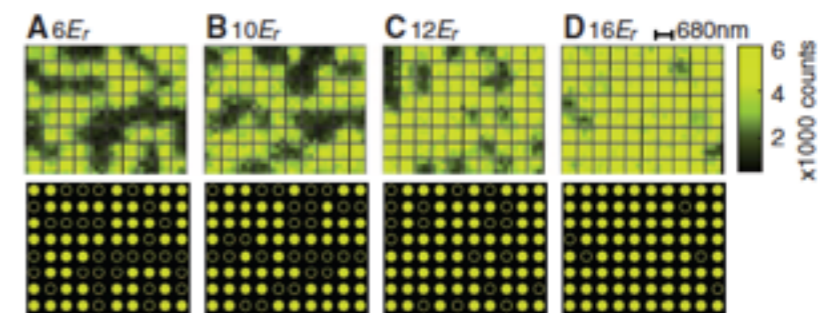
optical lattices / single-site detection:

- ▶ Mott-Insulator to Superfluid transition

MOTT-SUPERFLUID TRANSITION

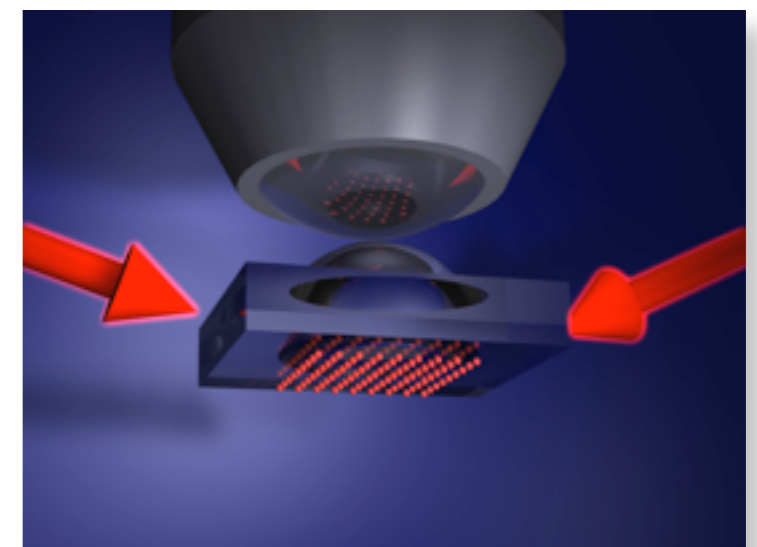


SHERSON ET AL., NATURE 467 (2010)



BAKR ET AL. SCIENCE 329 (2010)

QUANTUM GAS MICROSCOPE



BAKR ET AL. SCIENCE 329 (2010)

IMAGE: CUAWEB.MIT.EDU - GREINER GROUP

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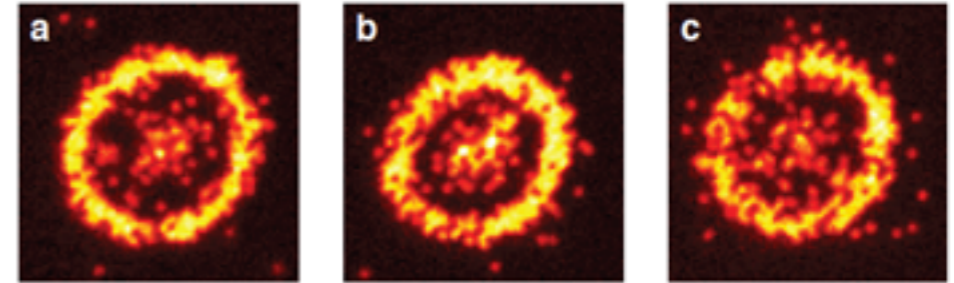
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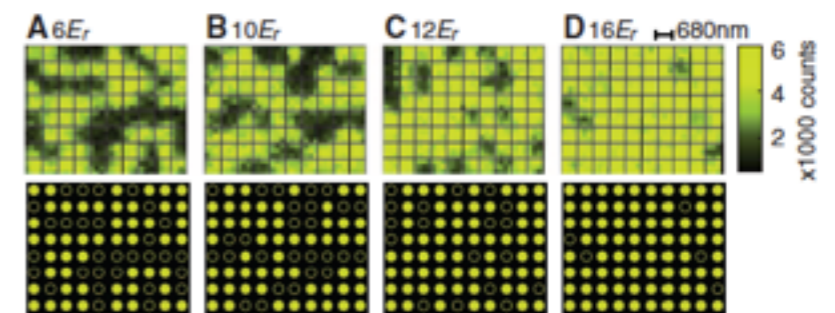
Feshbach resonances:

- ▶ unitary Fermi gas / BEC-BCS crossover
- ▶ ...

MOTT-SUPERFLUID TRANSITION



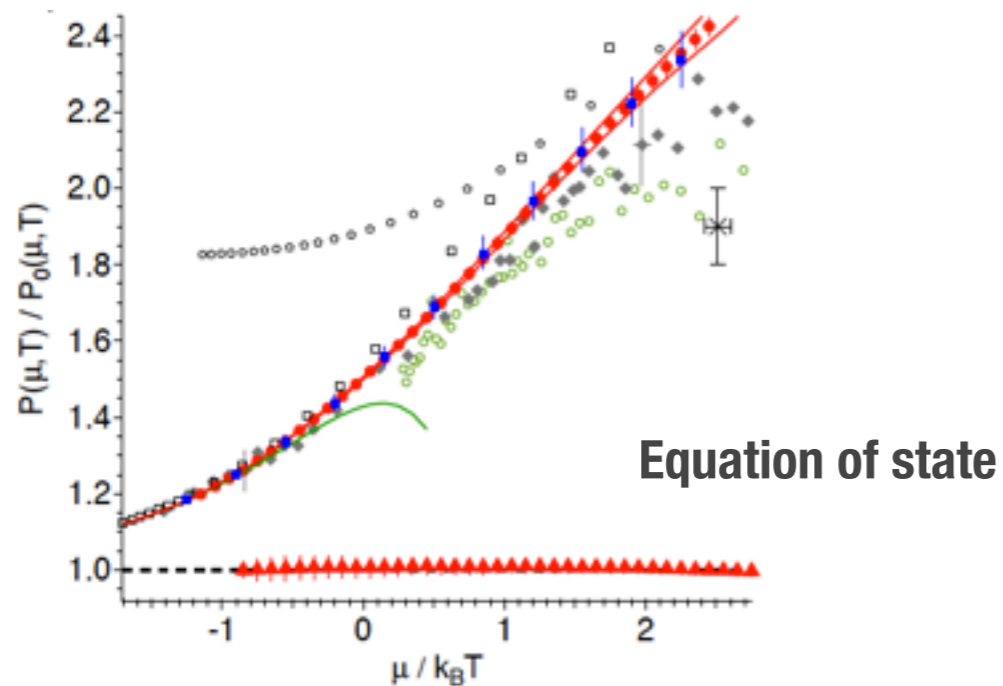
SHERSON ET AL., NATURE 467 (2010)



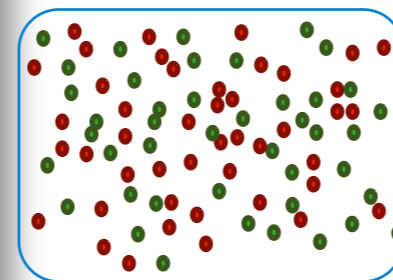
BAKR ET AL. SCIENCE 329 (2010)

UNITARY FERMION GAS

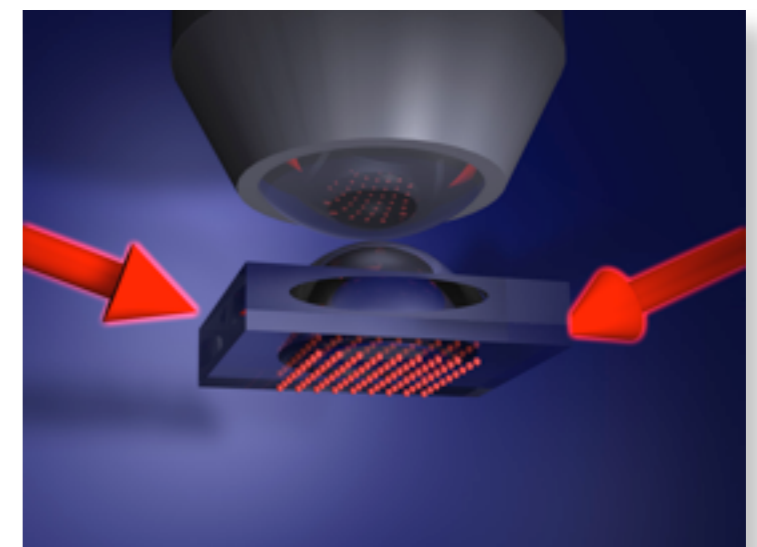
ZWIERLEIN (MIT)



Van Houcke et al. Nature Phys. 8 (2012)



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BAKR ET AL. SCIENCE 329 (2010)

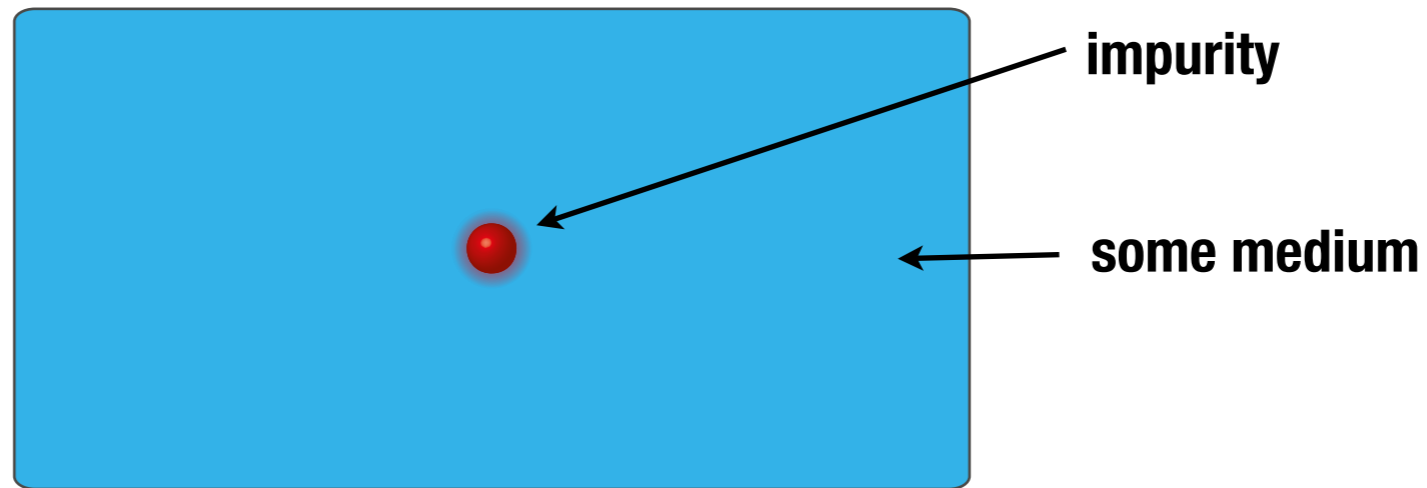
IMAGE: CUAWEB.MIT.EDU - GREINER GROUP

This talk: Impurity physics

Quantum impurities



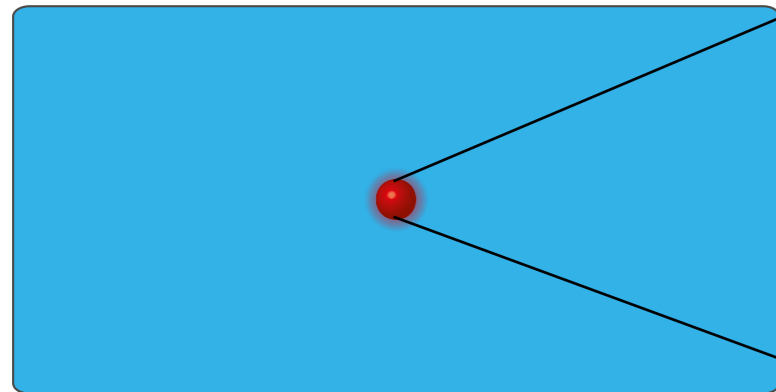
Paradigm of condensed matter physics



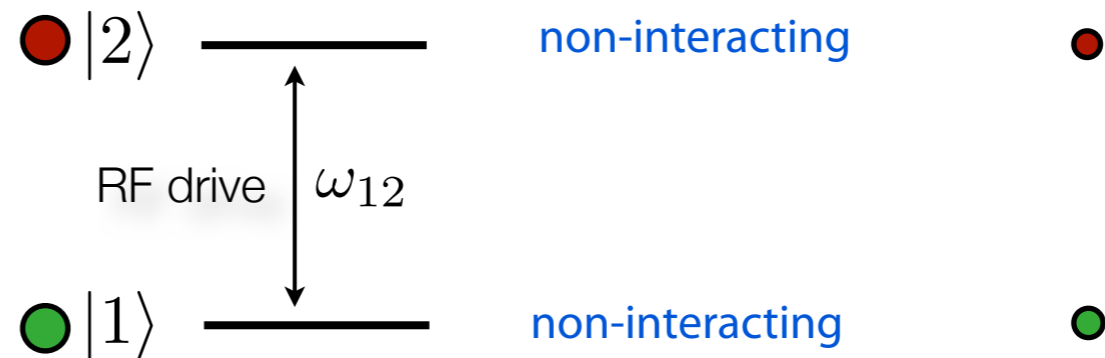
- ▶ appear in many flavors in condensed matter physics
- ▶ relatively simple system from many-body perspective: allow to advance theory in 'controlled way'
- ▶ system on the verge from few- to many-body physics

This talk: Impurity physics

- ▶ impurity physics appears also naturally in: quantum optics, quantum dots, NV centers, atomic clocks, cavity QED...



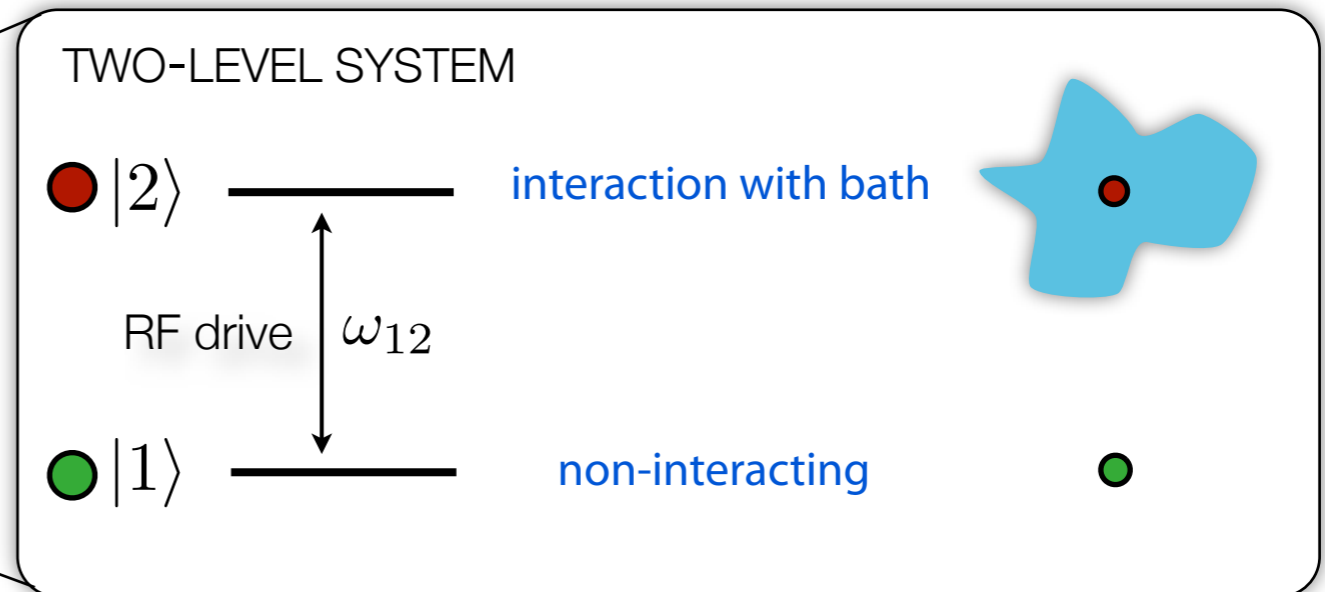
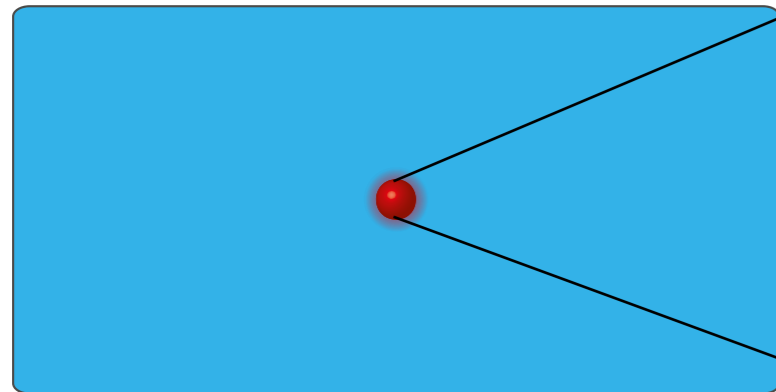
TWO-LEVEL SYSTEM



no interactions: ▶ coherent Rabi oscillations

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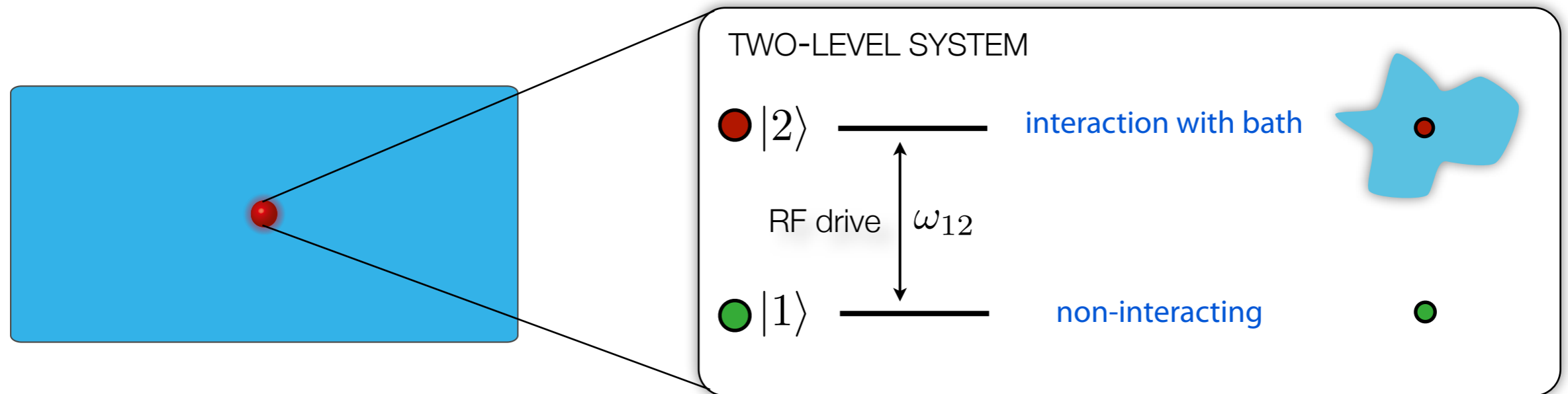


finite interactions: ▶ decoherence of Rabi oscillations



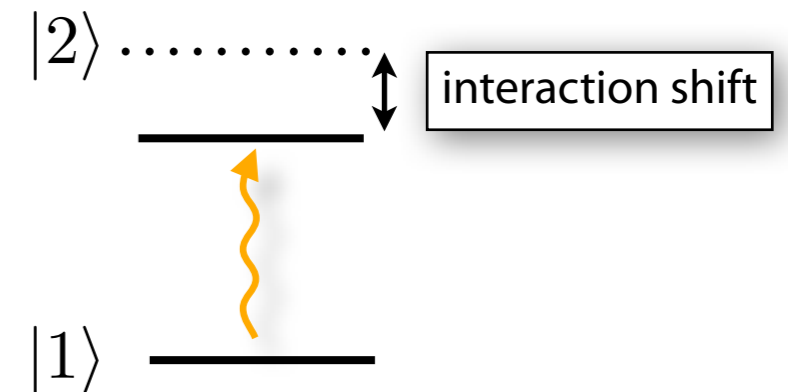
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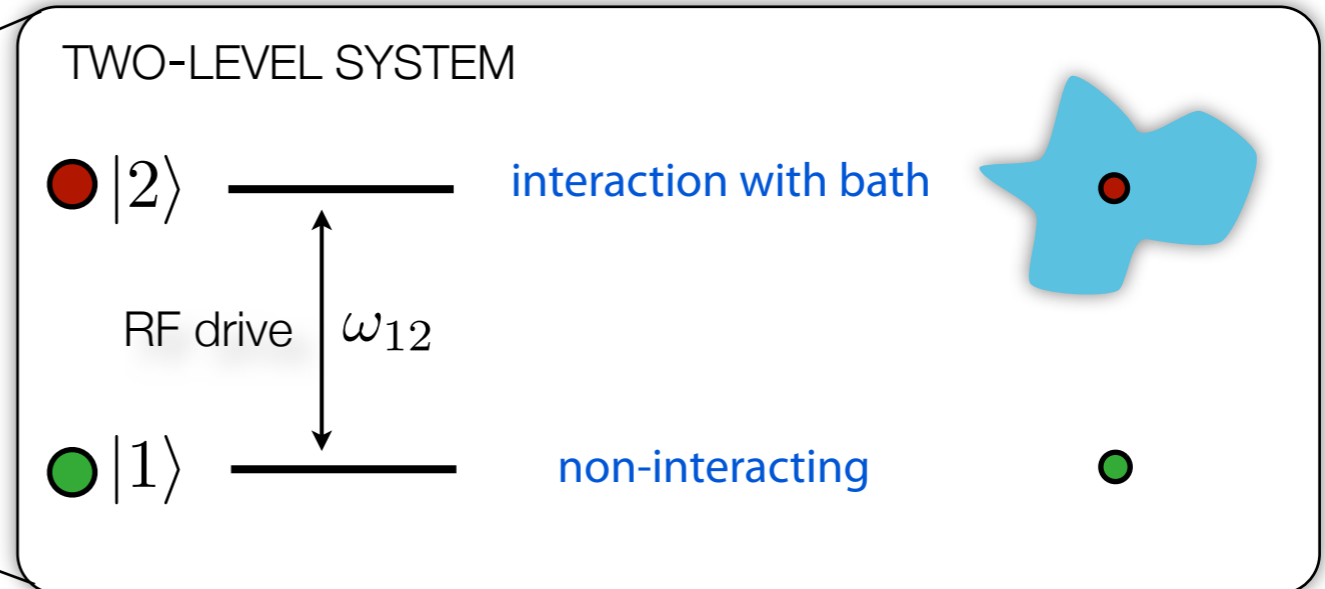
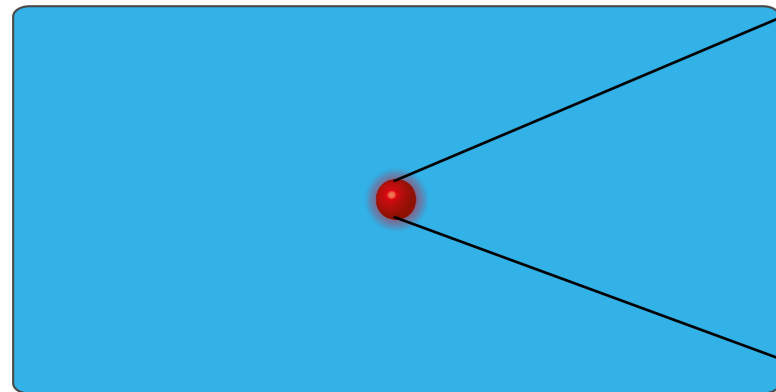
no interactions: ▶ coherent Rabi oscillations

- finite interactions:
- ▶ decoherence of Rabi oscillations
 - ▶ level shifts
 - ▶ line broadening



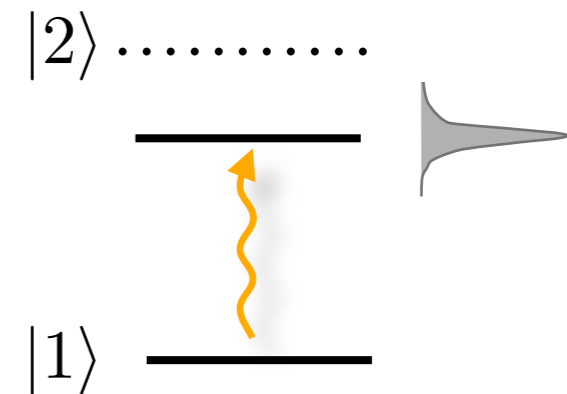
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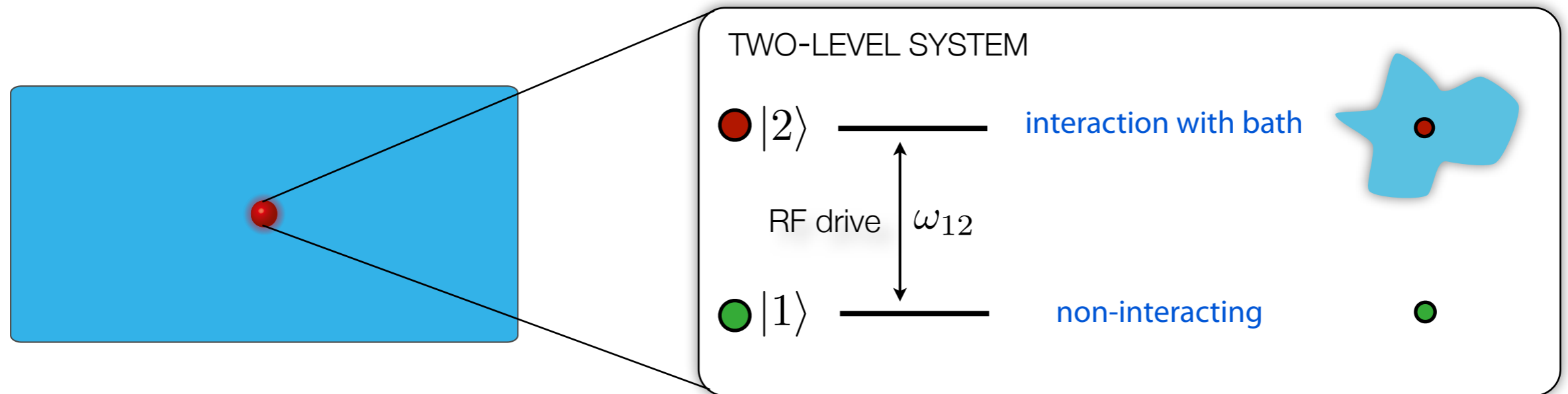
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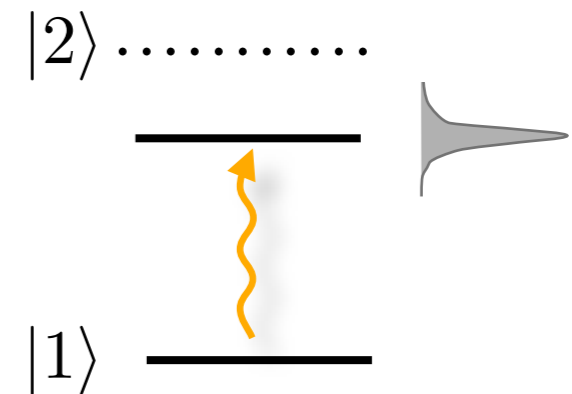
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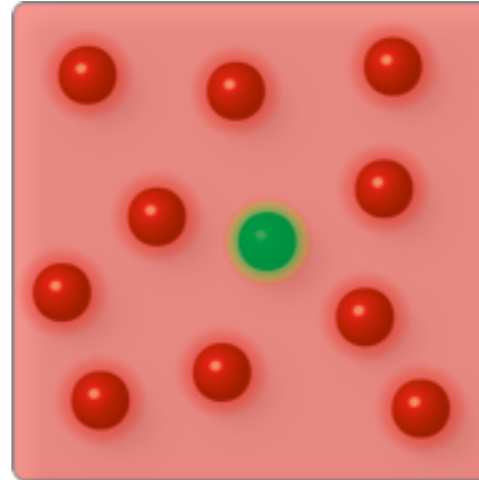
Theoretical challenge: Calculate properties of impurity (strongly) coupled to environment

Impurity experiments with ultracold atoms

Mainly studied experimentally:

THEORY REPULSIVE BRANCH PHYSICS: CUI, ZHAI, PRA 81 (2010), RS, ENSS, PRA 83 (2011), MASSIGNAN, BRUUN, EPJD 65 (2011)
REVIEW: MASSIGNAN, ZACCANTI, BRUUN REP. PROG. PHYS. 77, 034401 (2014)

Fermi polaron
[in continuum]



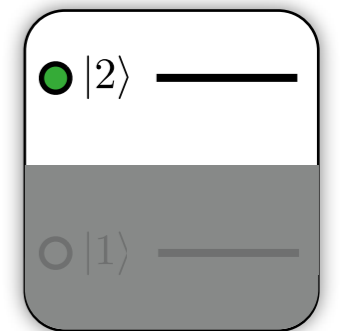
ideal Fermi gas

● ↑-atoms

mobile impurity

● ↓-atom

strongly coupled



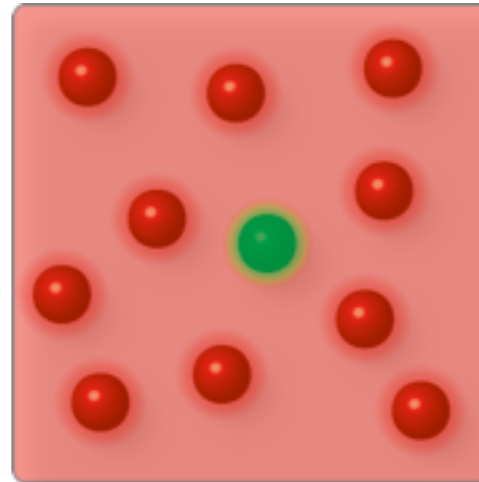
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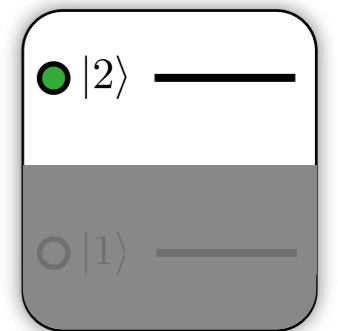
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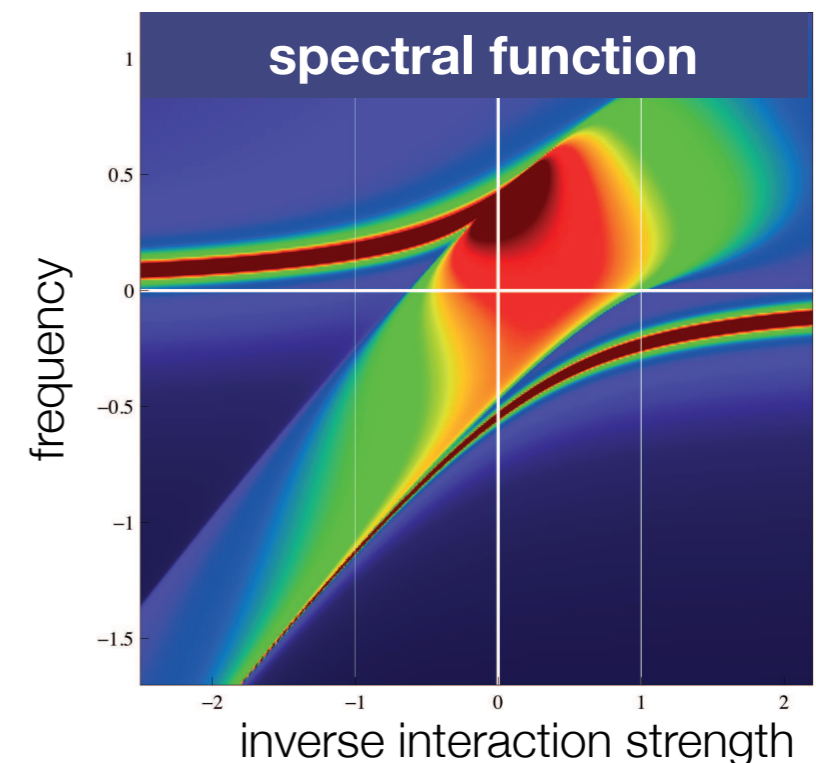
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- ▶ novel functional renormalization group method for non-perturbative **RG flows of spectral functions**

RS, ENSS, PRA 83 (2011)

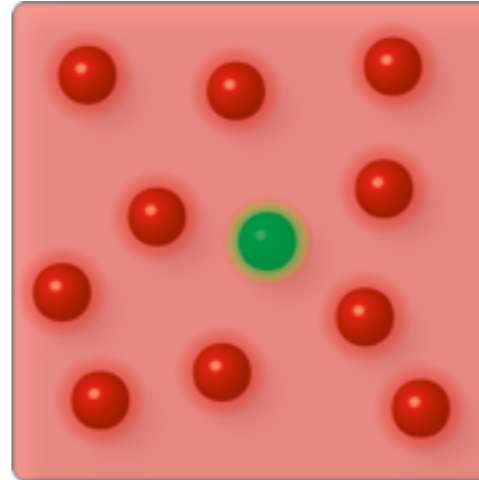


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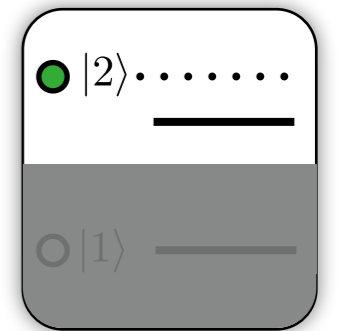
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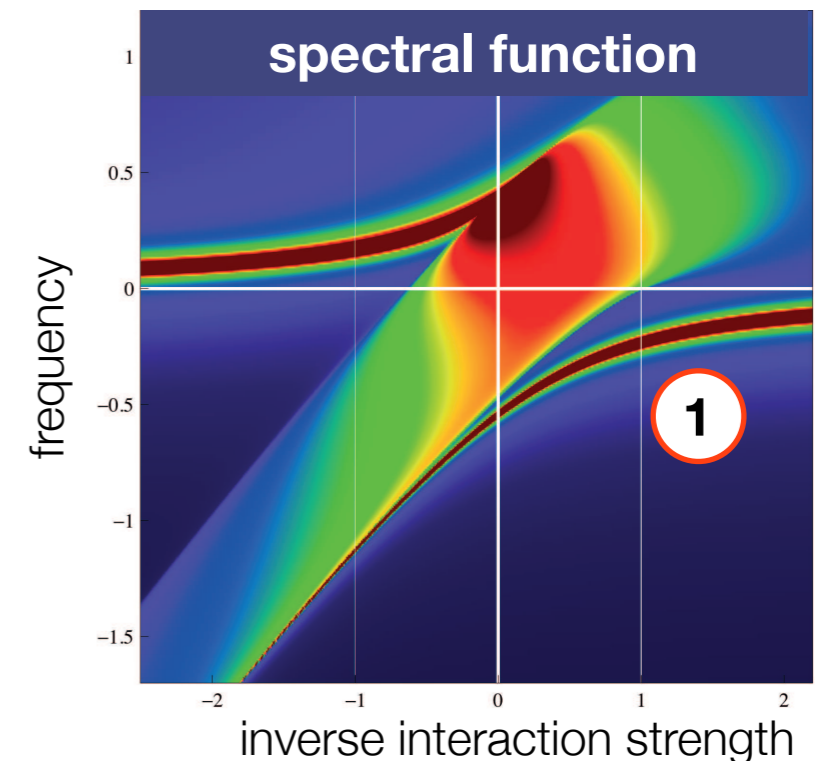
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① regime of weak attractive interactions

▶ energy shift of coherent level

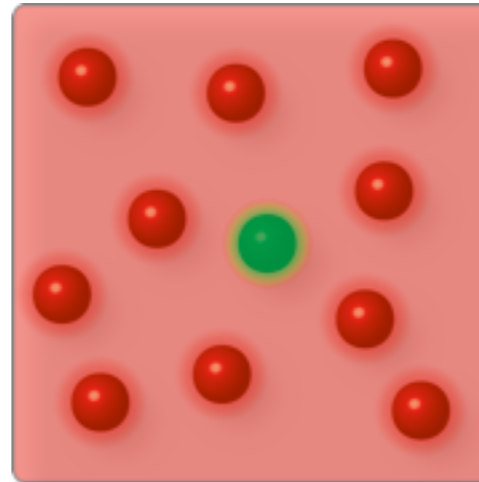


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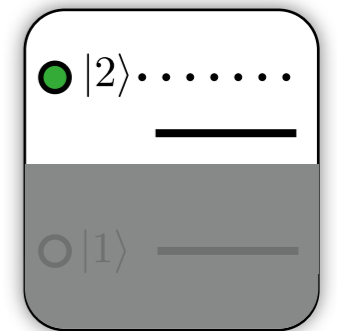
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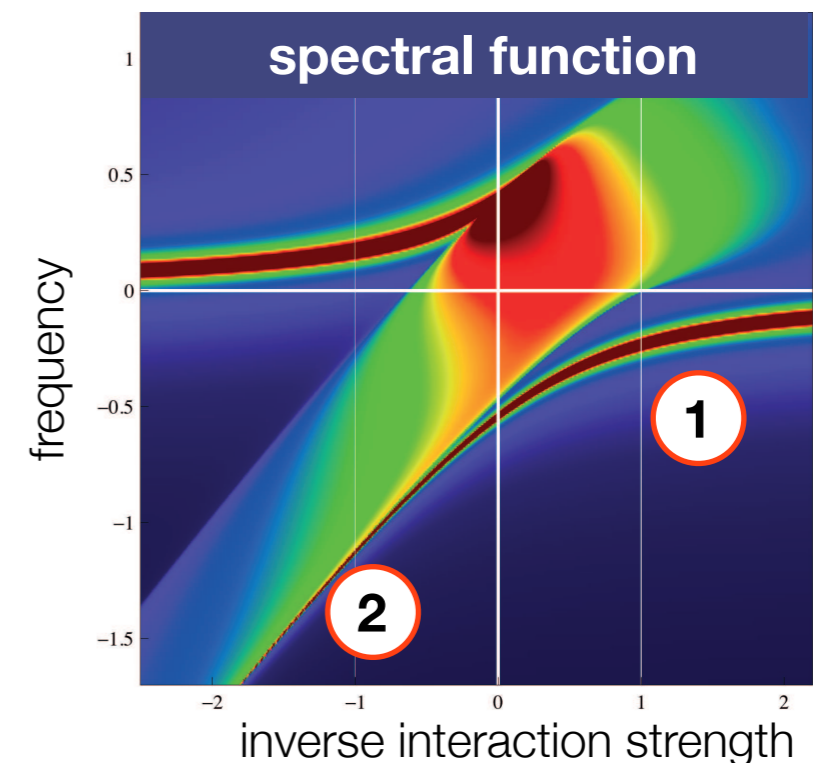


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 - ▶ energy shift of coherent level
- ② regime of strong attractive interactions

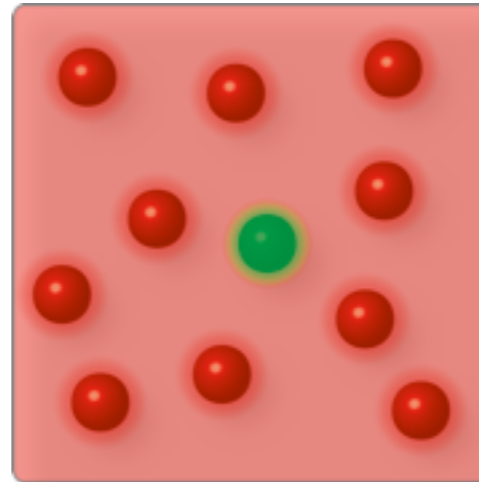


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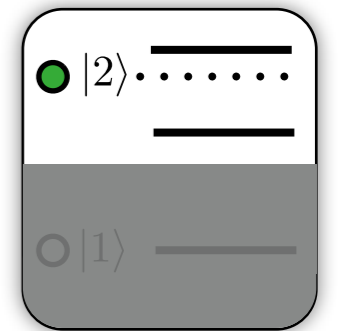
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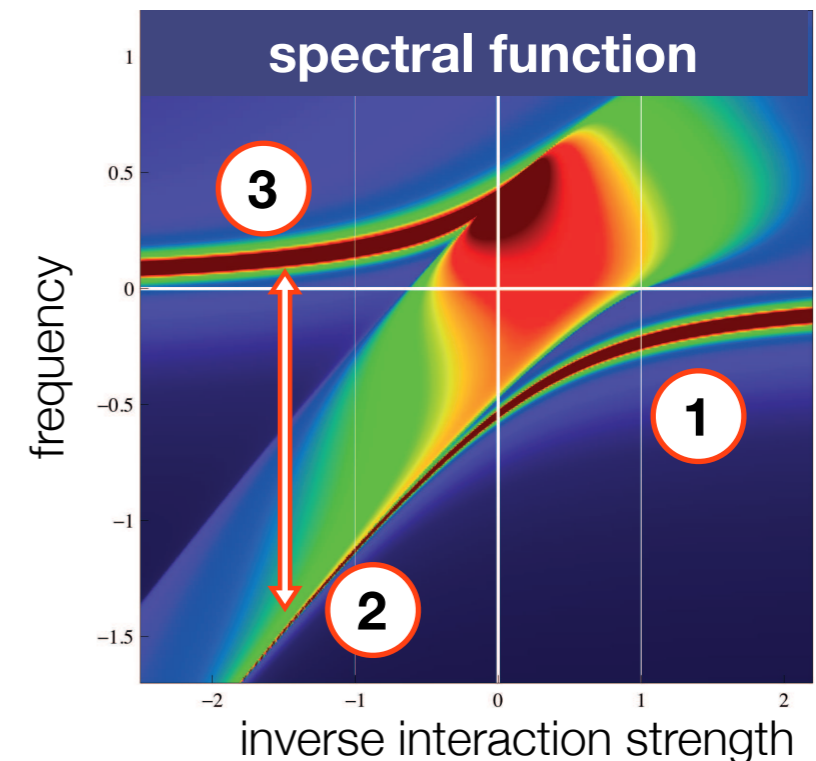
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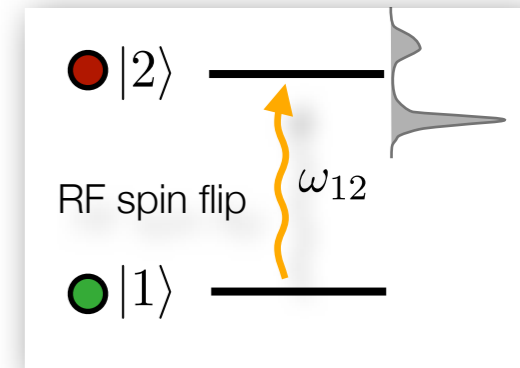
- ① regime of weak attractive interactions
 - ▶ energy shift of coherent level
- ② regime of strong attractive interactions
- ③ *emergent* effective repulsive interactions
 - ▶ single state has split into two branches
 - ▶ repulsive branch: ferromagnetic transition

cf. WORK BY TIN-LUN HO, AND MASSIGNAN, YU, BRUUN, PRL 110 (2013)



Experimental observation

our proposal: inverse radio-frequency [rf] spectroscopy RS, ENSS, PRA 83 (2011)



THREE SPATIAL DIMENSIONS

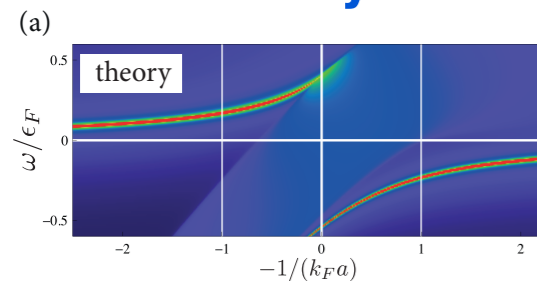
LETTER KOHSTALL ET AL., NATURE 485 (2012)

doi:10.1038/nature11065

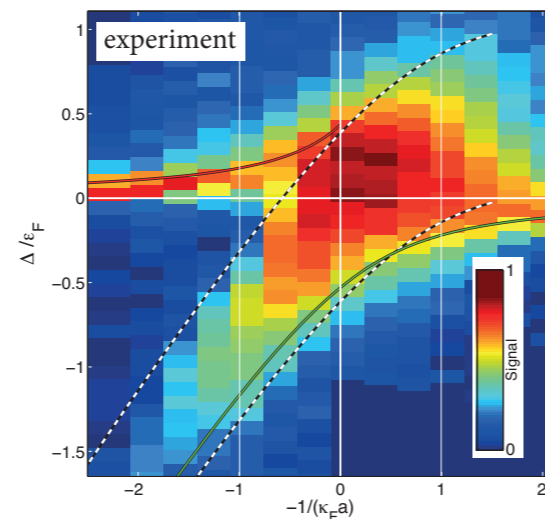
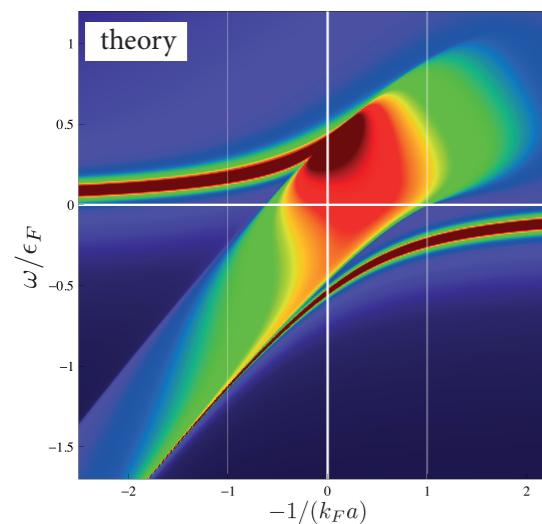
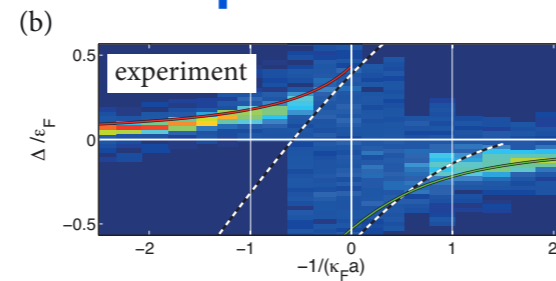
Metastability and coherence of repulsive polarons in a strongly interacting Fermi mixture

C. Kohstall^{1,2}, M. Zaccanti¹, M. Jag^{1,2}, A. Trenkwalder¹, P. Massignan³, G. M. Bruun⁴, F. Schreck¹ & R. Grimm^{1,2}

theory



experiment



RS, ENSS, PRA 83 (2011)

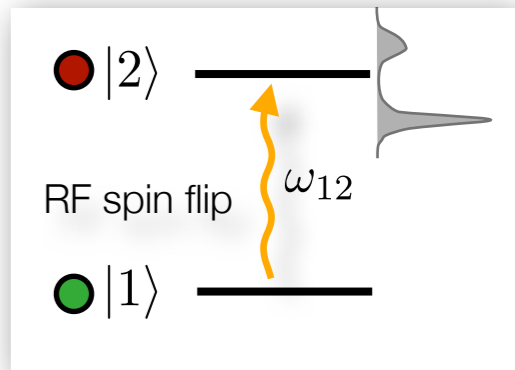
MASSIGNAN, BRUUN, EPJD 65 (2011)

RS, PHD THESIS (2013)

KOHSTALL ET AL., NATURE 485 (2012)

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TWO SPATIAL DIMENSIONS

LETTER

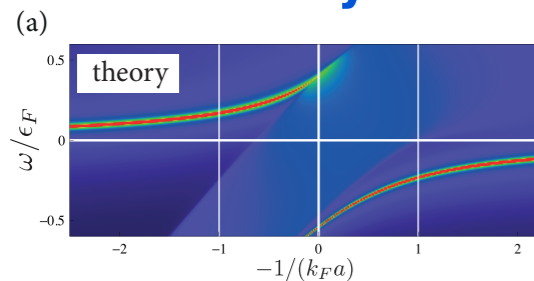
31 MAY 2012 | VOL 485 | NATURE | 619

doi:10.1038/nature11151

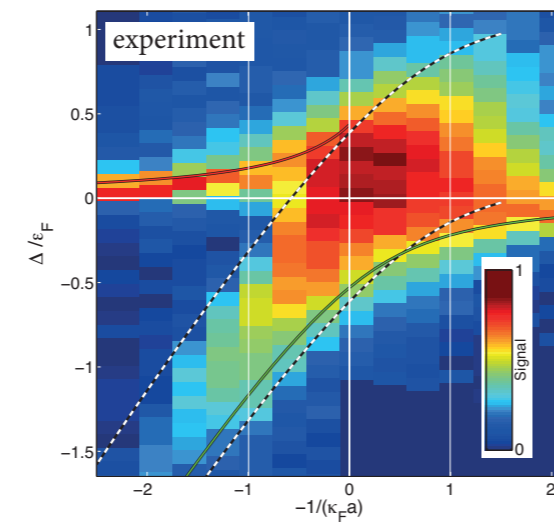
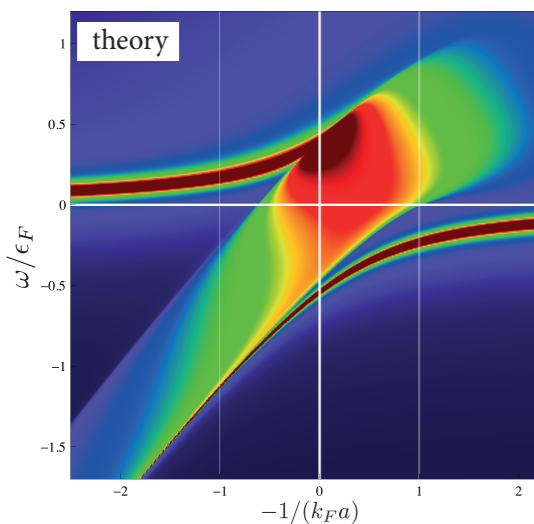
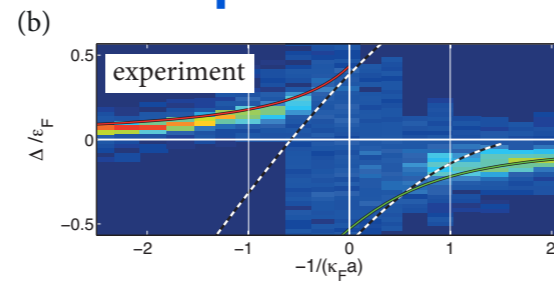
Attractive and repulsive Fermi polarons in two dimensions

Marco Koschorreck^{1*}, Daniel Pertot^{1*}, Enrico Vogt¹, Bernd Fröhlich¹, Michael Feld¹ & Michael Köhl¹

theory



experiment

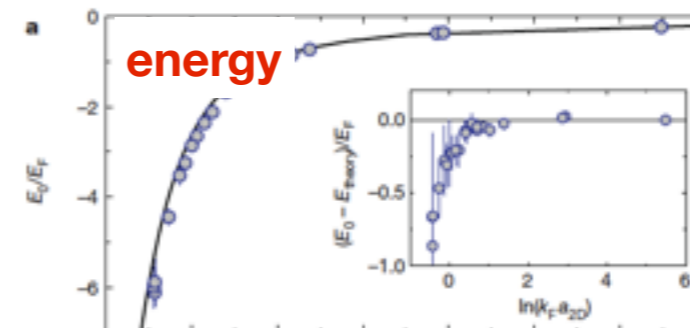


RS, ENSS, PRA 83 (2011)

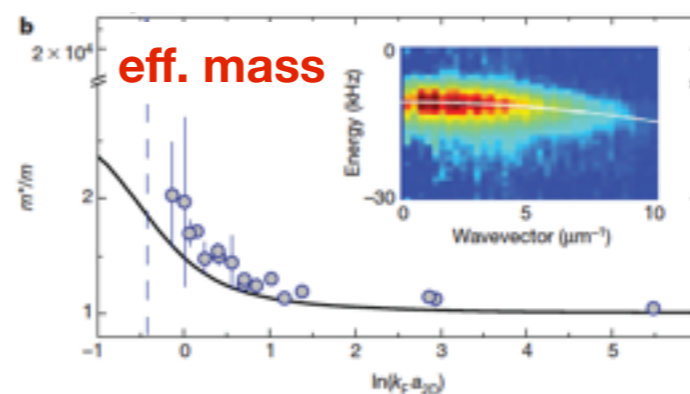
MASSIGNAN, BRUUN, EPJD 65 (2011)

RS, PHD THESIS (2013)

KOHSTALL ET AL., NATURE 485 (2012)



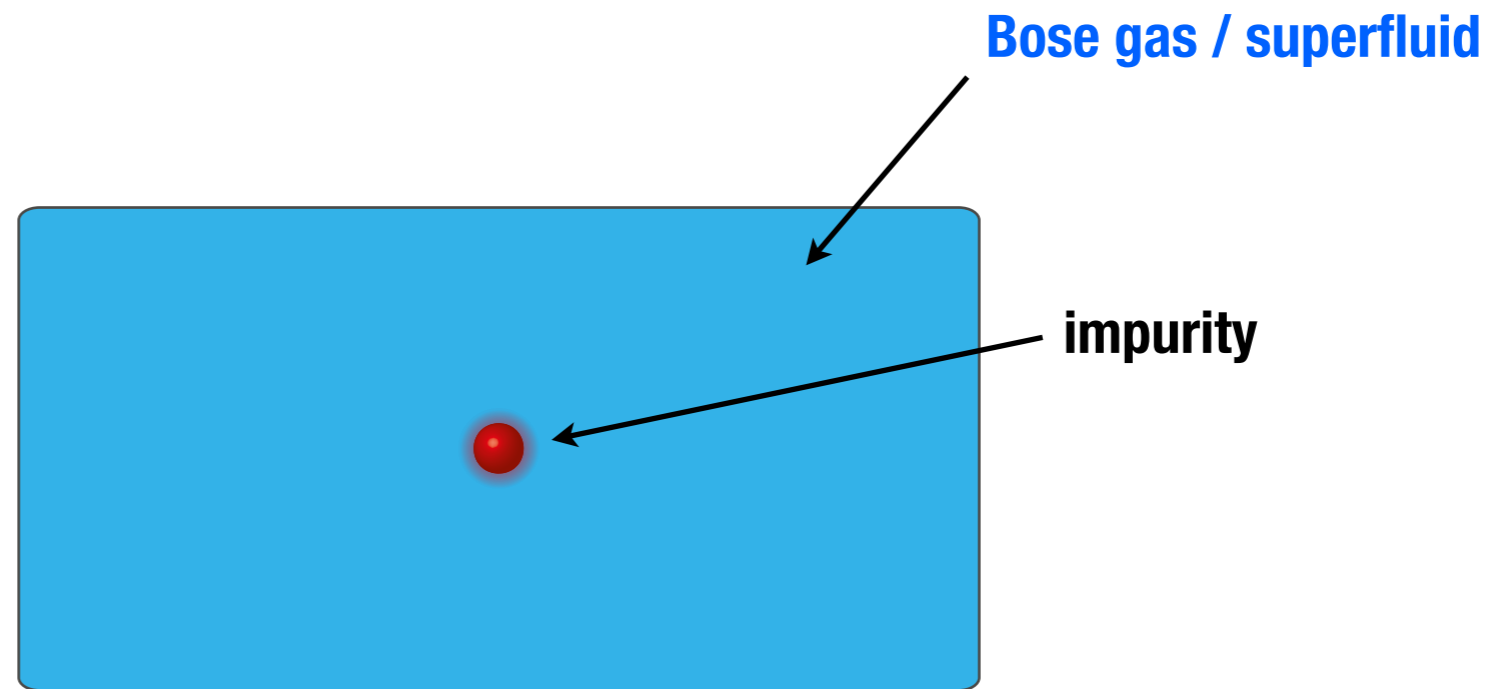
● experiment
KOSCHORRECK ET AL., NATURE 485 (2012)



— theory
RS, ENSS, PIETILA, DEMLER, PRA 85, 021602 (2012)

inverse interaction strength

Our recent work: **The Bose Polaron**

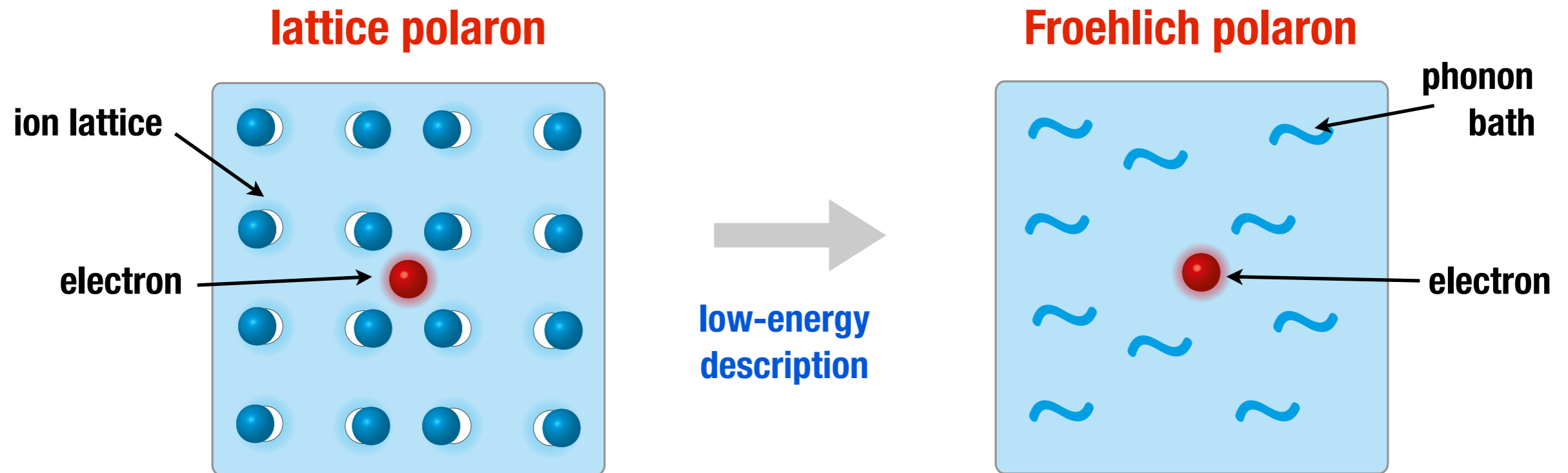


what happens if medium is Bose gas?

⇒ **Bose polaron**

A cond-mat motivation: The Froehlich polaron

a paradigm condensed matter model:



Fröhlich Hamiltonian FRÖHLICH, ADV. PHYS. 3, 325 (1954)

$$\hat{H} = \underbrace{\sum_{\mathbf{p}} \omega_{\mathbf{p}} \hat{b}_{\mathbf{p}}^{\dagger} \hat{b}_{\mathbf{p}}}_{\text{phonons}} + \underbrace{\sum_{\mathbf{p}} \epsilon_{\mathbf{p}} \hat{c}_{\mathbf{p}}^{\dagger} \hat{c}_{\mathbf{p}}}_{\text{impurity}} + \underbrace{\sum_{\mathbf{q}, \mathbf{p}} \alpha_{\mathbf{q}} \hat{c}_{\mathbf{p}+\mathbf{q}}^{\dagger} \hat{c}_{\mathbf{p}} (\hat{b}_{-\mathbf{q}}^{\dagger} + \hat{b}_{\mathbf{q}})}_{\text{impurity-phonon interaction}}$$

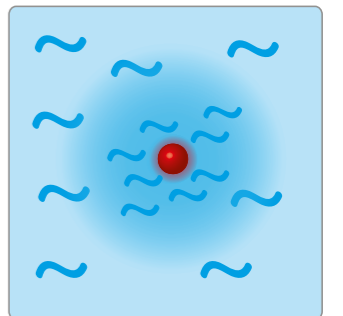
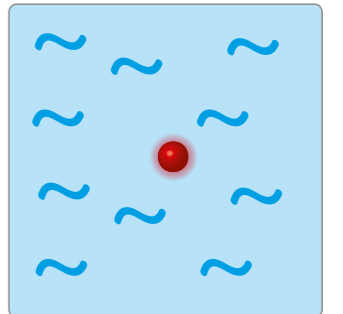
The Fröhlich polaron

Fröhlich Hamiltonian FRÖHLICH, ADV. PHYS. 3, 325 (1954)

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$\alpha > 0$

- ▶ impurity dressed by *phonon cloud* becomes the **'Fröhlich polaron'**
- ▶ enhanced effective mass, renormalized energy SEE E.G. MILLER ET AL. PHYS. REV. 127 ('62)

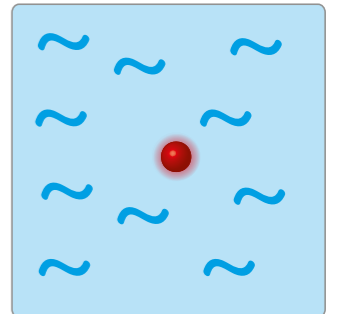


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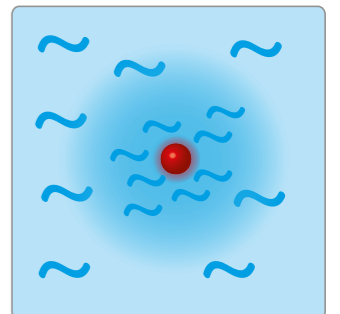
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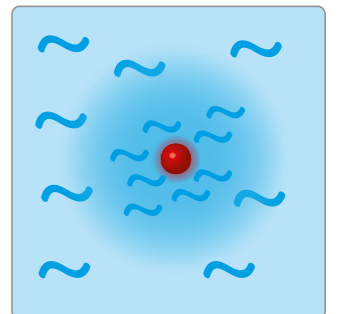
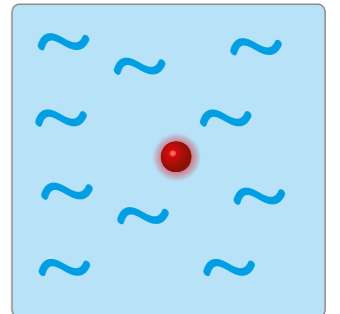
perturbation theory: $m^* = \frac{m}{1 - \alpha/6}$ $\xrightarrow[\text{large}]{\alpha > 0}$ ∞ \rightarrow **self-localization?**

The Fröhlich polaron

Fröhlich Hamiltonian FRÖHLICH, ADV. PHYS. 3, 325 (1954)

$$\hat{H} = \sum_{\mathbf{p}} \omega_p \hat{b}_{\mathbf{p}}^\dagger \hat{b}_{\mathbf{p}} + \sum_{\mathbf{p}} \epsilon_{\mathbf{p}} \hat{c}_{\mathbf{p}}^\dagger \hat{c}_{\mathbf{p}} + \sum_{\mathbf{q}, \mathbf{p}} \alpha_q \hat{c}_{\mathbf{p}+\mathbf{q}}^\dagger \hat{c}_{\mathbf{p}} (\hat{b}_{-\mathbf{q}}^\dagger + \hat{b}_{\mathbf{q}})$$

$\alpha > 0$

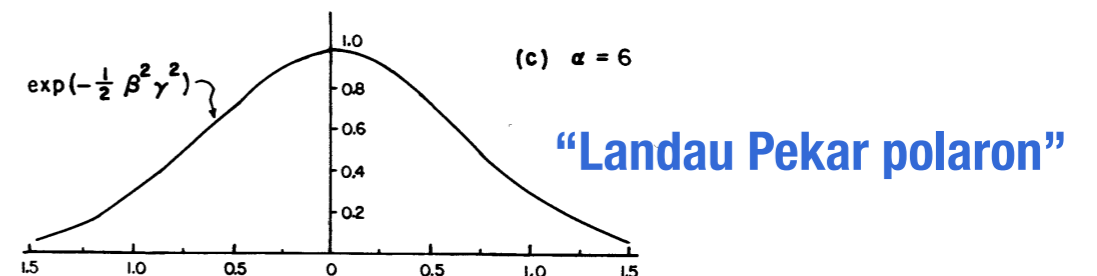


- ▶ impurity dressed by *phonon cloud* becomes the **‘Fröhlich polaron’**
- ▶ enhanced effective mass, renormalized energy SEE E.G. MILLER ET AL. PHYS. REV. 127 ('62)

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- ▶ **strong interactions:** variational wave function

LANDAU, PEKAR, JETP 18 (1948); FEYNMAN, COHEN, PHYS. REV. 102 (1956)



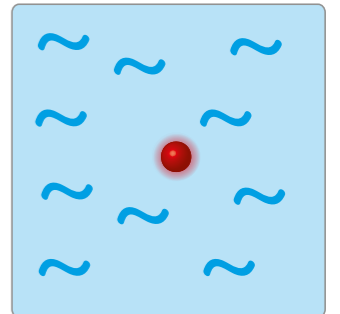
- ▶ describes localized particle
- ▶ yields energy smaller than pert. theory at strong coupling, **further evidence of self-localization**

The Fröhlich polaron

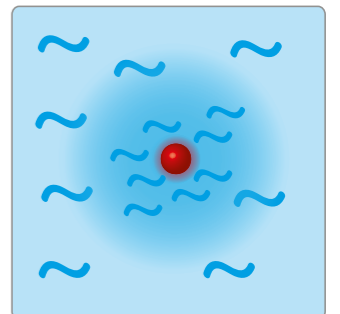
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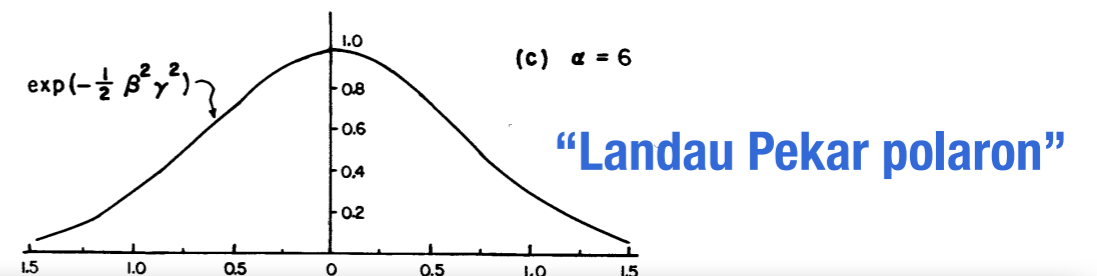


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- ▶ **strong interactions:** variational wave function
LANDAU, PEKAR, JETP 18 (1948); FEYNMAN, COHEN, PHYS. REV. 102 (1956)



superfluid + strong interactions:

Realizable with ultracold atoms!

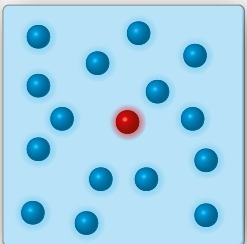
Let's do quantum simulation!

The Bose polaron with ultracold atoms

impurity in Bose gas: Bose polaron

e.g. take strongly imbalanced mixture of ultracold atoms CF. GENERAL CASE: B. LIU, J. HU INT. J. MOD. PHYS, 26 (2012)

$$S = \int d^4x \underbrace{\varphi^*(x)}_{\text{bosons}} \left(\partial_\tau - \frac{1}{2m_B} \nabla^2 - \mu_B \right) \underbrace{\varphi(x)}_{\text{bosons}} + \frac{g_B}{2} [\underbrace{\varphi(x)^* \varphi(x)}_{\text{bosons}}]^2$$
$$+ \underbrace{\psi^*(x)}_{\text{impurity}} \left(\partial_\tau - \frac{1}{2m_I} \nabla^2 - \mu_I \right) \underbrace{\psi(x)}_{\text{impurity}} + g_{IB} \underbrace{\psi(x)^* \psi(x)}_{\text{impurity}} \underbrace{\varphi(x)^* \varphi(x)}_{\text{bosons}}$$



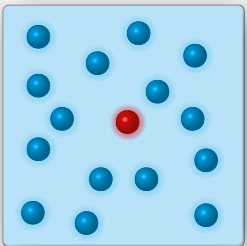
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bosons
impurity



weakly interacting BEC: Bogoliubov approximation for BEC

$$\varphi(\mathbf{x}, t) = \underbrace{\sqrt{n_B}}_{\text{mean-field}} + \underbrace{\phi(\mathbf{x}, t)}_{\text{fluctuations}}$$

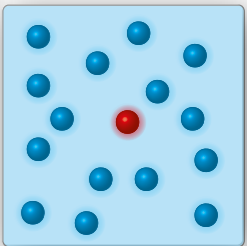
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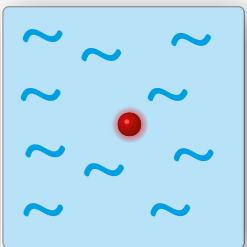


weakly interacting BEC: Bogoliubov approximation for BEC

$$\varphi(\mathbf{x}, t) = \underbrace{\sqrt{n_B}}_{\text{mean-field}} + \underbrace{\phi(\mathbf{x}, t)}_{\text{fluctuations}}$$

at weak coupling: Fröhlich Hamiltonian

$$S_{\text{eff}} = S_{\text{kin}}^{\text{Bos}} + S_{\text{kin}}^{\text{Imp}} + \underbrace{g_{IB} \int_x n_B(\mathbf{r}) |\psi(\mathbf{r})|^2}_{\text{MF energy shift}} + g_{IB} \int \left\{ \underbrace{\sqrt{n_B} \psi_{\mathbf{k}+\mathbf{q}}^* \psi_{\mathbf{k}} (\phi_{\mathbf{q}} + \phi_{-\mathbf{q}}^*)}_{\text{"Fröhlich term"}} \right\}$$



ELECTRON IN BEC: PFAU GROUP [STUTTGART] - BALEWSKI ET AL., NATURE 502 (2013)

Bose polaron at strong coupling

$$S_{\text{eff}} = S_{\text{kin}}^{\text{Bos}} + S_{\text{kin}}^{\text{Imp}} + g_{IB} \int \left\{ \sqrt{n_B} \psi_x^* \psi_x (\phi_x + \phi_x^*) + n_B \psi_x^* \psi_x \right. \quad \left. g_{IB} = \frac{2\pi\hbar^2}{m_r} a_{IB} \right.$$

"Fröhlich terms"

strong effective phonon-impurity interaction wanted

HEISELBERG ET AL., PRL 85 (2000)

CUCCHIENTTI, TIMMERMANS, PRL 96 (2006)

KALAS, BLUME, PRA 73 (2006)

WANG, PRL 96 (2006)

ENSS, ZWARGER, EPJB 68 (2009)

TEMPERE, OBERTHALER ET AL., PRB 80 (2009)

CASTEELS ET AL., PRA 83,84,86 (2011)

CASTEELS, CAUTEREN, TEMPERE, DEVREESE, LASER PHYS. 21 (2011)

CASTEELS, TEMPERE, DEVREESE, PRA 84 (2011)

CASTEELS, TEMPERE, DEVREESE, PRA 86 (2012)

DASENBROOK, KOMNIK, PRB 87 (2013)

BLINOVA, BOSHIER, TIMMERMANS, PRA 88 (2013)

SHASHI, GRUSDT, ABANIN, DEMLER, PRA 89, 053617 (2014)

GRUSDT, SHCHADILOVA, RUBTSOV, DEMLER, 1410.2203 (2014)

...

Bose polaron at strong coupling

$$S_{\text{eff}} = S_{\text{kin}}^{\text{Bos}} + S_{\text{kin}}^{\text{Imp}} + g_{IB} \int \left\{ \sqrt{n_B} \psi_x^* \psi_x (\phi_x + \phi_x^*) + n_B \psi_x^* \psi_x \right. \quad \left. g_{IB} = \frac{2\pi\hbar^2}{m_r} a_{IB} \right.$$

"Fröhlich terms"

strong *effective* phonon-impurity interaction **comes at a prize**

Bose polaron at strong coupling

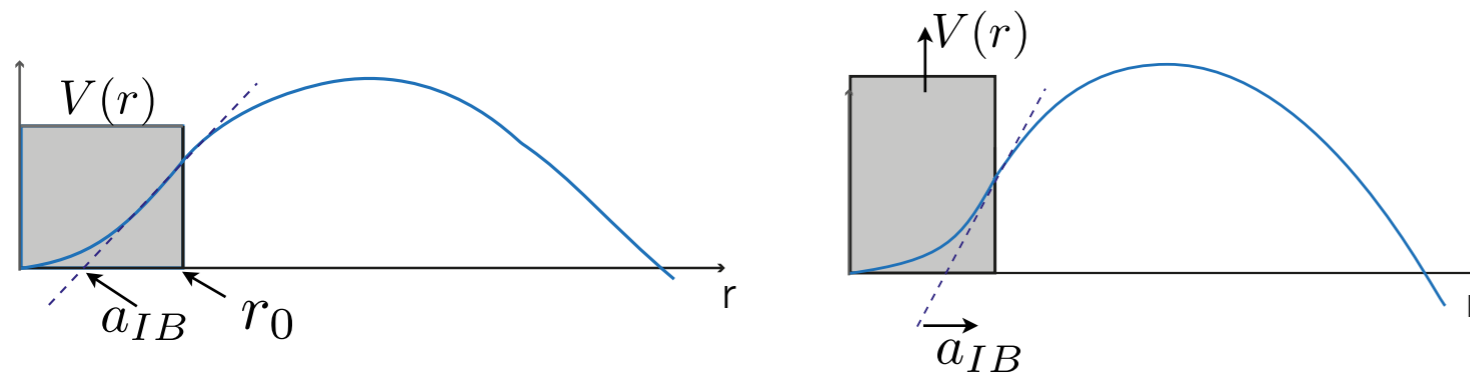
$$S_{\text{eff}} = S_{\text{kin}}^{\text{Bos}} + S_{\text{kin}}^{\text{Imp}} + g_{IB} \int \left\{ \sqrt{n_B} \psi_x^* \psi_x (\phi_x + \phi_x^*) + n_B \psi_x^* \psi_x \right.$$

"Fröhlich terms"

$$g_{IB} = \frac{2\pi\hbar^2}{m_r} a_{IB}$$

How to achieve large, positive scattering length a_{IB} ?

1. way: large microscopic repulsion



not sufficient:

for infinite repulsion

$$a_{\text{max}} = r_0$$

Bose polaron at strong coupling

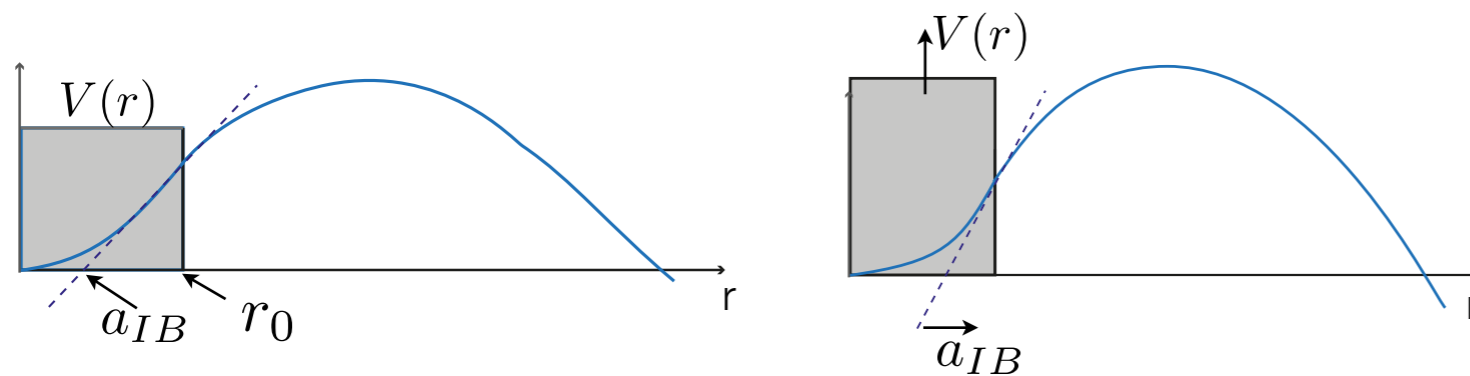
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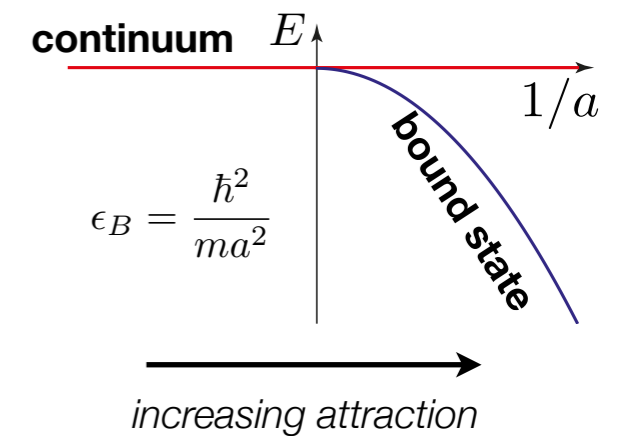
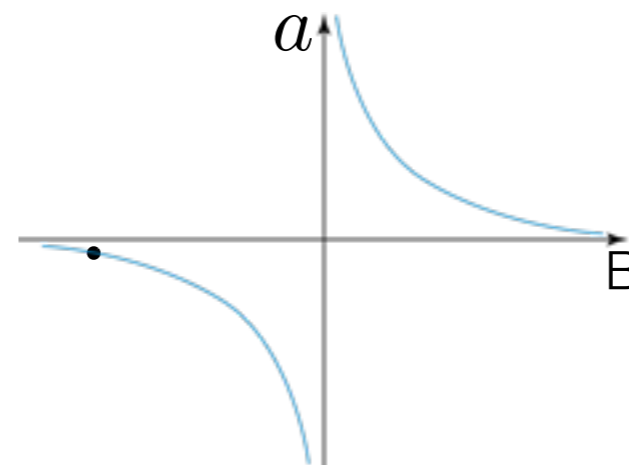
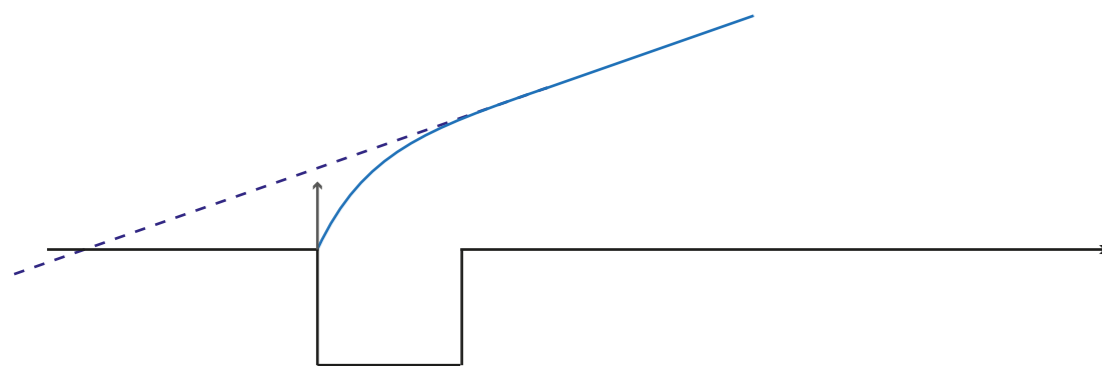


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2. way: use resonance [here shape resonance, cold atoms: typ. Feshbach resonance]



Bose polaron at strong coupling

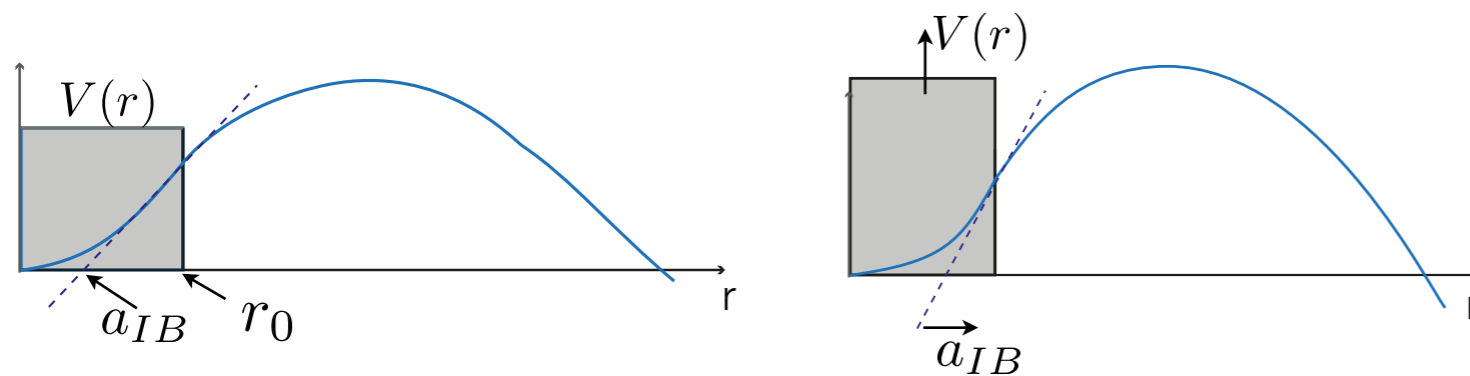
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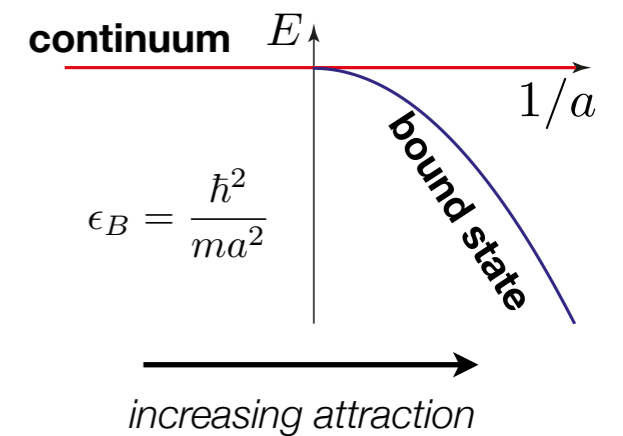
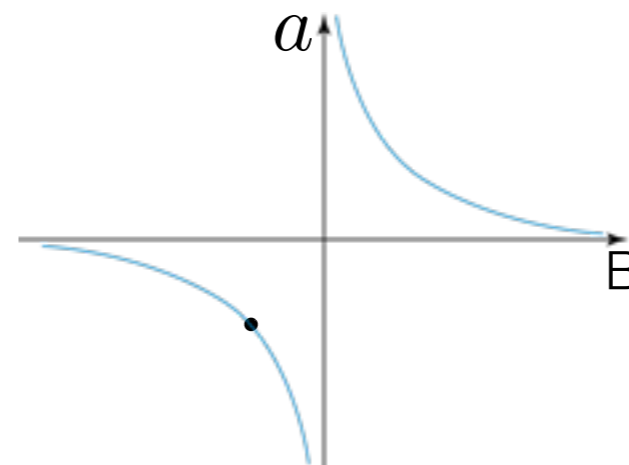
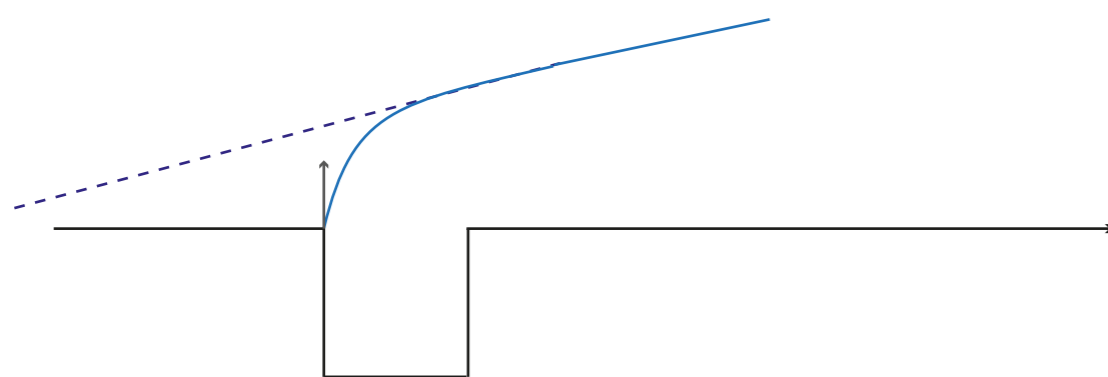


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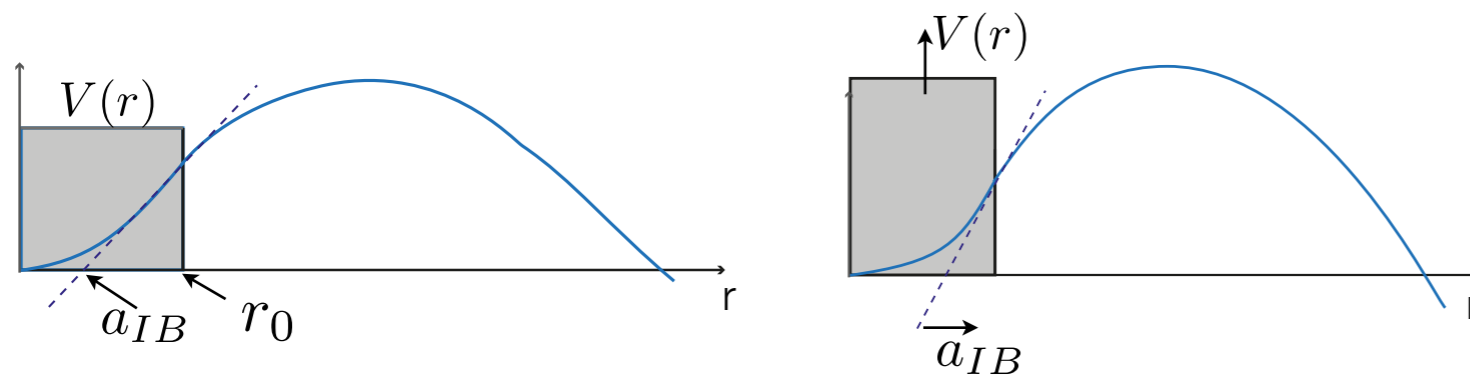
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"Fröhlich terms"

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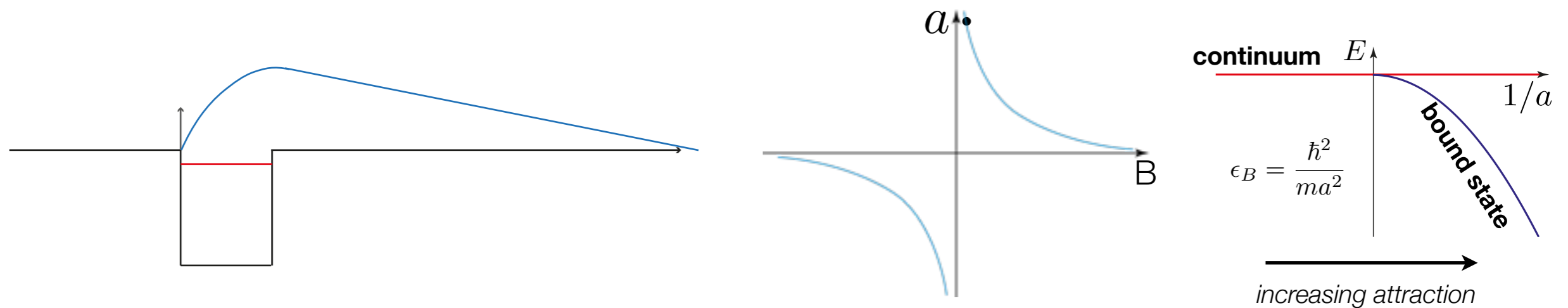


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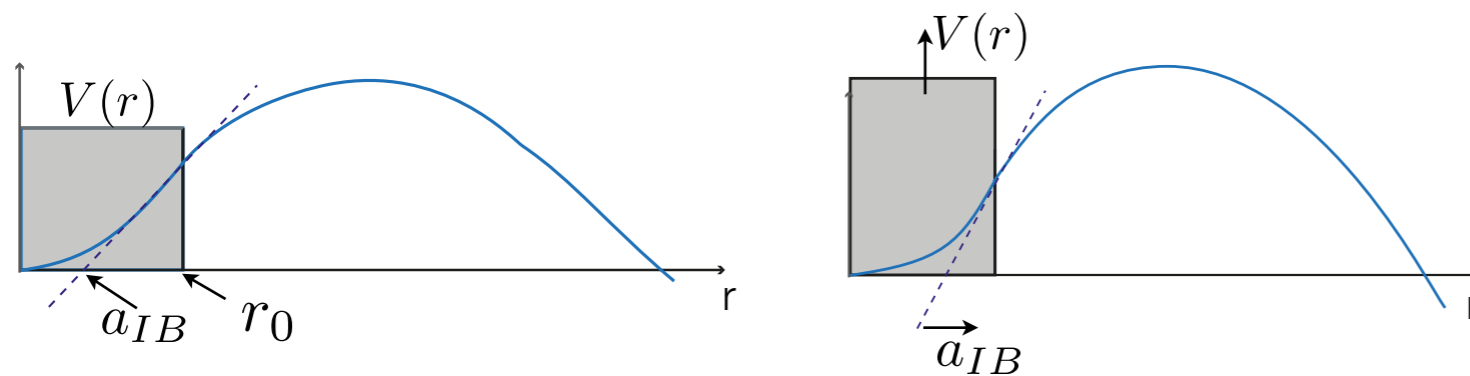
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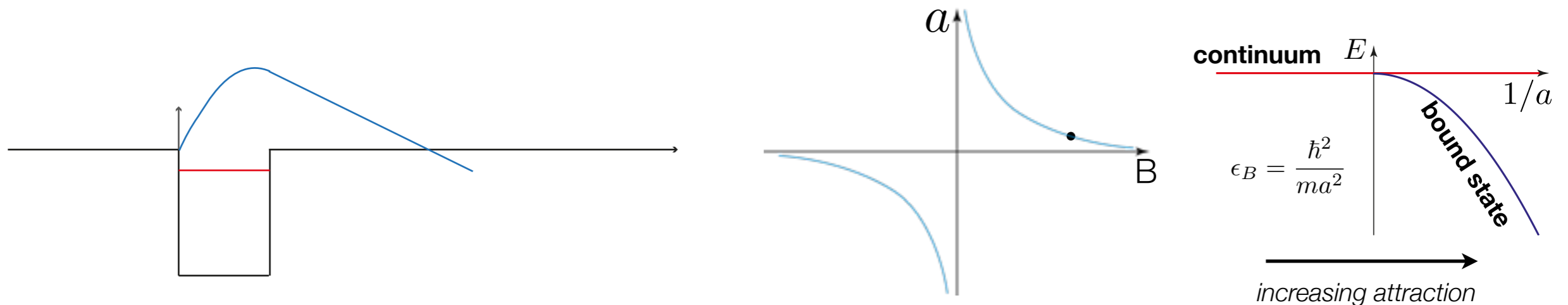


not sufficient:

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2. way: use resonance [here shape resonance, cold atoms: typ. Feshbach resonance]



arbitrarily large a_{IB} possible

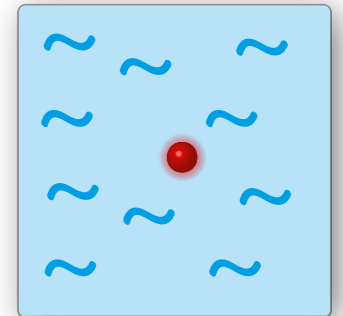
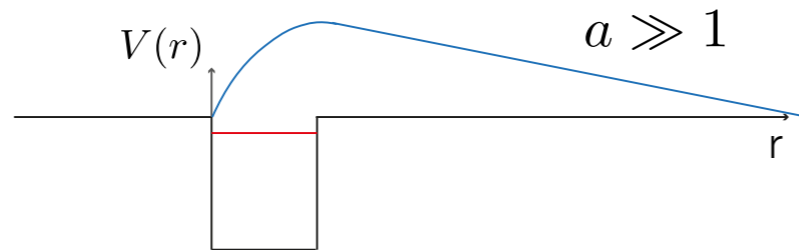
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"Fröhlich terms"

strong effective phonon-impurity interaction comes at a prize

1. microscopic attraction needed



$$g_{IB} \neq \frac{2\pi\hbar^2}{m_r} a_{IB} \quad \text{'mean-field replacement' invalid!}$$

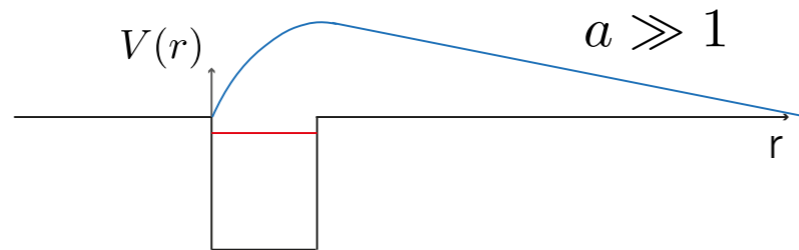
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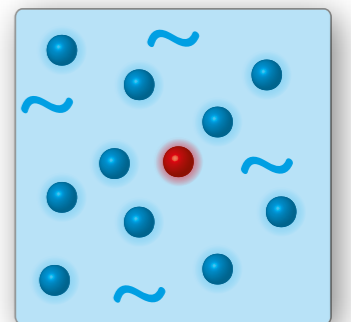
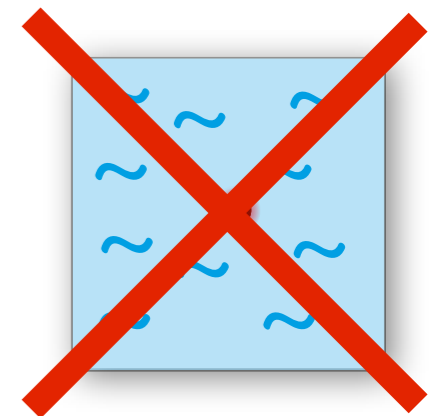
"Fröhlich terms"

strong effective phonon-impurity interaction comes at a prize

1. microscopic attraction needed



$$g_{IB} \neq \frac{2\pi\hbar^2}{m_r} a_{IB} \quad \text{'mean-field replacement' invalid!}$$



2. pairing fluctuations become relevant

→ MF approach & Fröhlich Hamiltonian becomes invalid

in RG language: *Fröhlich*: weak coupling RG fixed point

cold atoms at Feshbach resonance: strong coupling RG fixed point

SIMILAR RG ANALYSIS FOR BOSE-FERMI MIXTURE: B. LIU, J. HU INT. J. MOD. PHYS, 26 (2012)

Our work: Bose polaron from a truly attractive model

simple quantum field-theory approach

$$S = \int d^4x \underbrace{\varphi^*(x)}_{\text{bosons}} \left(\partial_\tau - \frac{1}{2m_B} \nabla^2 - \mu_B \right) \varphi(x) + \underbrace{\psi^*(x)}_{\text{impurity}} \left(\partial_\tau - \frac{1}{2m_I} \nabla^2 - \mu_I \right) \psi(x) + \frac{g_B}{2} \underbrace{[\varphi^*(x)\varphi(x)]^2}_{\text{boson-boson interaction}} + \tilde{g}_{IB} \underbrace{\psi^*(x)\psi(x)}_{\text{impurity}} \underbrace{\varphi^*(x)\varphi(x)}_{\text{bosons}}$$

↪ **attractive interaction**

assume homogeneous, weakly interacting BEC

RATH, RS, PRA 88 (2013)

$$\varphi(\mathbf{x}, t) = \underbrace{\rho_0^{1/2}}_{\text{mean-field}} + \underbrace{\phi(\mathbf{x}, t)}_{\text{fluctuations}}$$

Bogoliubov approximation for bosons: keep all terms up to quadratic in ϕ, ϕ^*

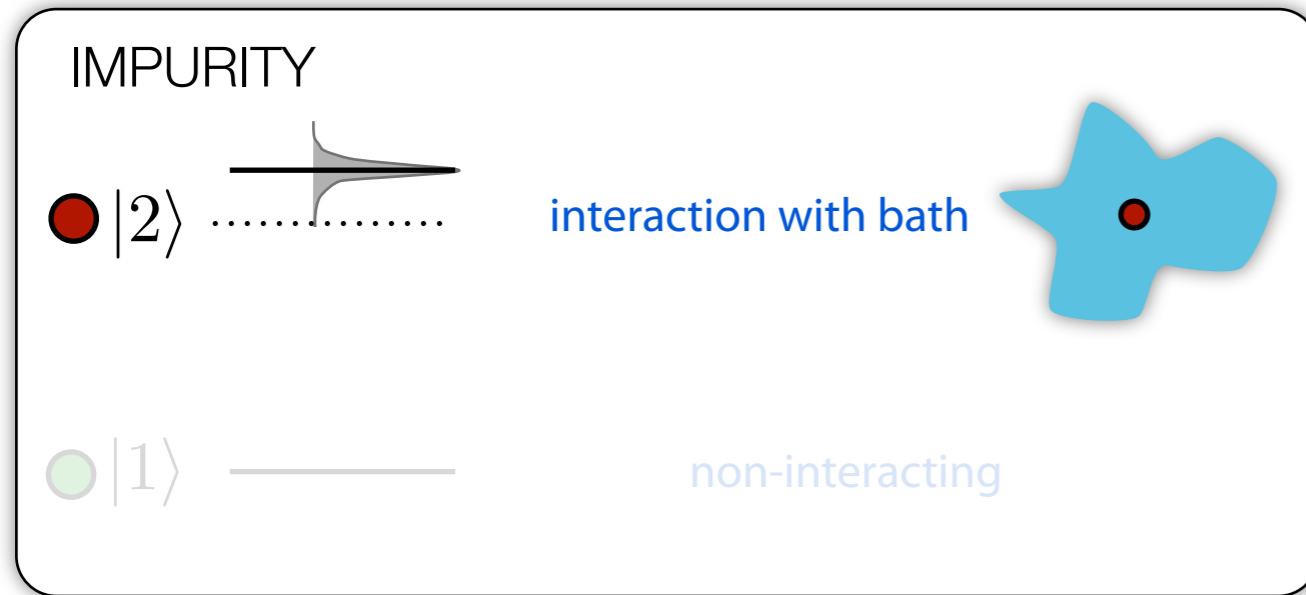
$$S_{\text{eff}} = \int_{\omega, \mathbf{p}} \left\{ \frac{1}{2} \begin{pmatrix} \phi_p^* \\ \phi_{-p} \end{pmatrix} \begin{pmatrix} - [G_\phi^{(0)}(-p)]^{-1} & g_{\phi\phi}\rho_0 \\ g_{\phi\phi}\rho_0 & - [G_\phi^{(0)}(p)]^{-1} \end{pmatrix} \begin{pmatrix} \phi_p \\ \phi_{-p}^* \end{pmatrix} + \psi_p^* \left(-i\omega + \frac{\mathbf{p}^2}{2m_\psi} - \mu_\psi \right) \psi_p \right\} + \tilde{g}_{\phi\psi} \int_x \left[\underbrace{\psi_x^* \psi_x \phi_x^* \phi_x}_{\text{pairing fluctuations}} + \underbrace{\sqrt{\rho_0} \psi_x^* \psi_x (\phi_x + \phi_x^*) + \rho_0 \psi_x^* \psi_x}_{\text{"Fröhlich terms"}} \right]$$

Unlike previous approaches, we keep **pairing fluctuations**

"Fröhlich terms"

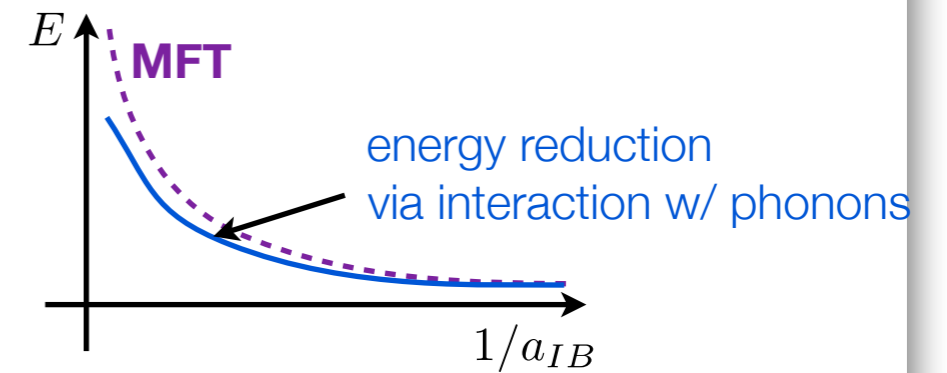
The Question

What is the spectrum of the model?



condensed matter/ pure Froehlich model

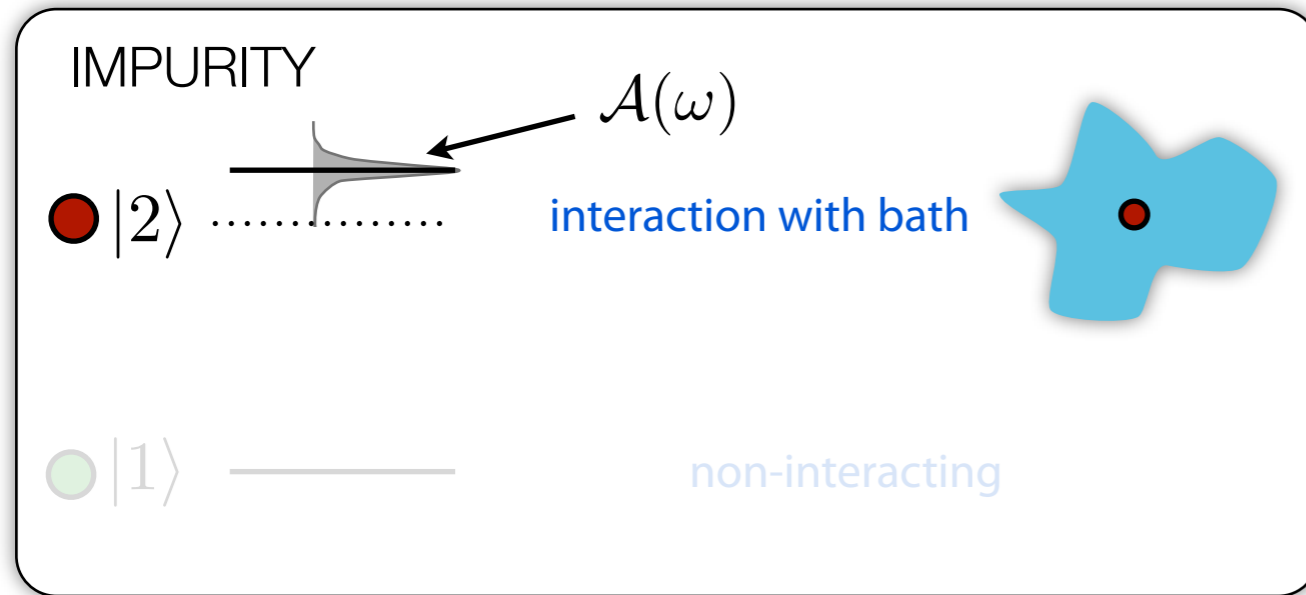
predicts sharp quasi-particle excitation



do cold atoms forget underlying microscopic physics?

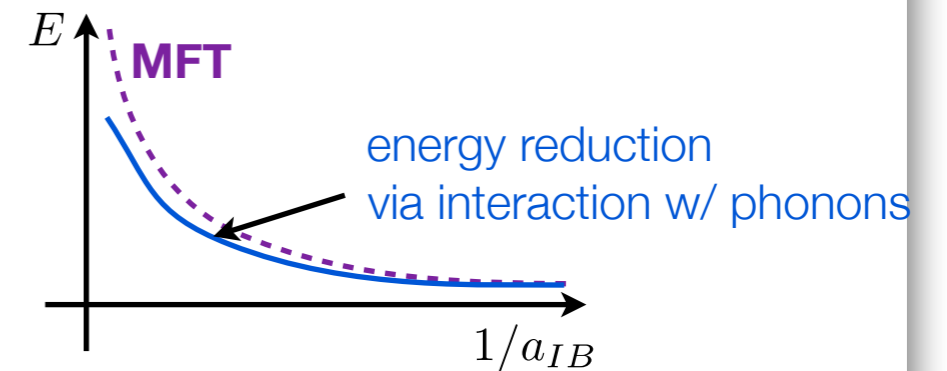
The Question

What is the spectrum of the model?



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do cold atoms forget underlying microscopic physics?

Quantity to address this question in quantum field theory:

Spectral function (gives access to radio-frequency response etc...)

$$\mathcal{A}(\omega, \mathbf{p}) = -2 \text{Im} G^{\text{R}}(\omega, \mathbf{p})$$

T-matrix approximation

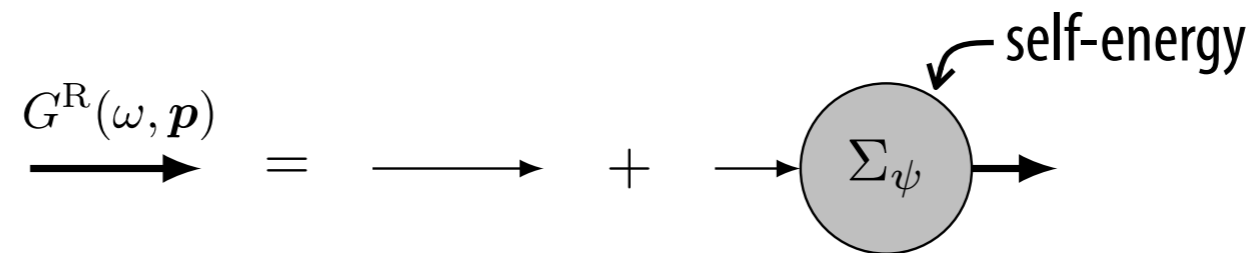
SEE ALSO FOR FERMIONS: **RS, ENSS, PRA 83 (2011)**

impurity spectral function

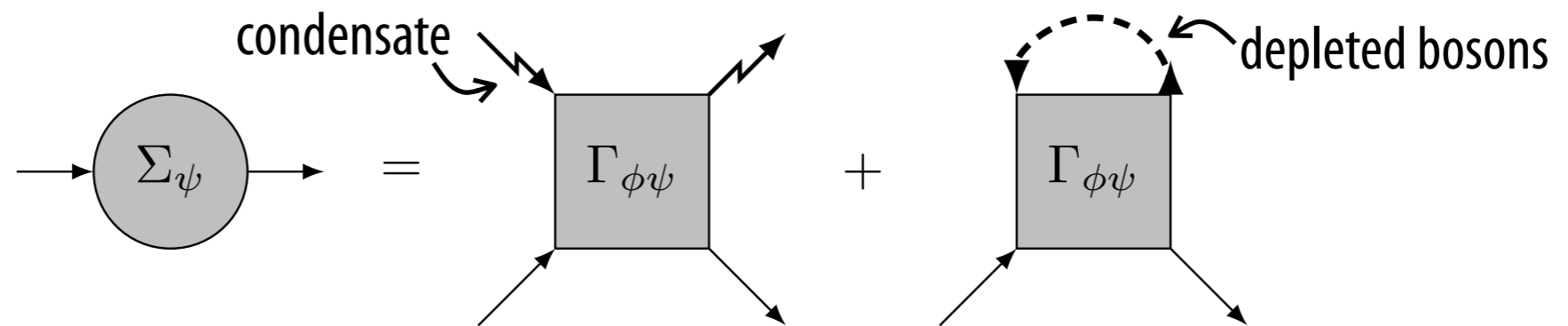
$$A_{\text{pol}}(\omega, \mathbf{p}) = -2 \text{Im} G^{\text{R}}(\omega, \mathbf{p})$$

from Dyson equation

full Green's function:



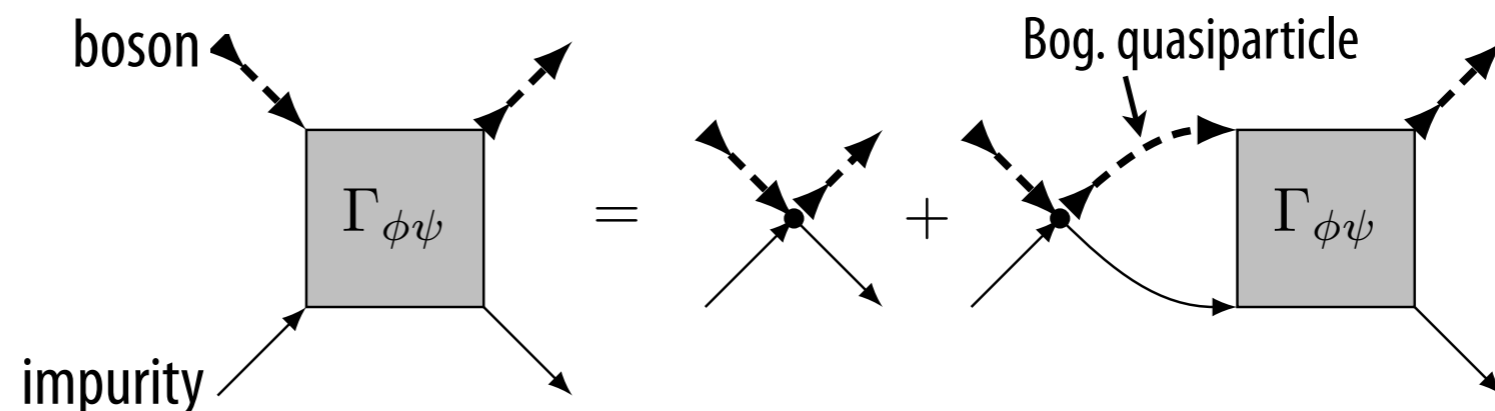
self-energy:



prerequisite: recover exact two-body solution [unlike previous works]

resummed perturbation theory

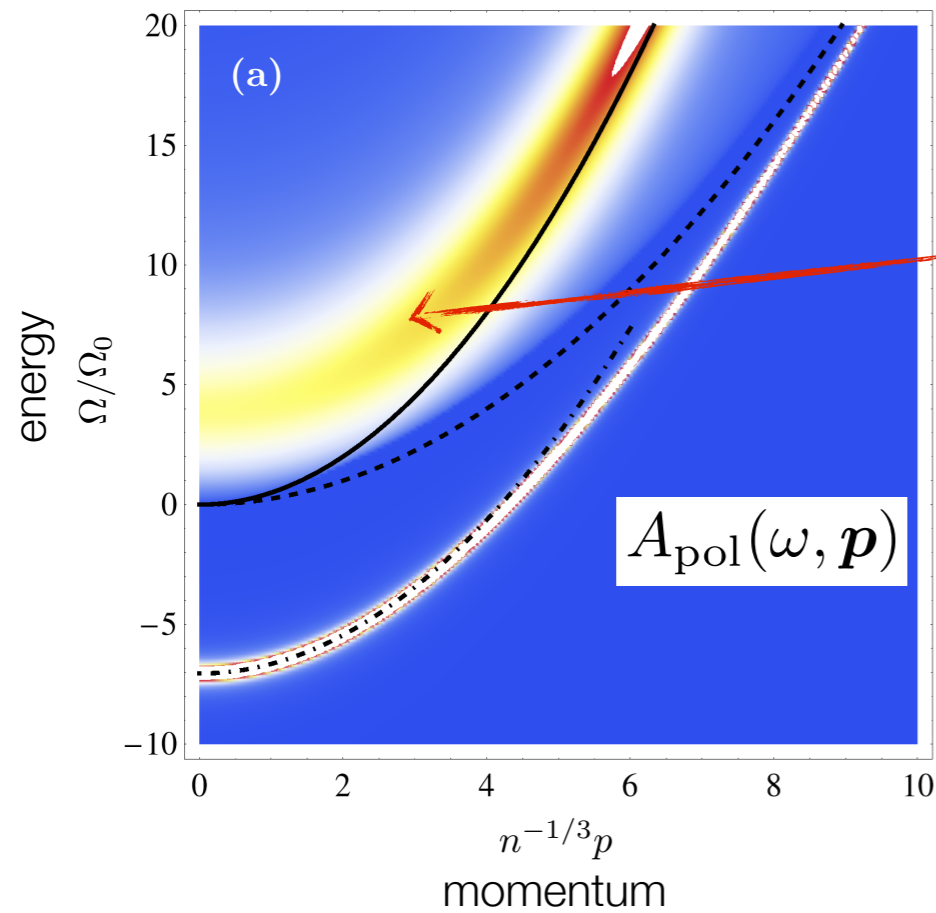
T-matrix equation:



Result for momentum resolved spectral function

$$n^{1/3} a_{IB}^{-1} = 1$$

unlike condmat: **two** coherent quasi-particle excitations!

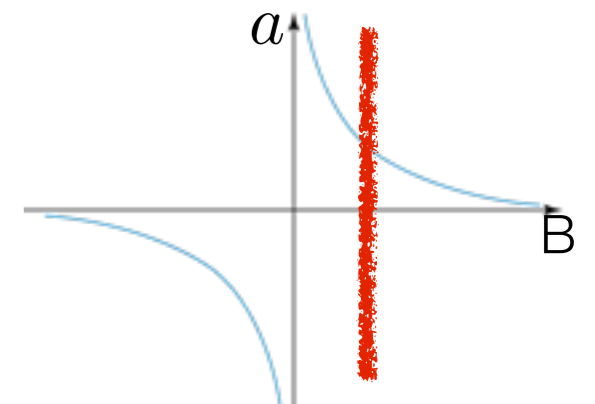


almost “standard” repulsive polaron $a > 0$

- ▶ at positive energy
- ▶ enhanced effective mass
- ▶ finite lifetime!
- ▶ largely reduced quasi-particle weight

RATH, RS, PRA 88 (2013)

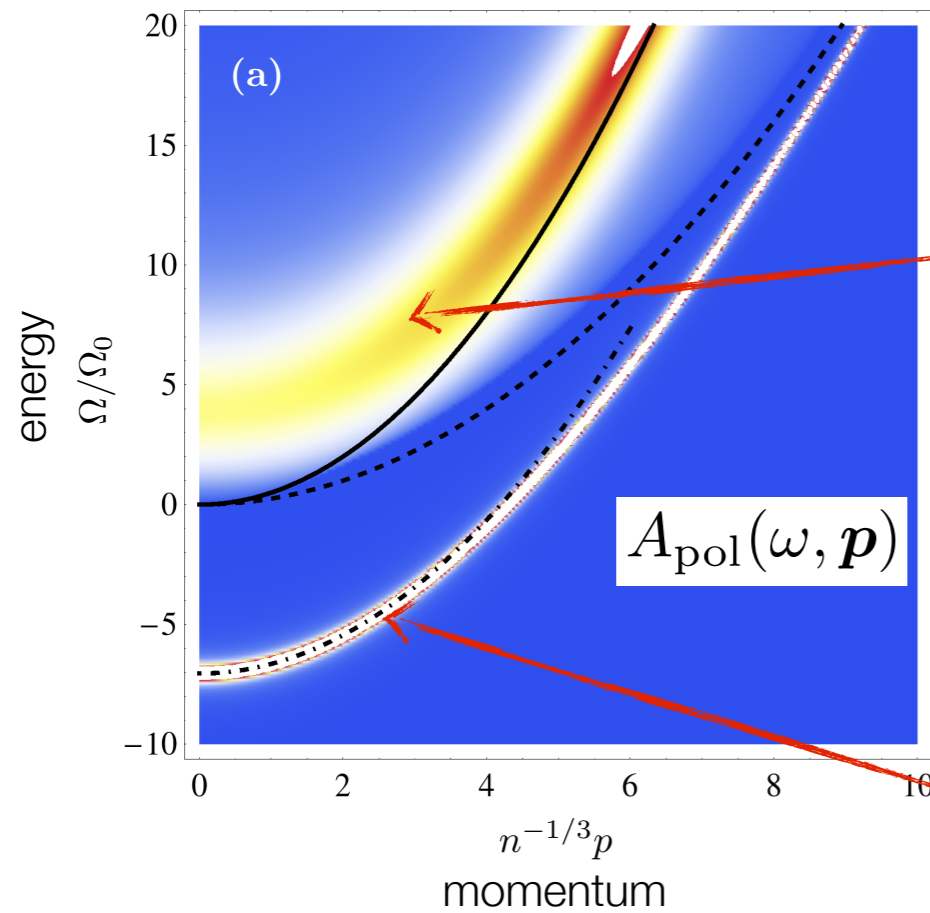
Feshbach resonance



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RATH, RS, PRA 88 (2013)

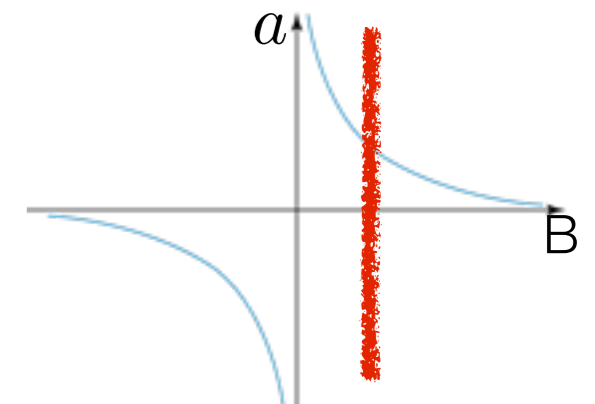
almost “standard” repulsive polaron $a > 0$

- ▶ at positive energy
- ▶ enhanced effective mass
- ▶ finite lifetime!
- ▶ largely reduced quasi-particle weight

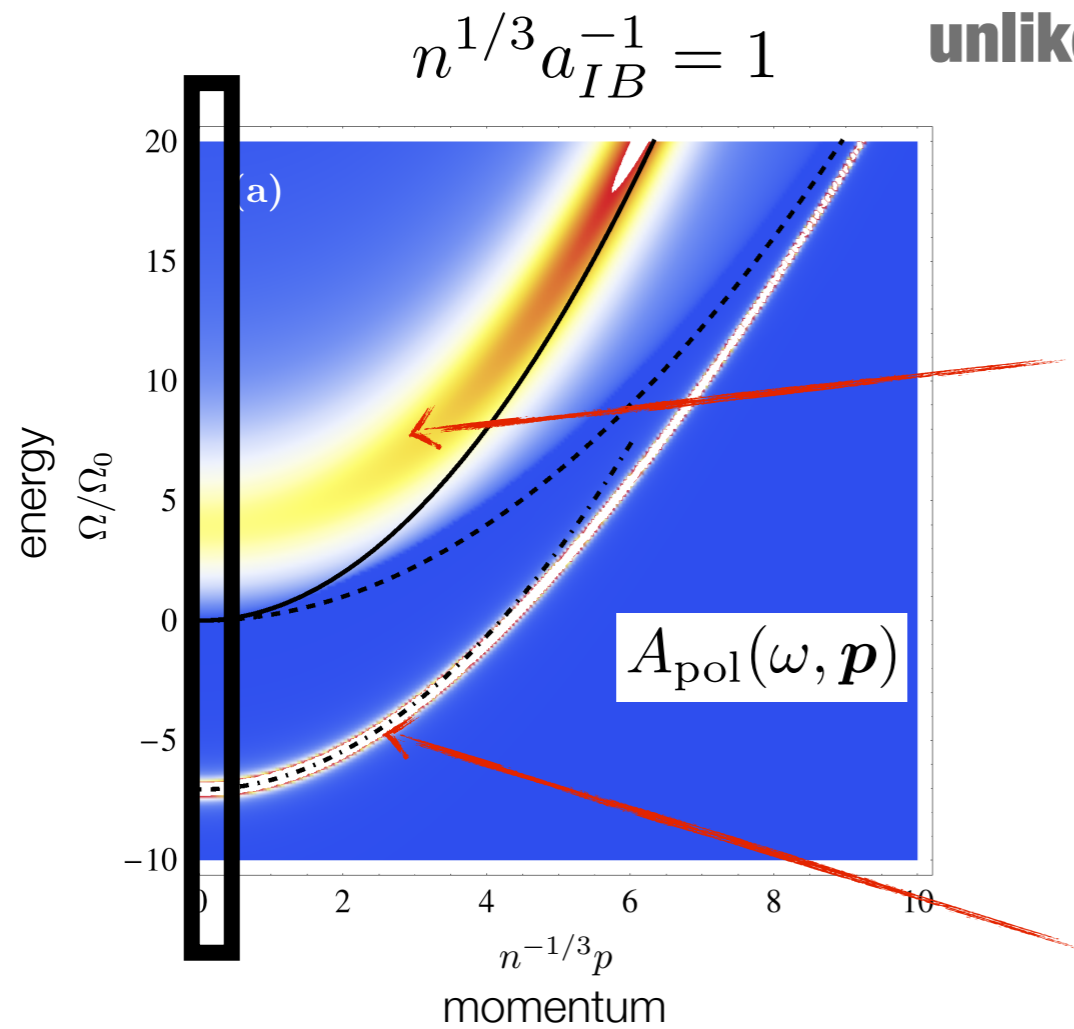
“new” attractive polaron $\forall a$

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Feshbach resonance



Result for momentum resolved spectral function



RATH, RS, PRA 88 (2013)

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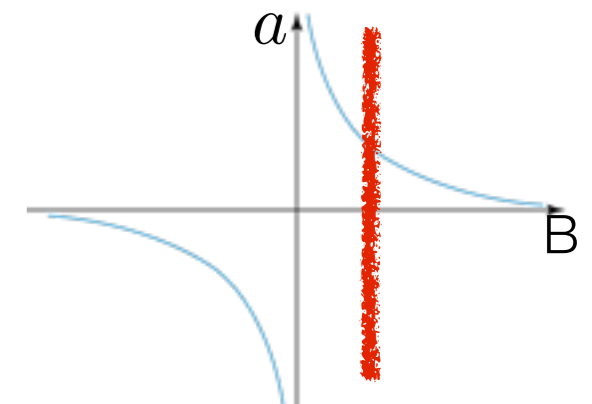
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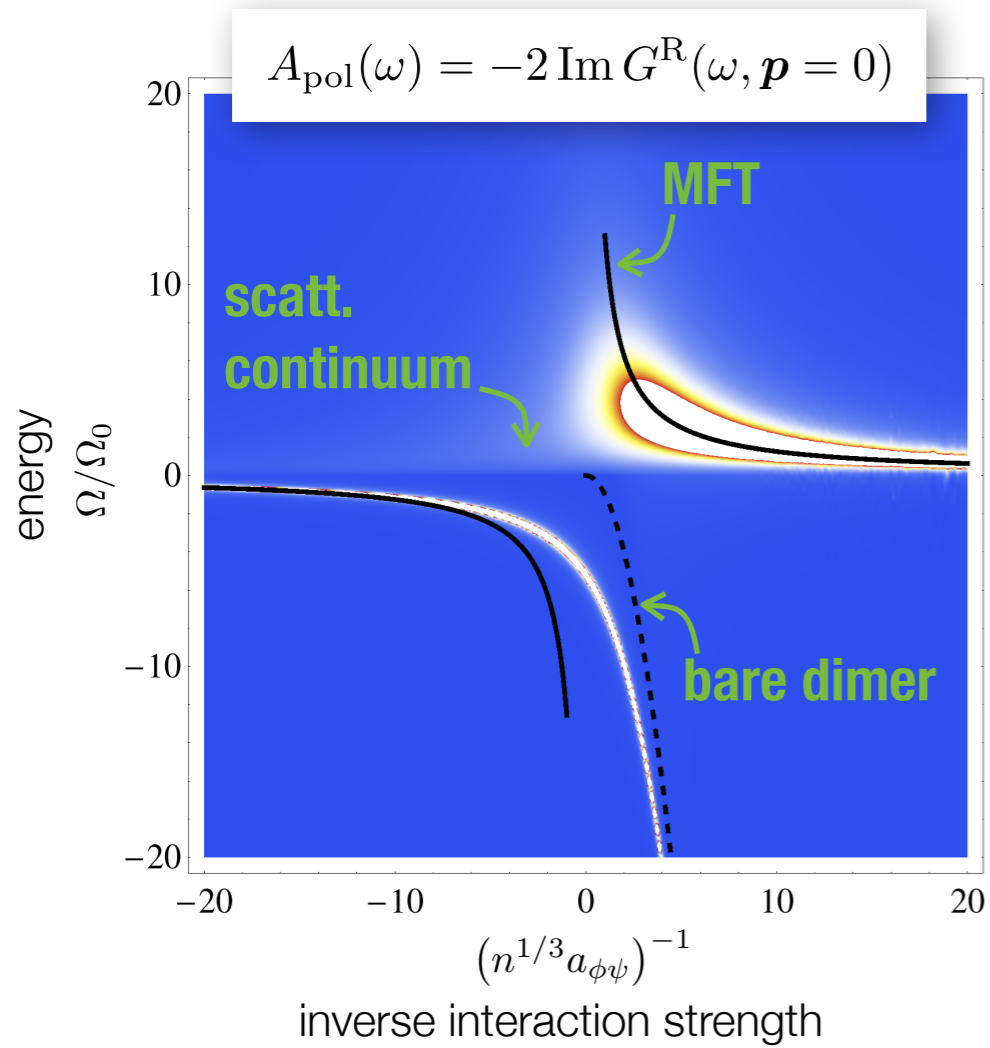
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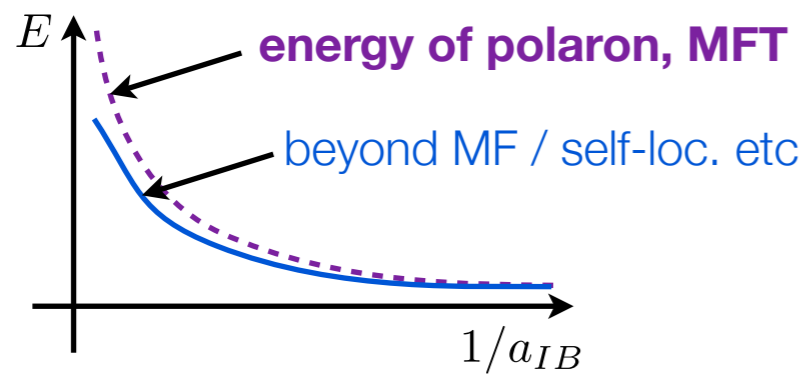
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Energy spectrum for impurity at rest

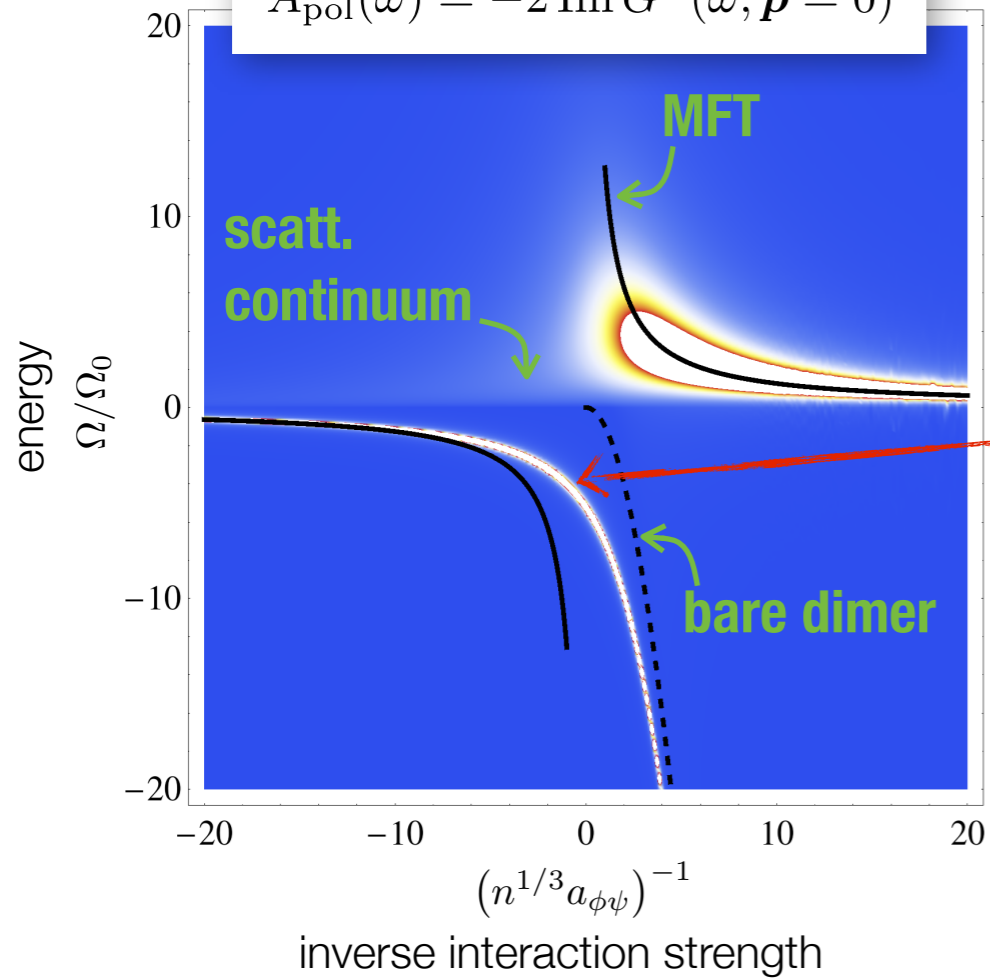


previous “quantum simulation proposals”



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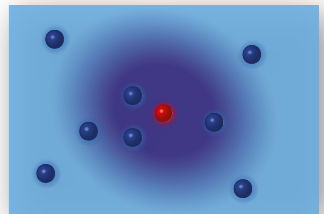
$$A_{\text{pol}}(\omega) = -2 \text{Im} G^{\text{R}}(\omega, \mathbf{p} = 0)$$



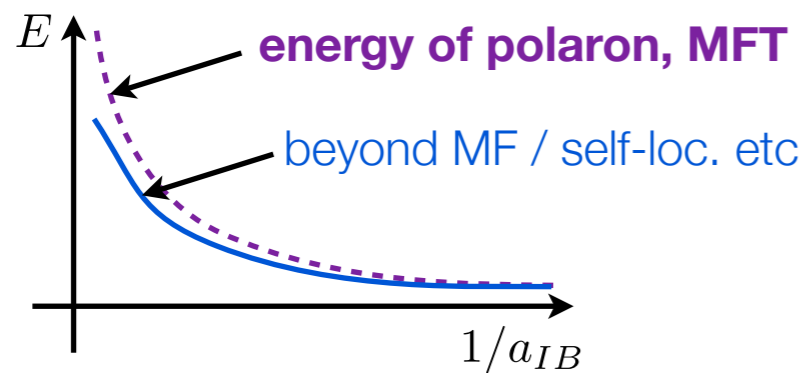
attractive polaron

- ▶ stable ground state at all scattering lengths

OBSERVED AT WEAK COUPLING PFAU GROUP [STUTTGART]: BALEWSKI ET AL., NATURE 502 (2013)

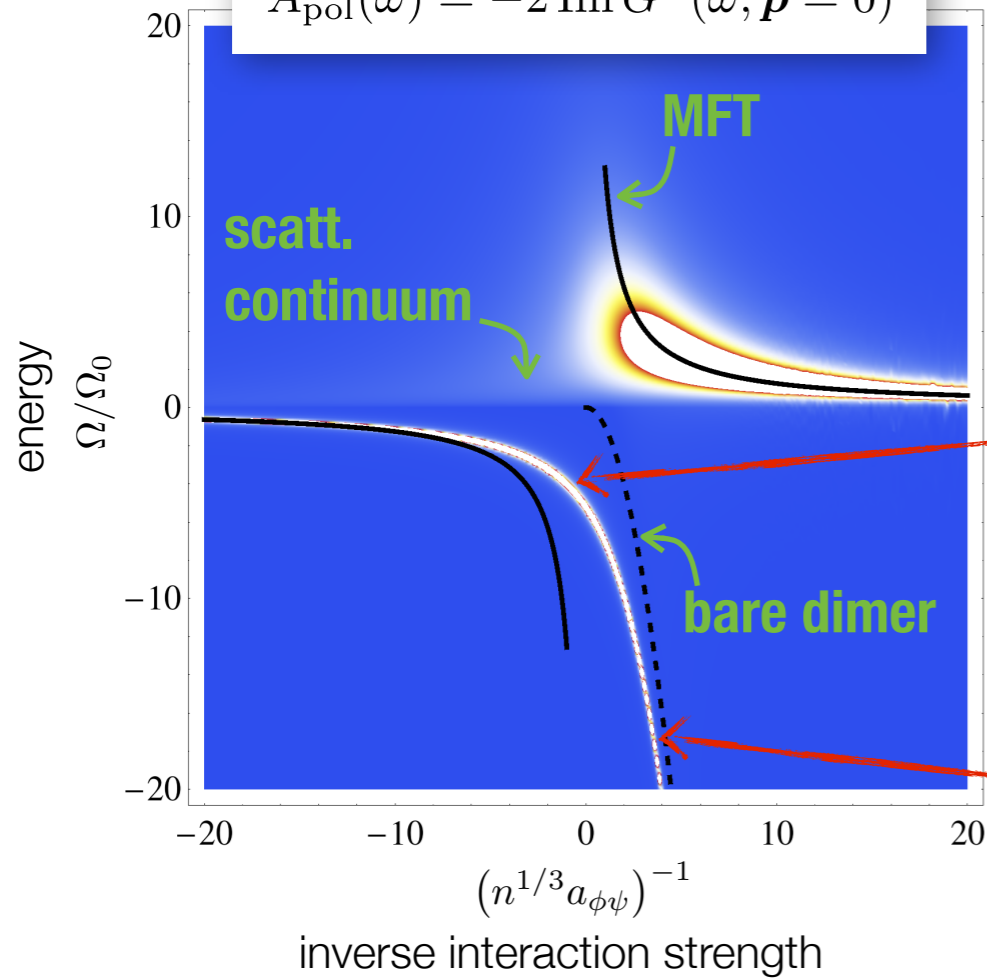


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Energy spectrum for impurity at rest

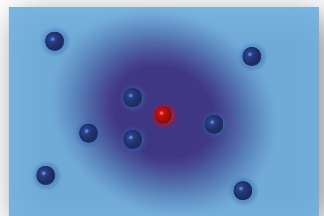
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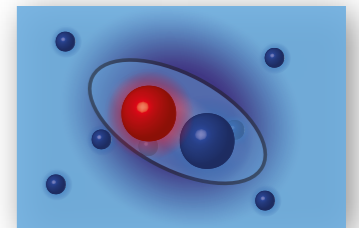
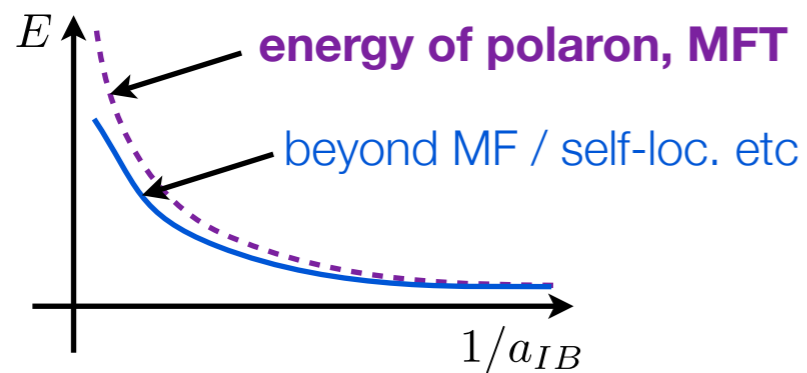
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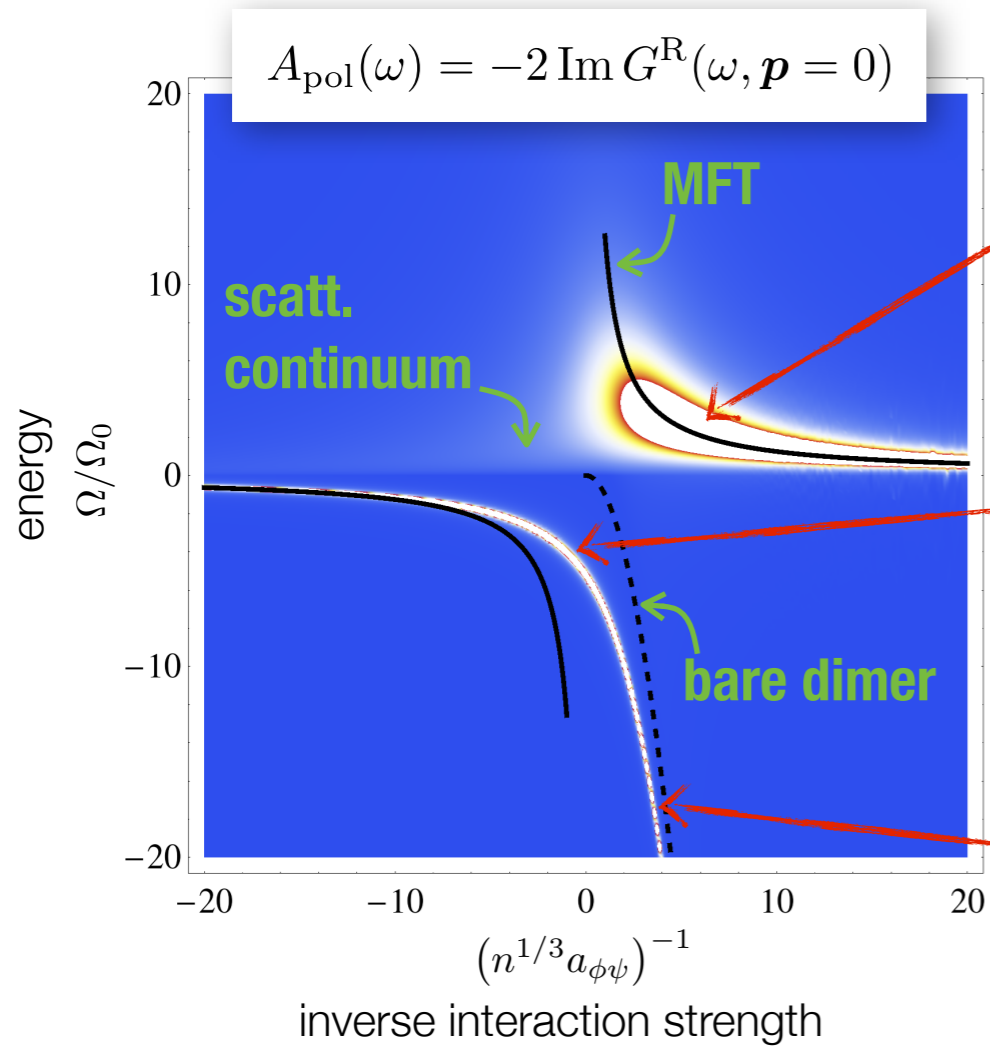
smooth **crossover to molecular state** - hybridization
 - different from transition for Fermi polaron - RATH, RS, PRA 88 (2013)

DISCUSSED IN CONTEXT OF B/F MIXTURES BY
 MARCHETTI, ... PARISH, PRB 78 (2008)

previous "quantum simulation proposals"

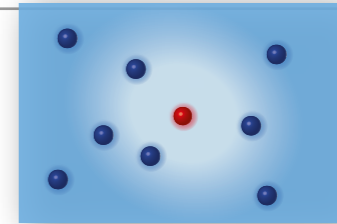


Energy spectrum for impurity at rest



repulsive polaron

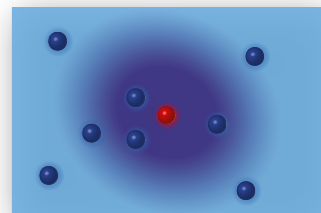
- ▶ extremely unstable in strong-coupling regime!
[molecule formation]
- ▶ self-localization challenging to observe



attractive polaron

- ▶ stable ground state at all scattering lengths

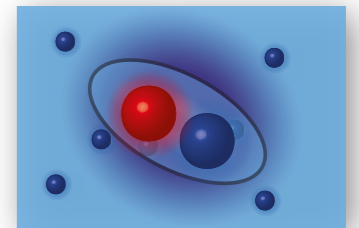
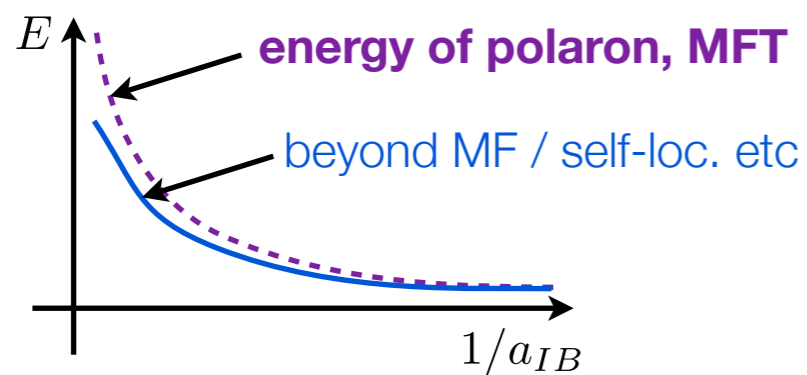
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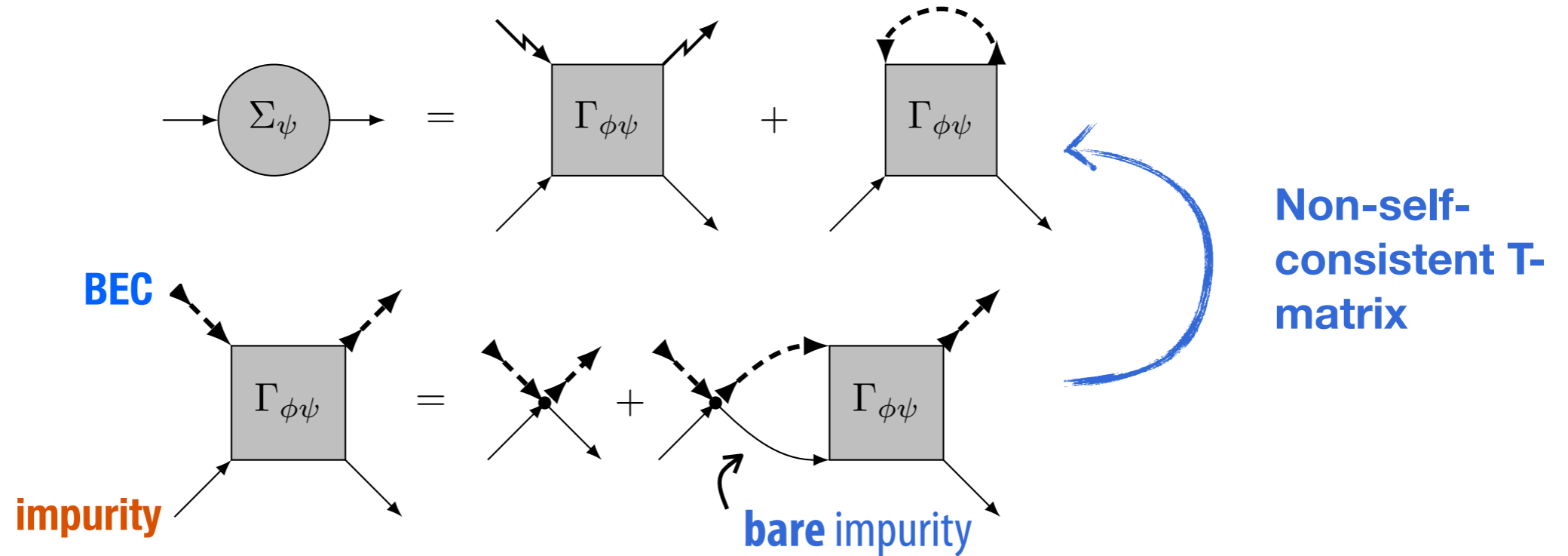
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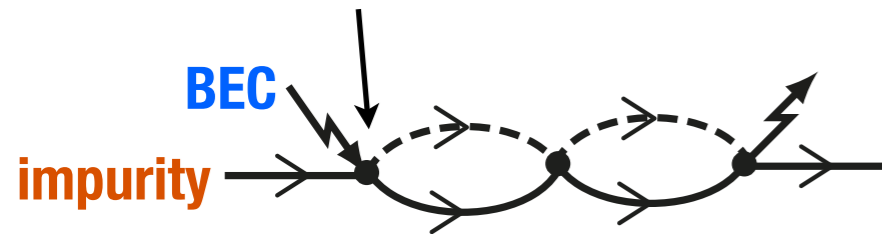


Self-consistent T-matrix

So far: Non-selfconsistent T-matrix approach

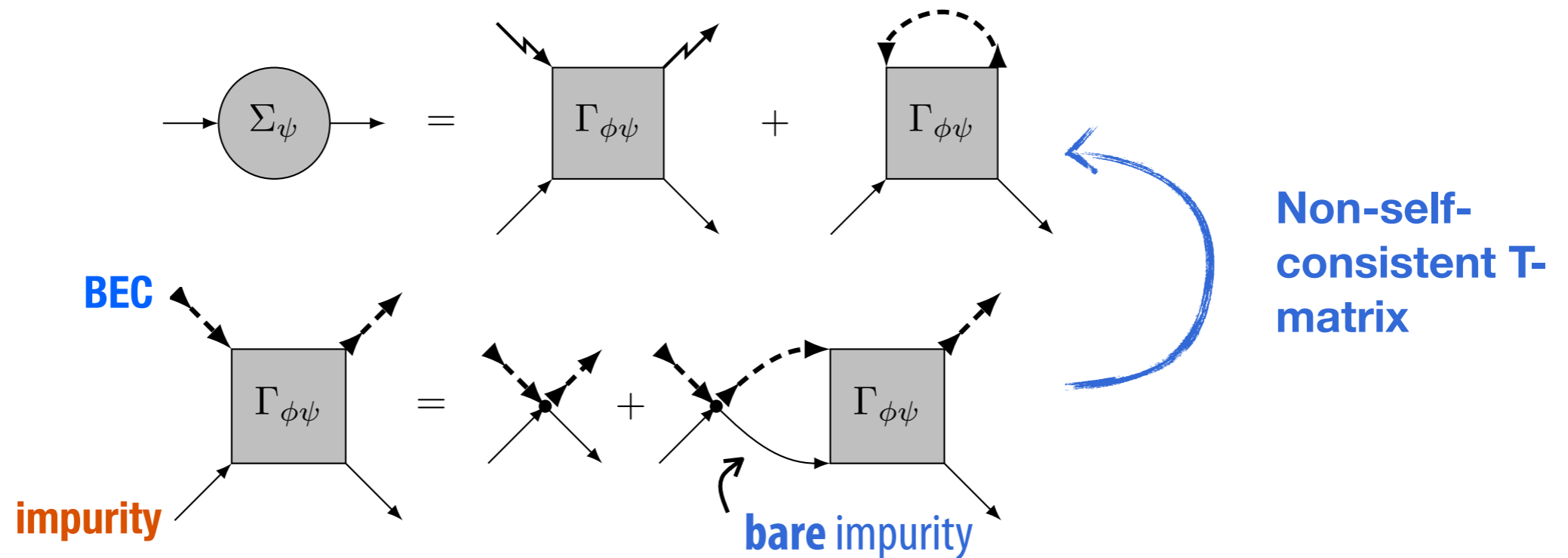


► single boson taken out of condensate

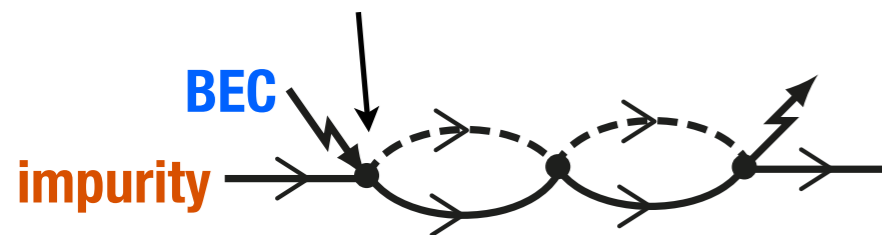


Self-consistent T-matrix

So far: Non-selfconsistent T-matrix approach



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- ▶ equivalent to **simple variational wave function**

$$|\psi_0\rangle = \sqrt{Z} \hat{c}_0^\dagger |\text{BEC}\rangle + \sum_{\mathbf{k}} \mathcal{A}(\mathbf{k}) \hat{c}_{-\mathbf{k}}^\dagger \hat{b}_{\mathbf{k}}^\dagger |\text{BEC}\rangle$$

impurity

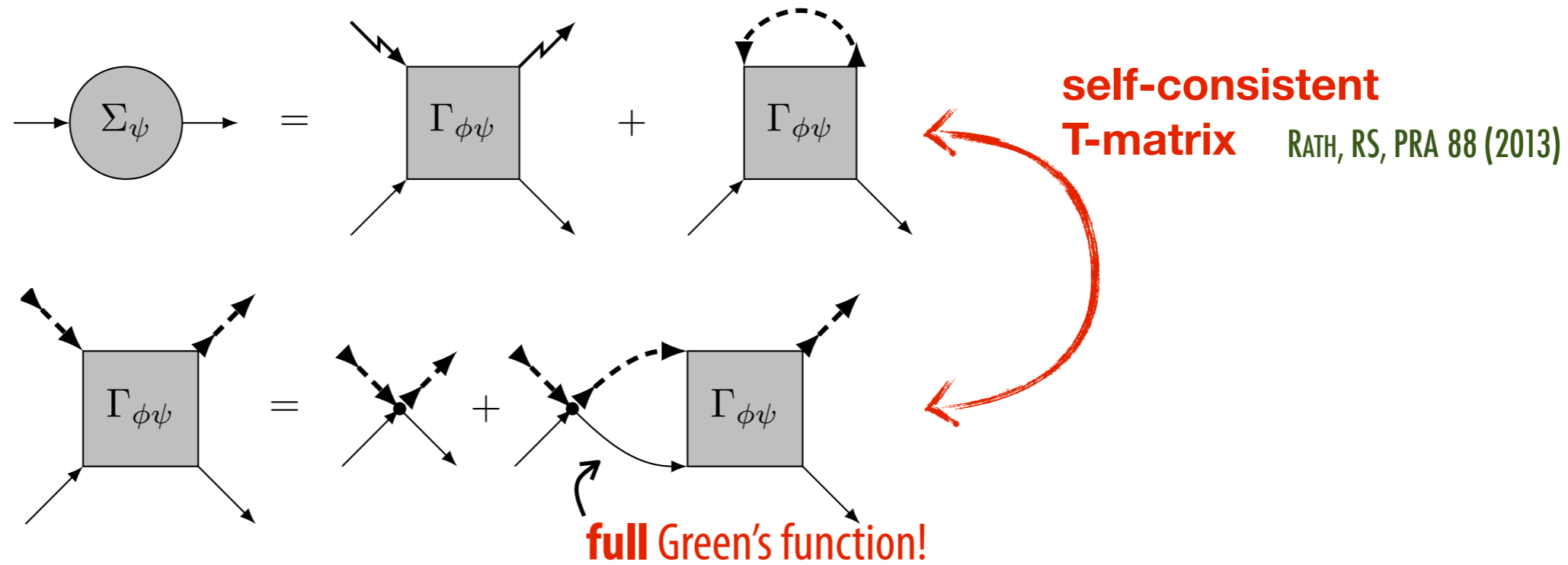
boson excited out of condensate

- ▶ captures very simple 'entanglement' between BEC and impurity

recently studied in: LI AND DAS SARMA, ARXIV:1404.4054

Self-consistent T-matrix

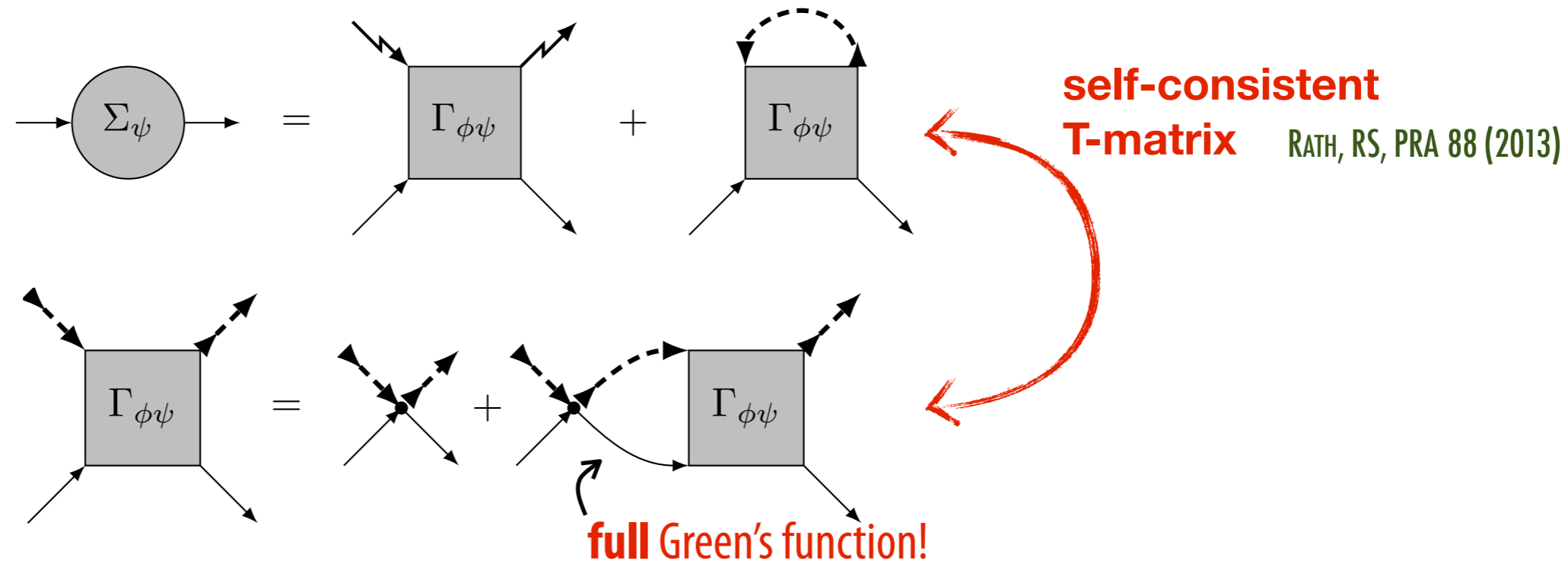
Selfconsistent T-matrix approach



- ▶ solved numerically using algorithm developed for *functional renormalization group approach* for RG flow of full spectral functions RS, ENSS, PRA 83 (2011)

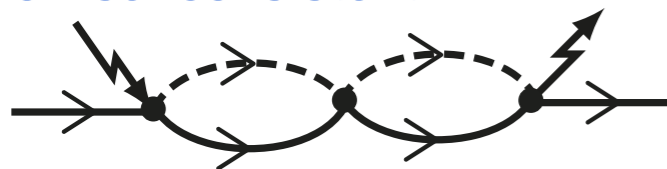
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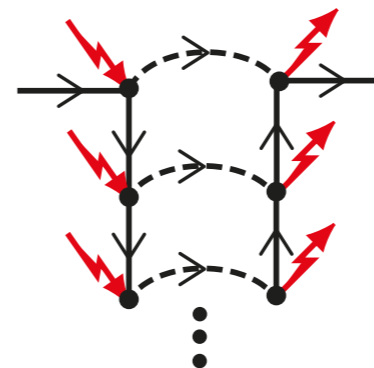
- ▶ solved numerically using algorithm developed for *functional renormalization group approach* for RG flow of full spectral functions
RS, ENSS, PRA 83 (2011)
- ▶ accounts for *infinitely many virtual excitations* of bosons out of the coherent condensate state

Non-selfconsistent:



- ▶ **single boson taken out of condensate**

Selfconsistent:

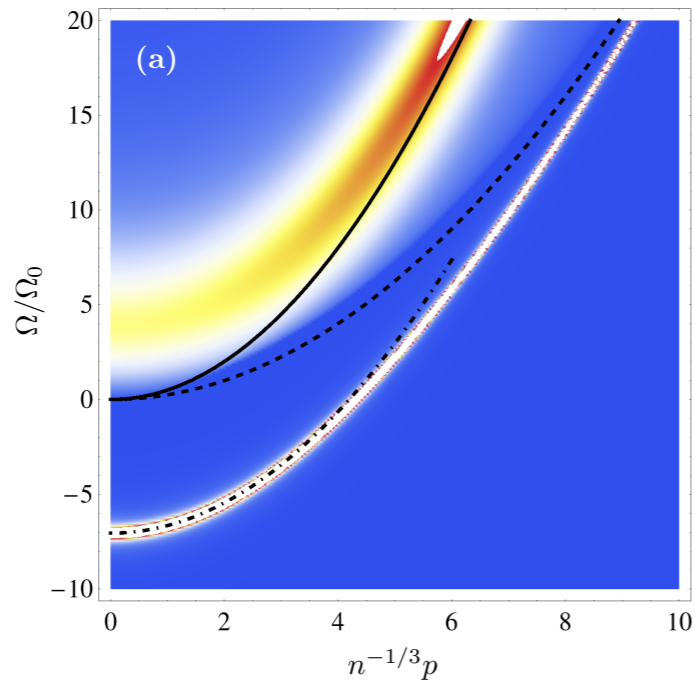


- ▶ **infinite number of bosons taken out of condensate** - way beyond product wave functions for BEC

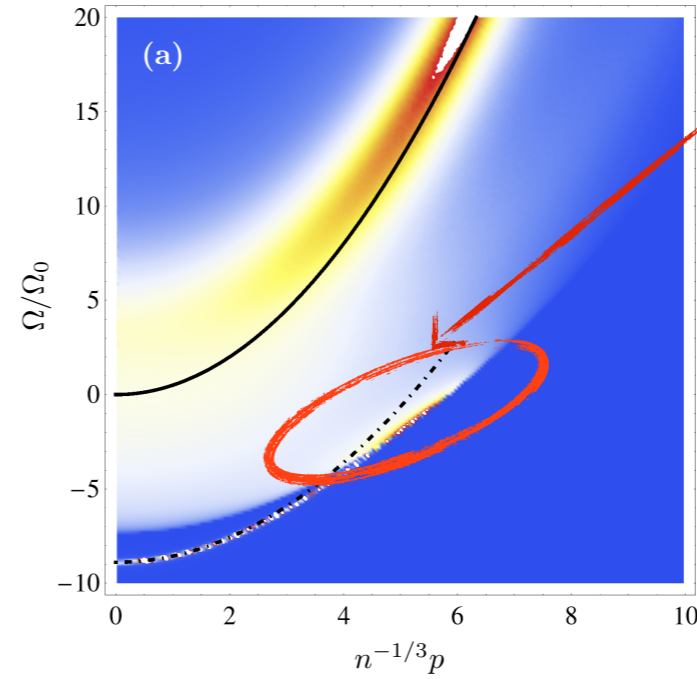
Self-consistent T-matrix - Results

Momentum
resolved
Spectral
function

Non-selfconsistent



Selfconsistent



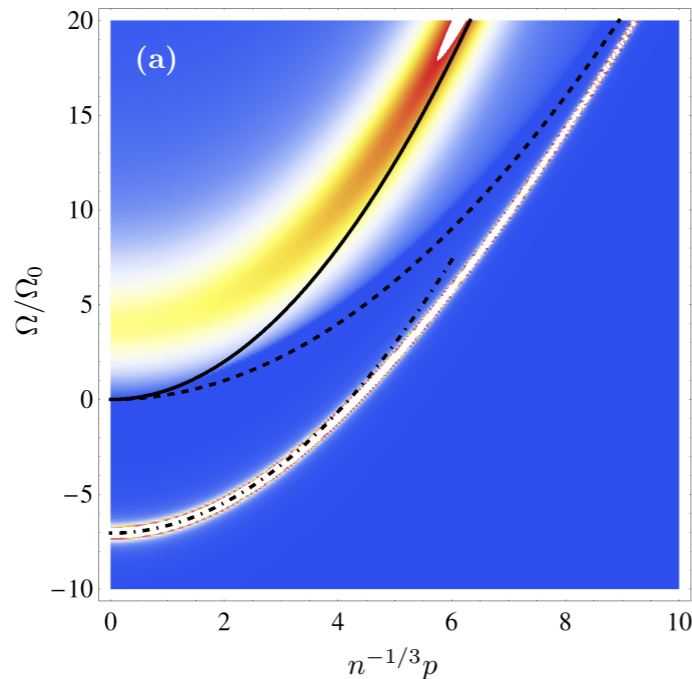
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*momentum
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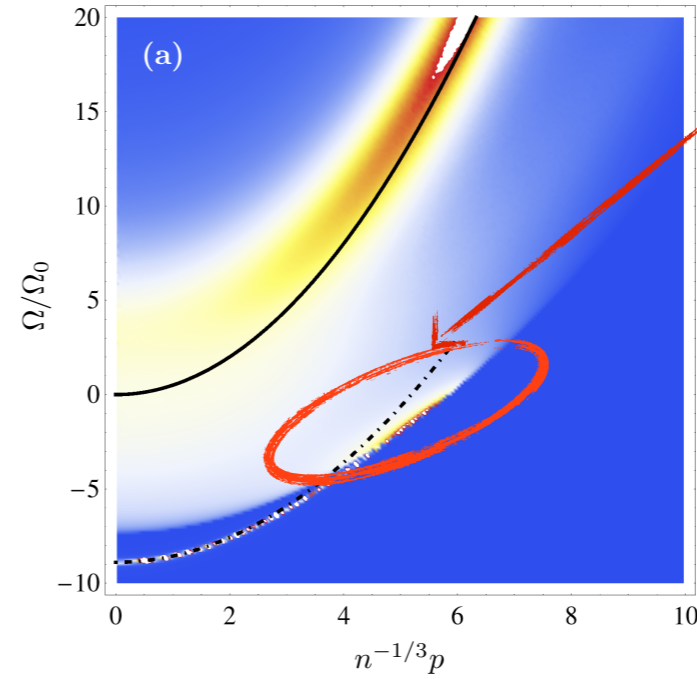
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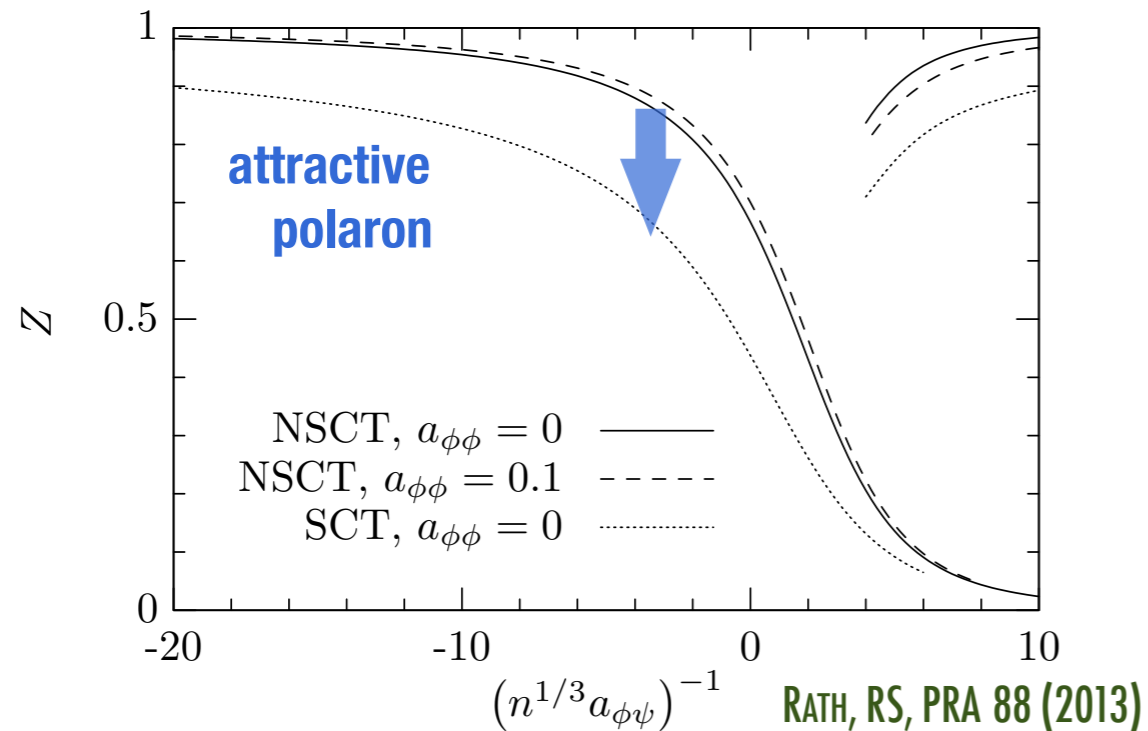
Selfconsistent



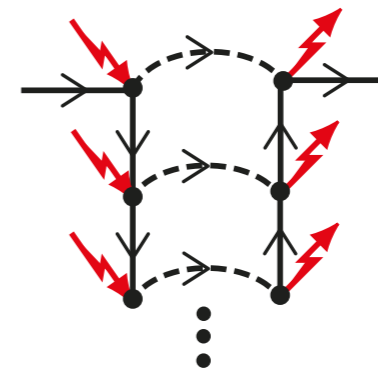
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Suppression of quasi-particle weight



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► for Fermi polaron much smaller deviation!

PROKEFIEV, SVISTUNOV, PRB 77 (2008)

RS, ENSS, PRA 83, 063620 (2011)

Proposal for experimental observation

E.G. $^{40}\text{K}/^{41}\text{K}$ MIXTURE AT $B=543\text{ G}$

SEE MIT GROUP: WU ET AL. PRA 84 (2012)

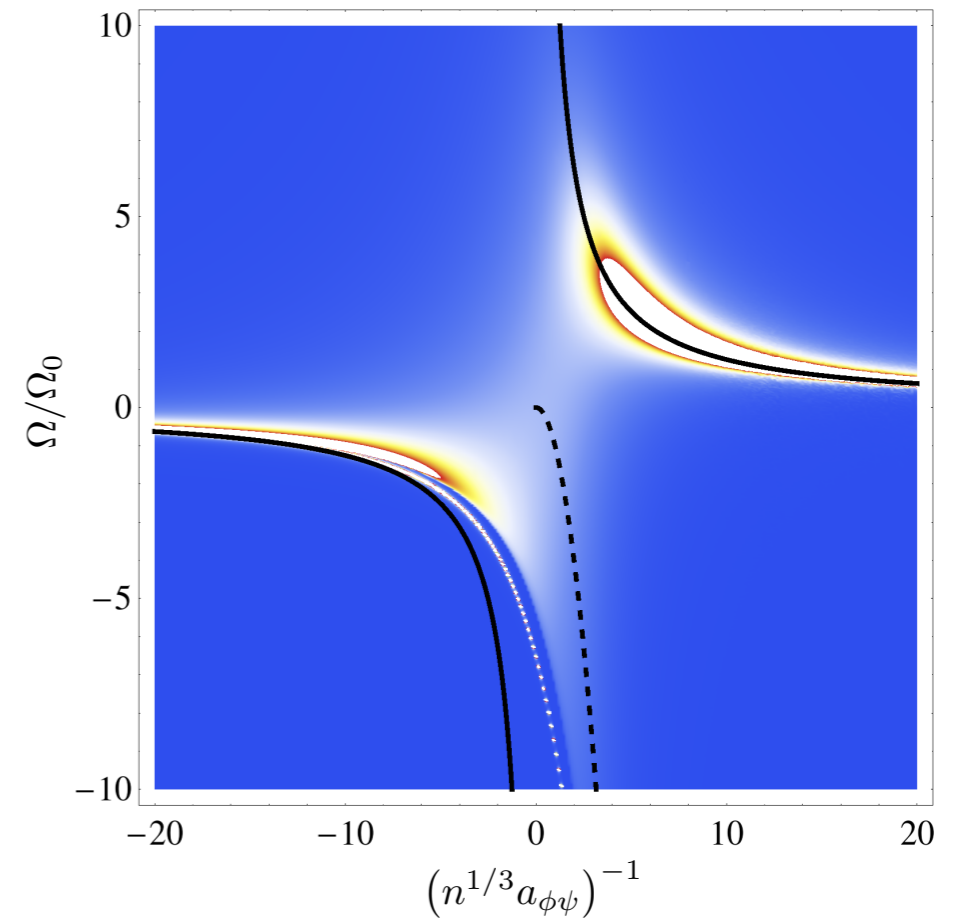
Challenge

▶ Efimov effect + statistics: Bose-Fermi mixtures unstable due to enhanced three-body recombination

SEE E.G. RS, RATH, ZWERGER, EJB 85 (2012)

▶ possible BEC deformation due to large interactions

EFIMOV IN SPIN-ORBIT BOSONS: SHI, CUI, ZHAI PRL 112 (2014)



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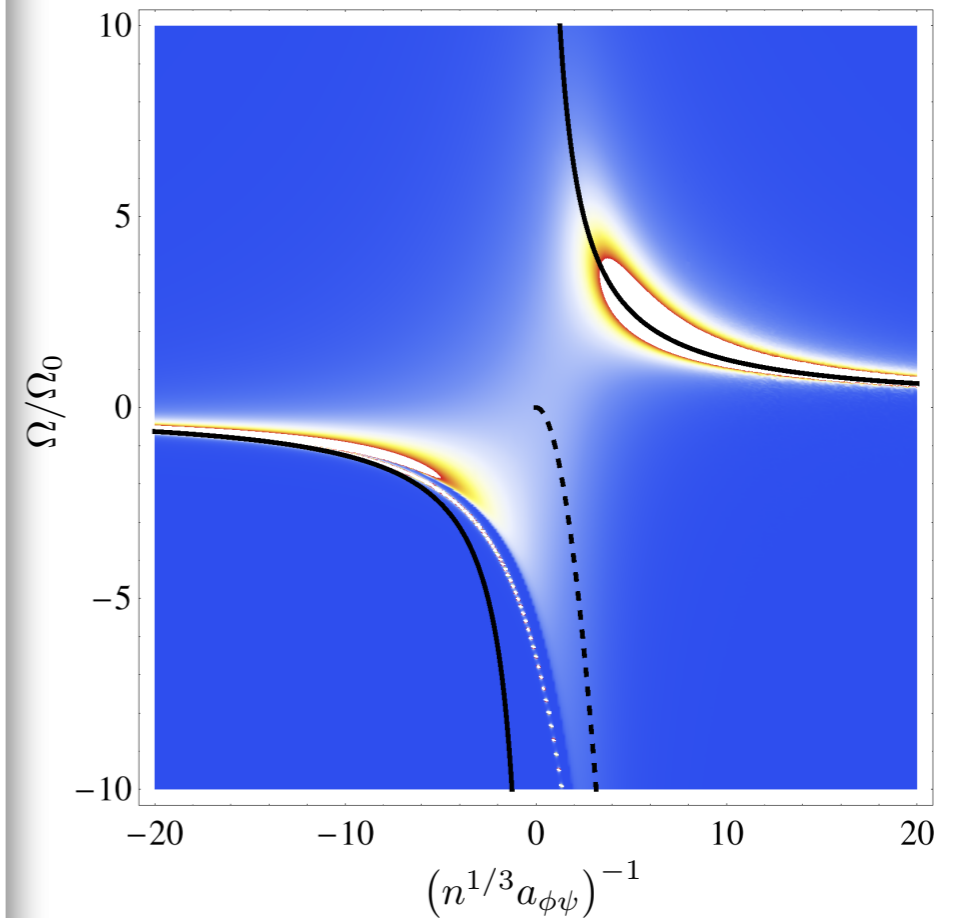
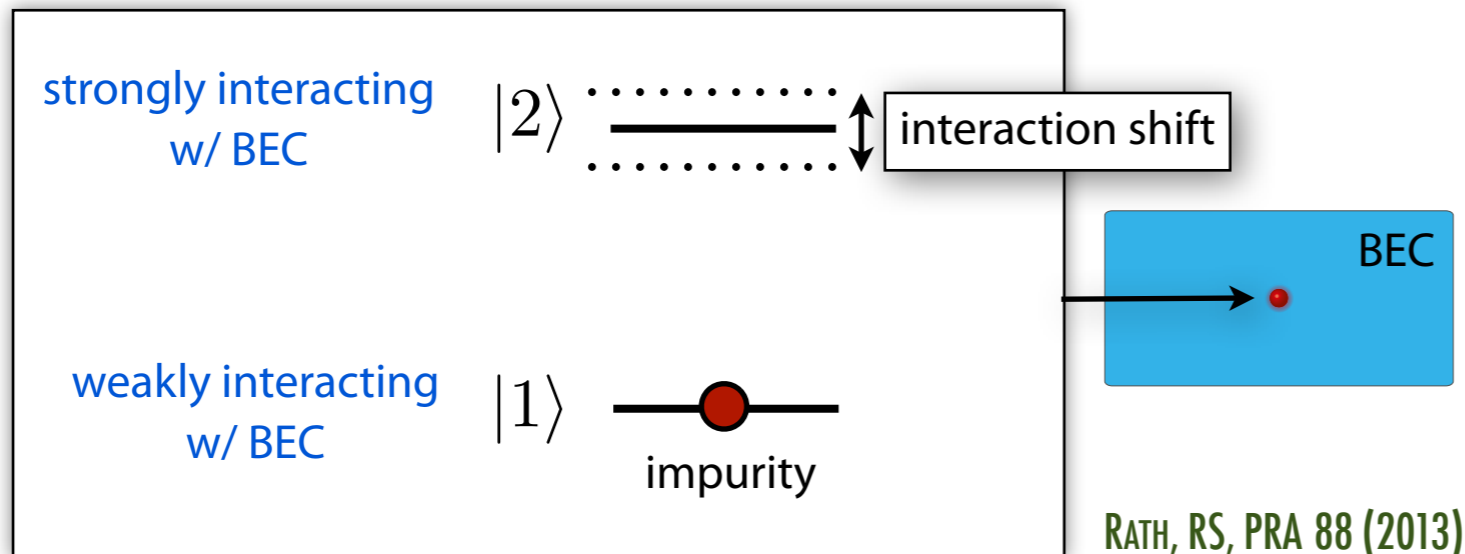
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Resolution: Inverse RF spectroscopy



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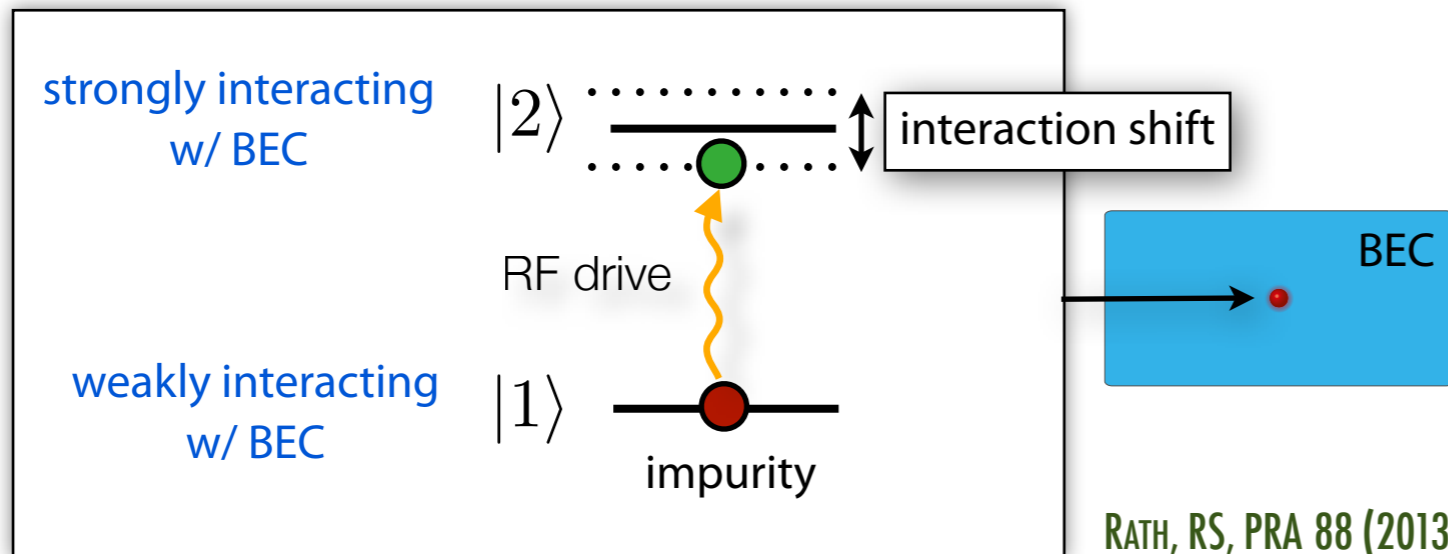
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RATH, RS, PRA 88 (2013)

THEORY: RS, ENSS, PRA 83, 063620 (2011)

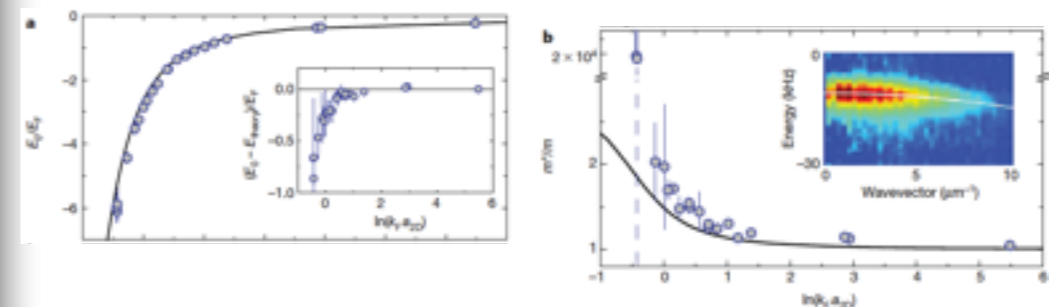
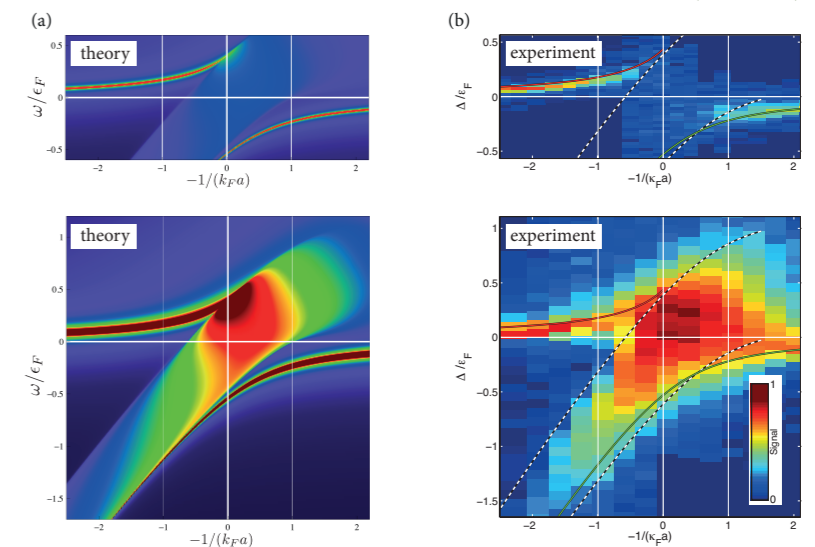
MASSIGNAN, BRUUN, EPJD 65 (2011)

RS, ENSS, PIETILA, DEMLER, PRA 85, 021602 (2012)

EXPERIMENTS: KOHSTALL ET AL., NATURE 485, 615 (2012)

KOSCHORREK ET AL., NATURE 485, 619 (2012)

- ▶ maps out impurity spectral function
- similar procedure proposed and used for fermions



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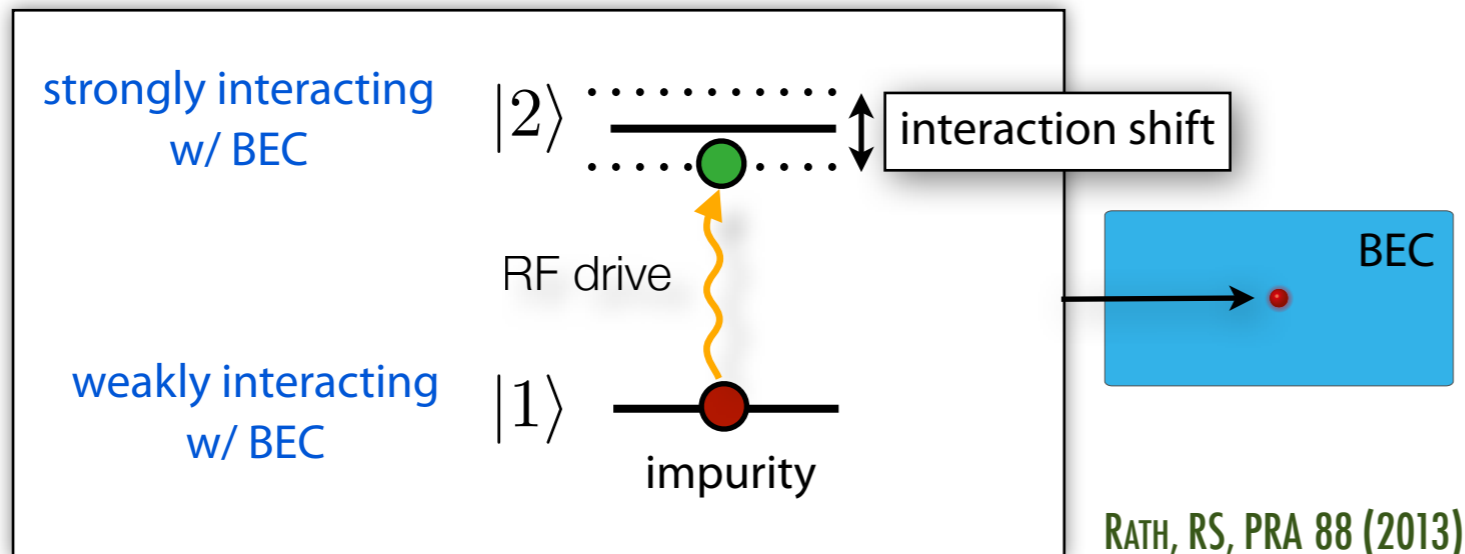
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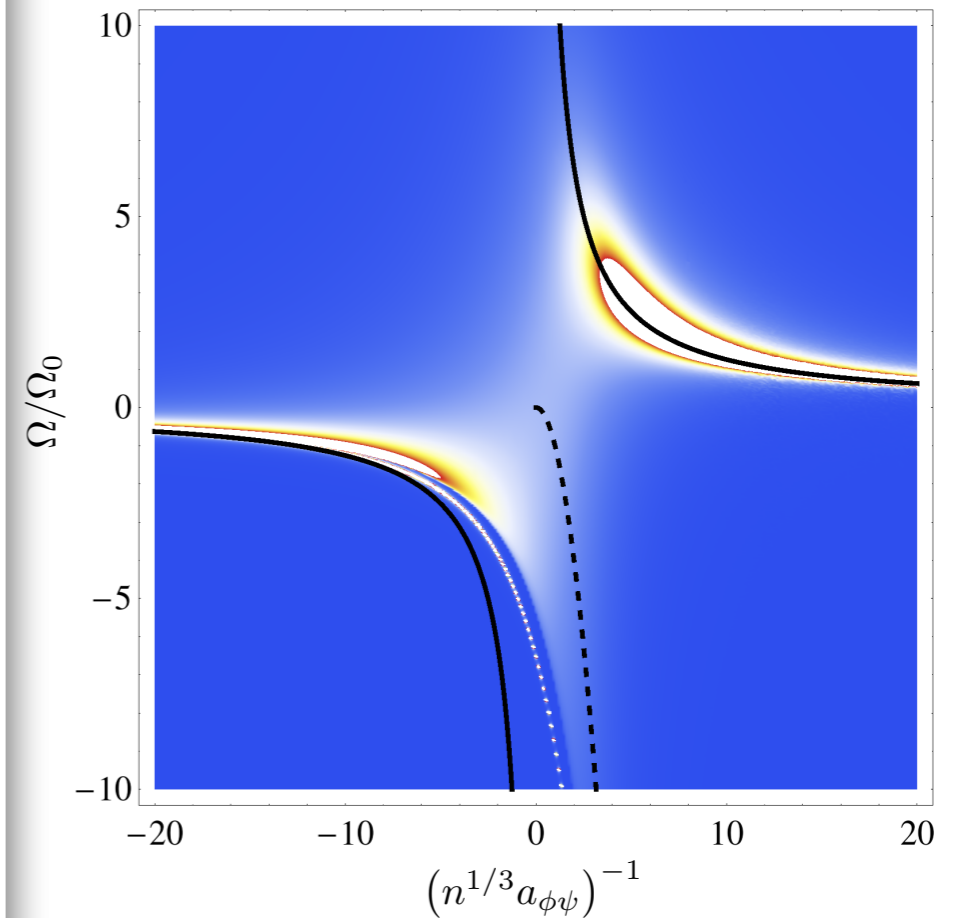
- ▶ possible BEC deformation due to large interactions

Resolution: Inverse RF spectroscopy



- ▶ maps out impurity spectral function
- similar procedure proposed and used for fermions
- ▶ Efimov states off-resonant
- ▶ BEC deformation irrelevant as

$$\tau_{BEC} = \frac{\hbar}{\mu_B} = \frac{\hbar}{g_B n_B} \gg \tau_{\text{Rabi}}$$



Summary

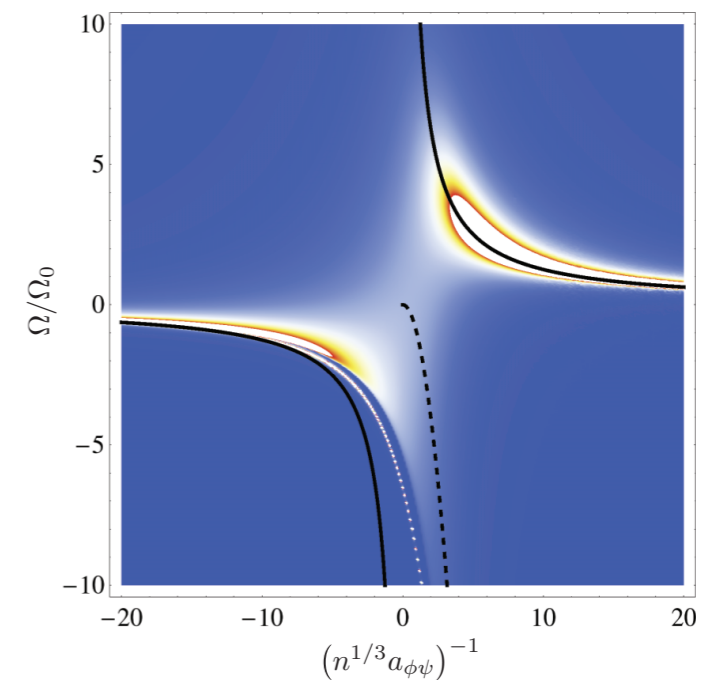
▶ Using a field theoretical approach we studied impurity in ultracold BEC close to Feshbach resonance

▶ Impurity is dressed by bosonic excitations \Rightarrow Bose polaron

▶ Spectrum exhibits two distinct quasi-particle branches

➔ attractive polaron

➔ repulsive (Fröhlich) polaron



▶ We predict finite lifetime of repulsive polaron:

Quantum simulation of Fröhlich model challenging with ultracold atoms

▶ We propose experimental procedure to measure polaron properties via radio-frequency spectroscopy

Outlook

1 non-equilibrium physics

- ▶ Our theory describes the polaron right after the drive to the final state
Here the repulsive polaron is in a highly excited, non-equilibrium state
- ▶ What happens on time scales longer than those of RF experiments?

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↳ Repulsive Bose polaron as probe of nonequilibrium physics

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 - ▶ Detailed study of impurity-molecule crossover

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Thank you!

