

Recent progress of 2D topological insulators in both experiments and theories



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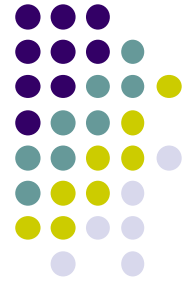
School of Physics, Nanjing University

June 17, 2015

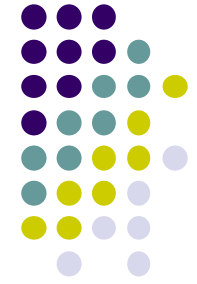


Outline

- Background
 - The family of Hall effects
- 2D topological insulators in theories and experiments
 - Quantum Wells type
 - Graphene type
 - 2D limit from 3D topological insulators
- Outlook
- Summary



New states of matter in condensed matter physics



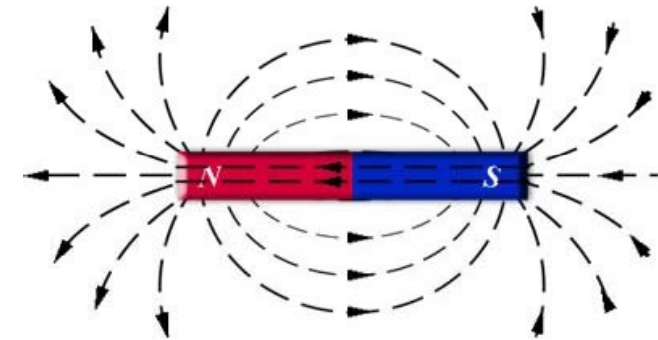
We have known many kinds of fundamental states of matter, including metals, insulators, superconductors, magnets, ... And most of them can be differentiated by **breaking symmetries**.



Crystal: break the translation symmetry

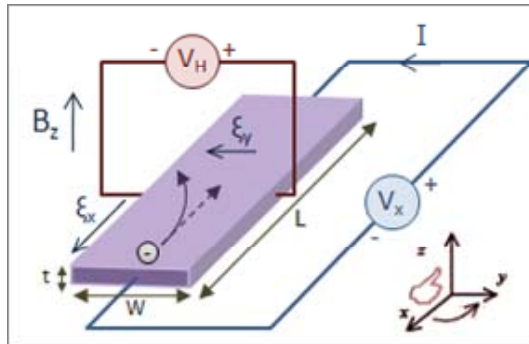


Superconductors: break the U(1) gauge symmetry

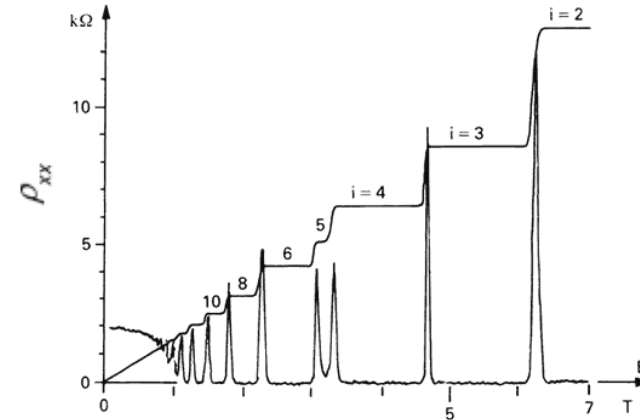


Magnets: break the rotation symmetry

The Hall effect and quantum Hall effect

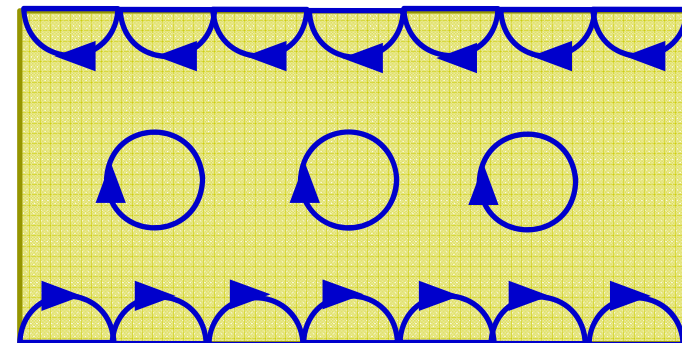
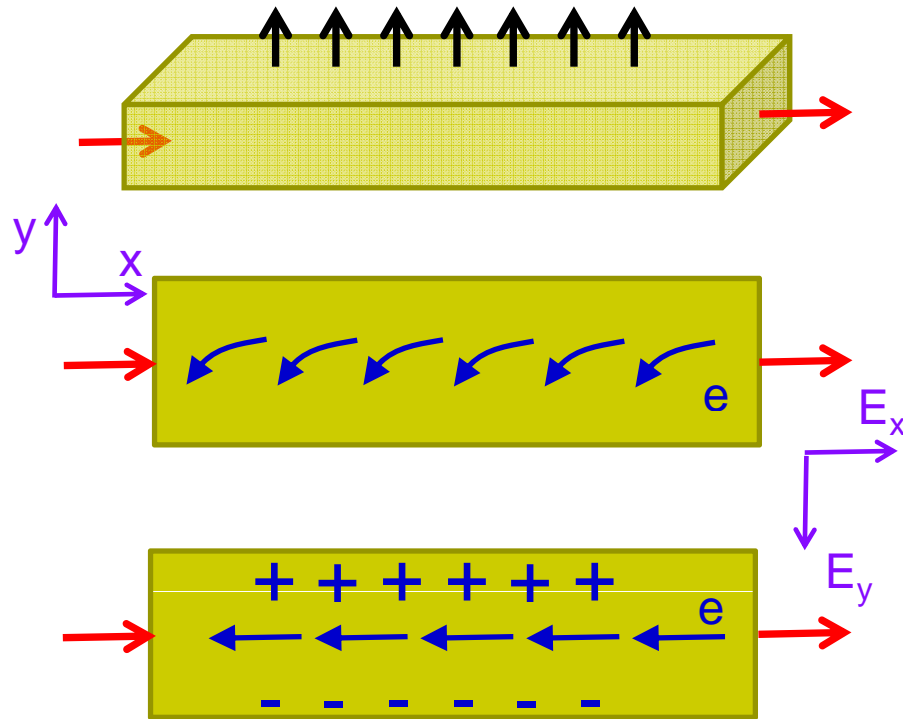


$$R_H = \frac{E_y}{j_x B}$$

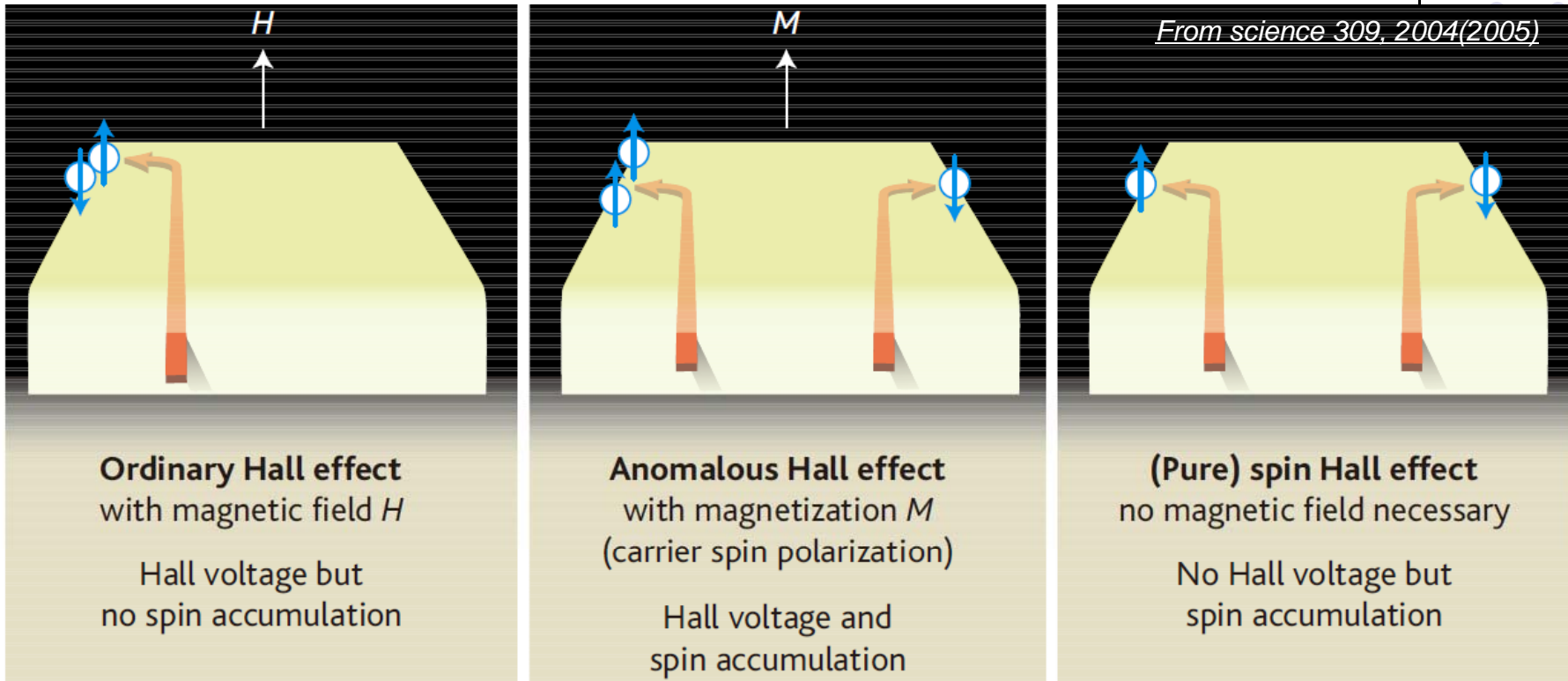


$$\sigma = \nu \frac{e^2}{h}$$

In mathematics, known as the first Chern number, related to Berry's phase.



The family of Hall effects



From science 309, 2004(2005)

Quantum Hall
effect 1980

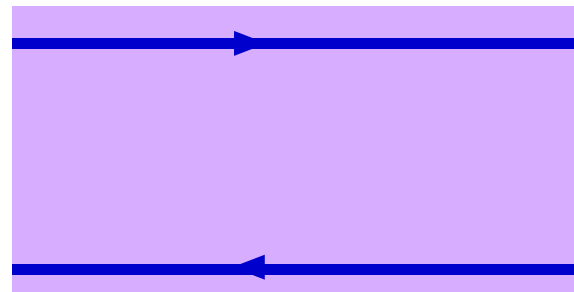
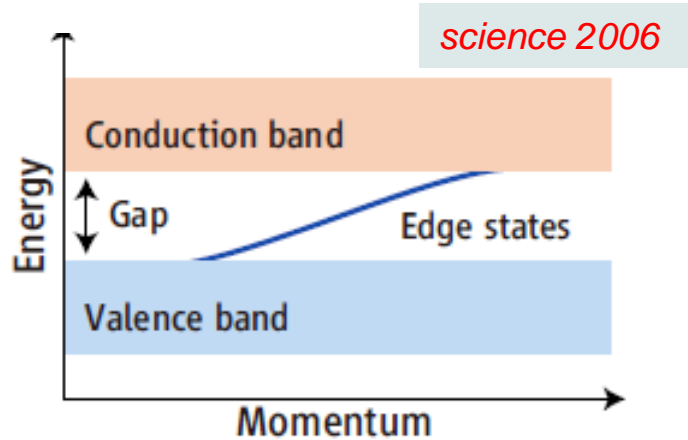
Quantum
Anomalous Hall
effect 2013

Quantum Spin
Hall effect 2007

The quantum anomalous Hall (QAH) states



✓ No external magnetic field

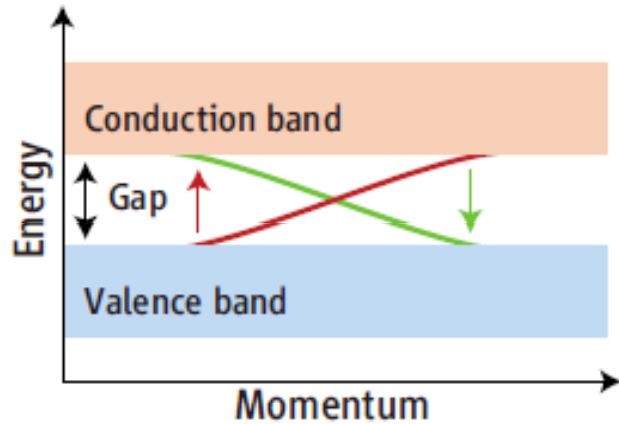


Dissipationless charge current!

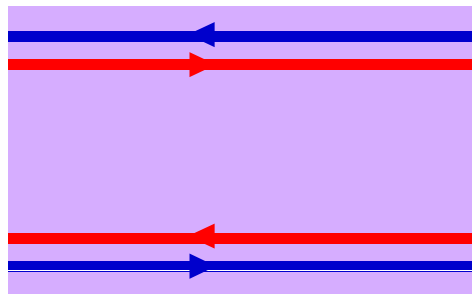
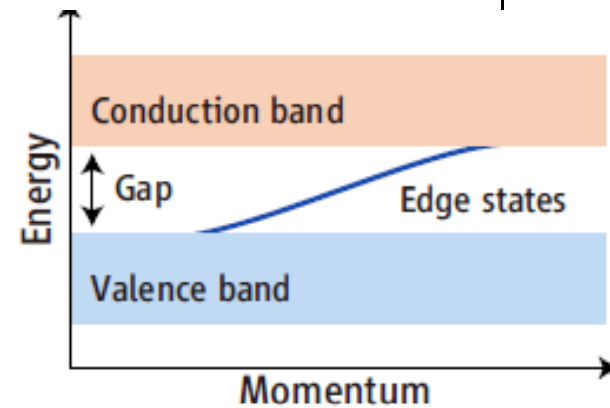


In QAH systems, electrons move just like cars on a highway!

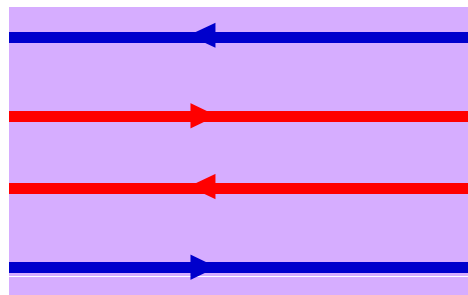
QAH states could get born in quantum spin Hall (QSH) states



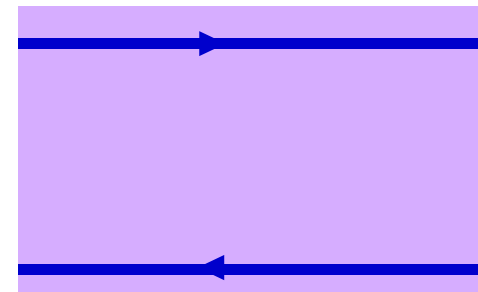
Time reversal
symmetry breaking



Time reversal
symmetry



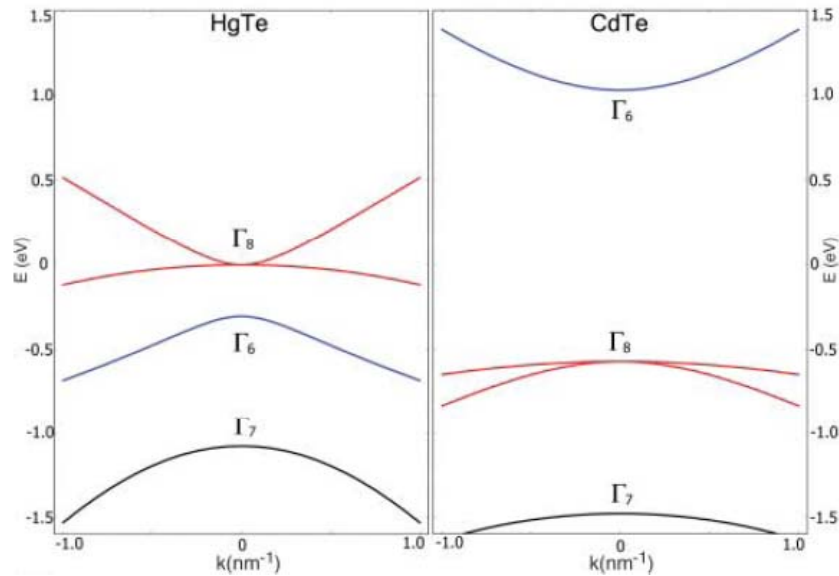
Time reversal Symmetry breaking





2D topological insulators in quantum Wells

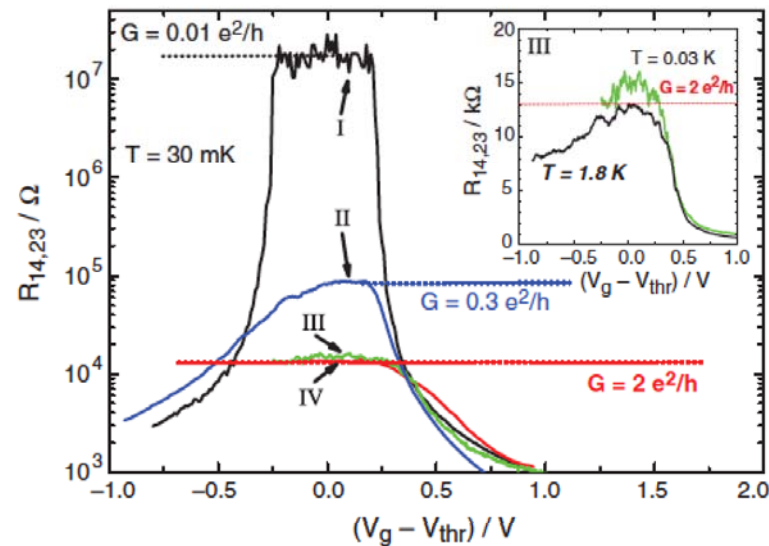
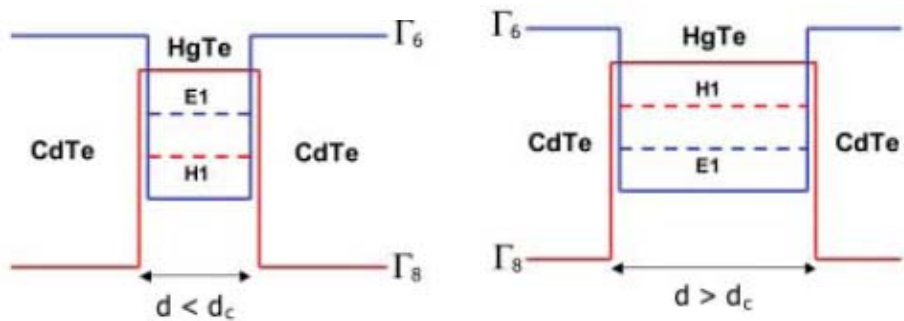
HgTe/CdTe Quantum Wells



$$H_{\text{eff}}(k_x, k_y) = \begin{pmatrix} H(k) & 0 \\ 0 & H^*(-k) \end{pmatrix},$$

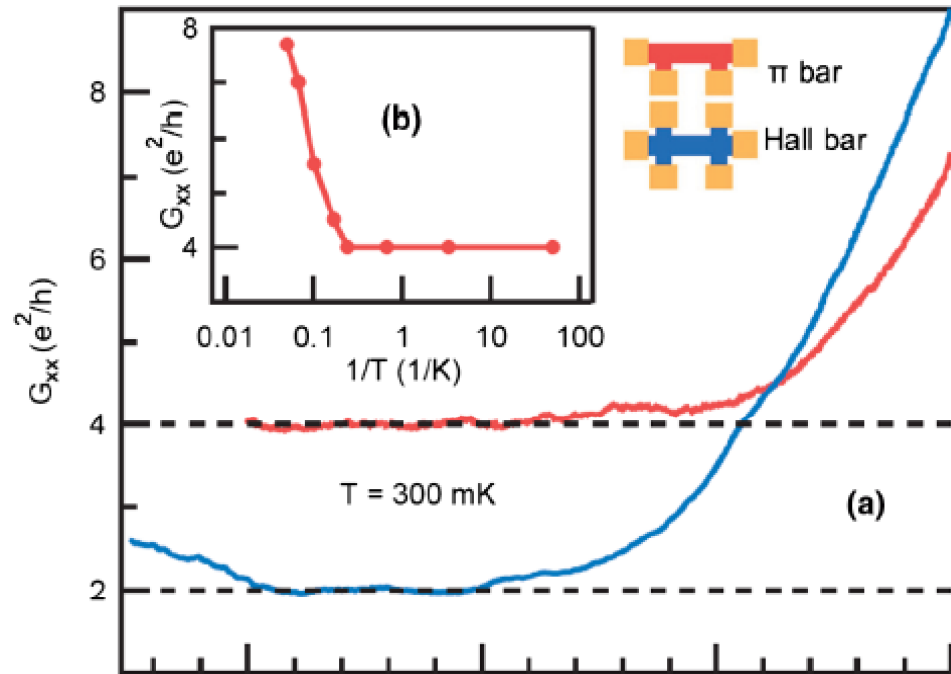
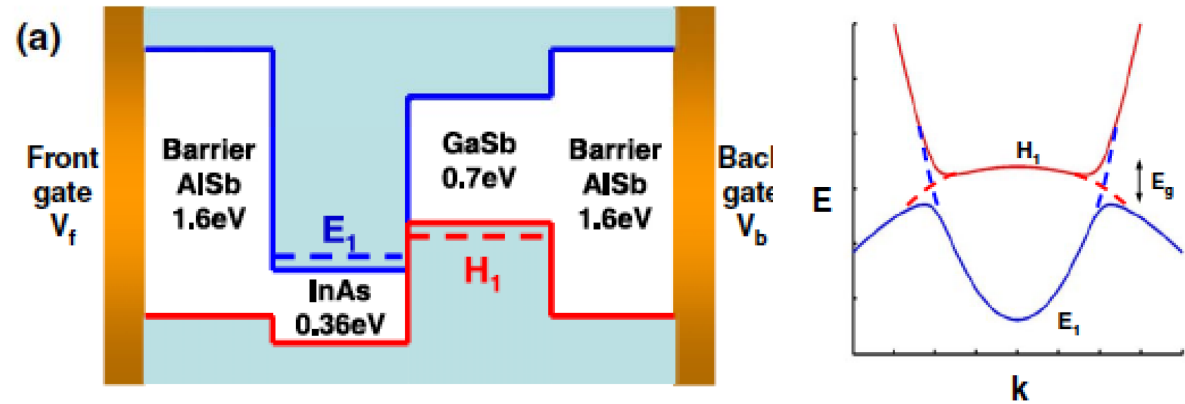
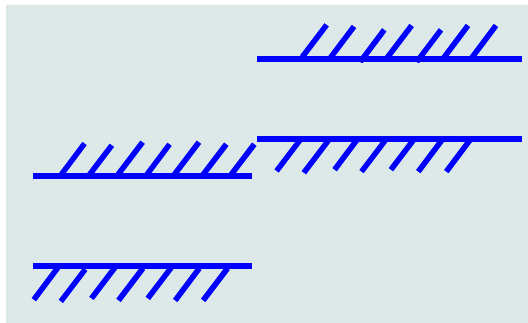
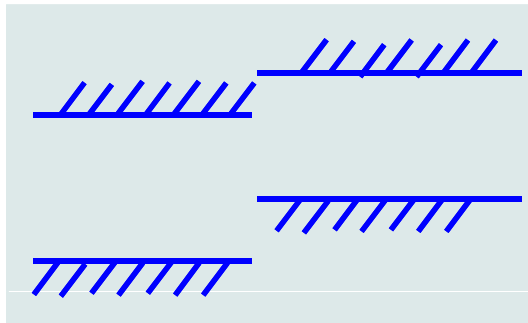
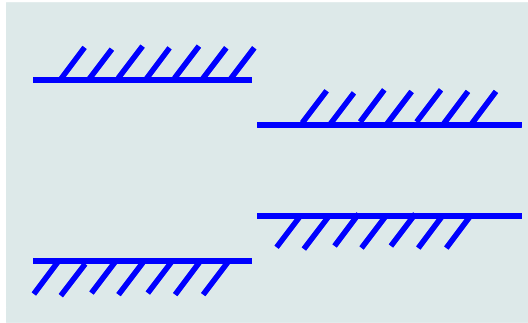
$$H(k) = \varepsilon(k) + d_i(k)\sigma_i \quad (2)$$

$|E1, m_j = 1/2\rangle, |H1, m_j = 3/2\rangle,$
 $|E1, m_j = -1/2\rangle, |H1, m_j = -3/2\rangle$



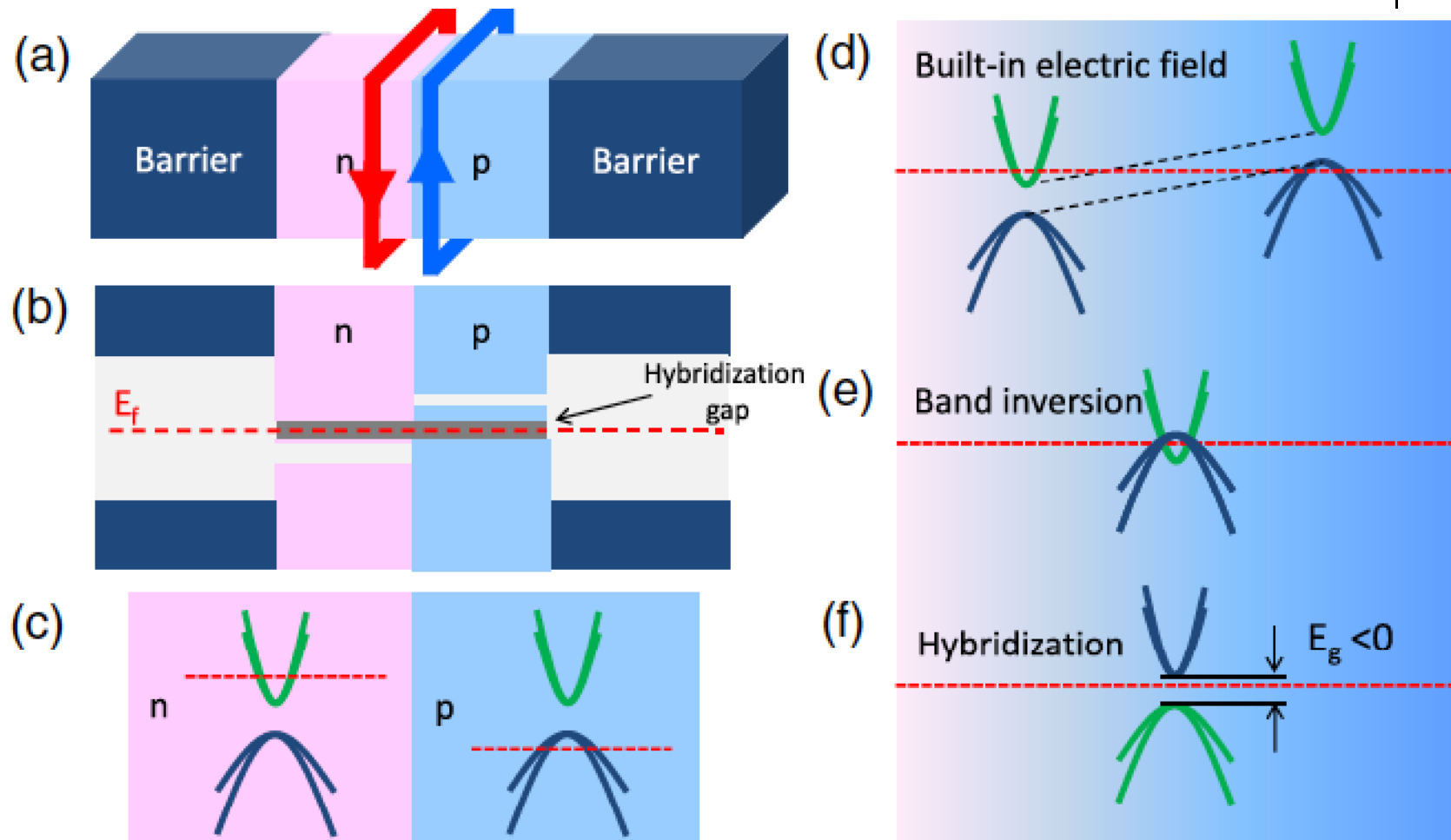
[Bernevig et al. Science 2006](#); [Markus et al. Science 2007](#)

Type-II quantum wells

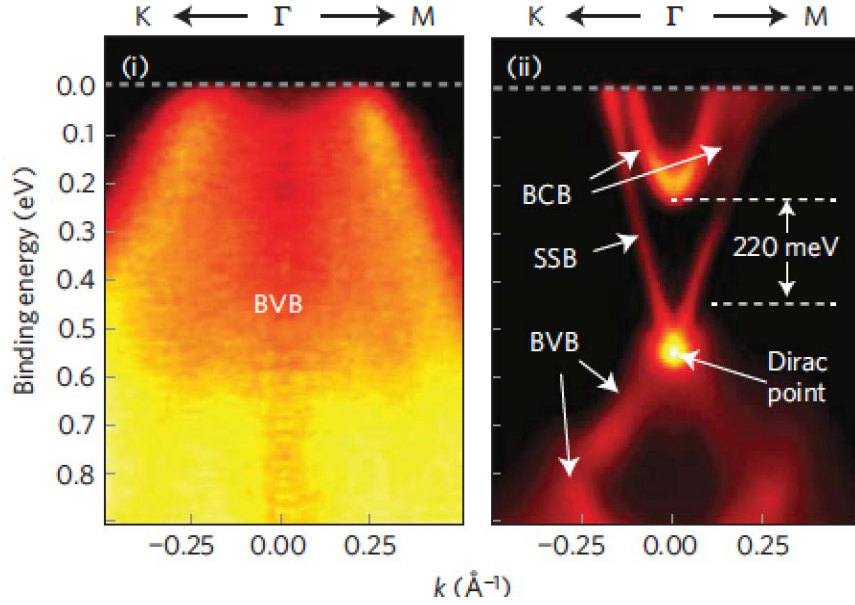
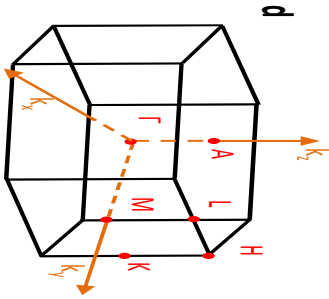
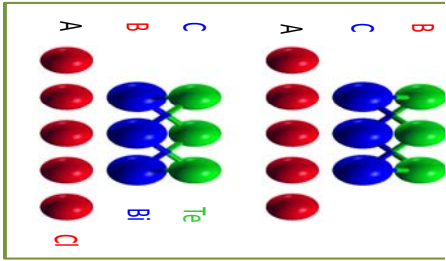


[Liu et al. PRL 2008; Du et al. PRL 2013](#)

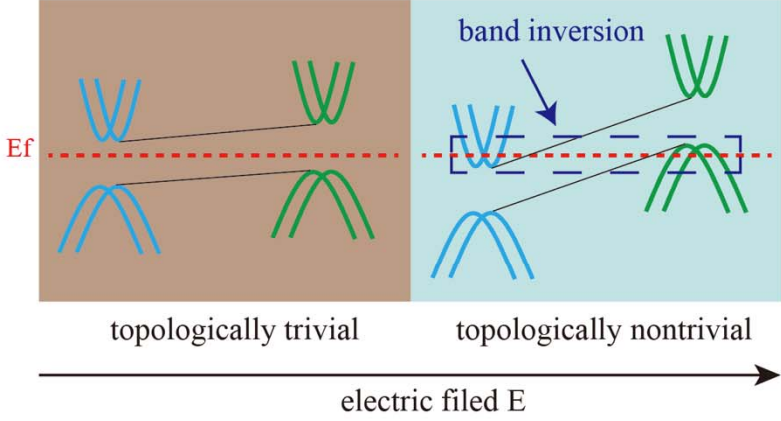
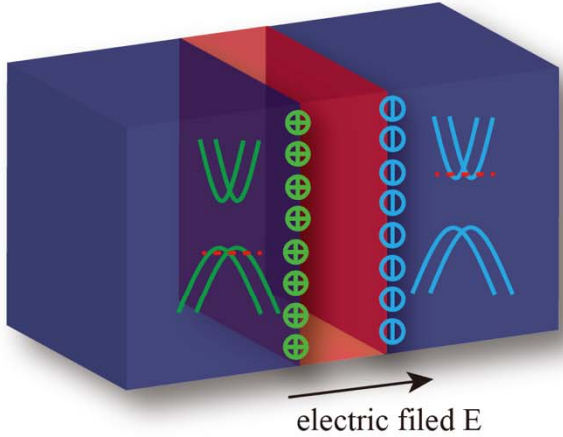
pn doping to make type-II quantum Wells



Polar quantum Wells



[Chen et al. Nature Physics 2013](#)



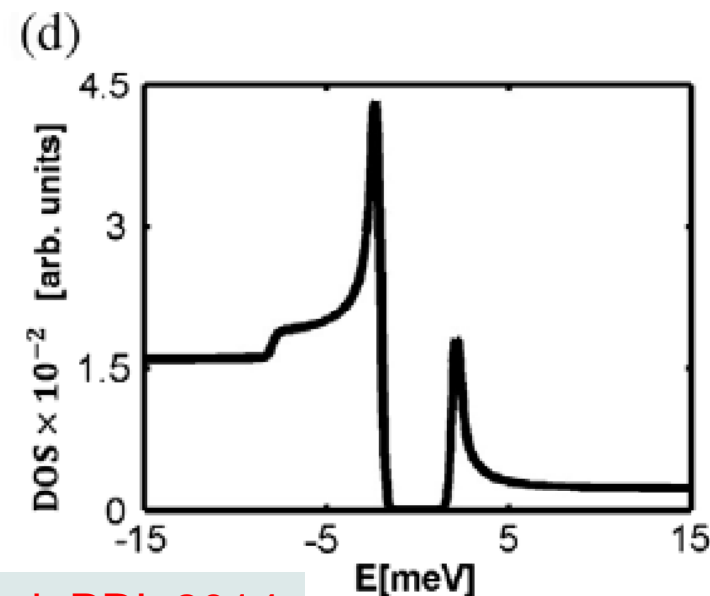
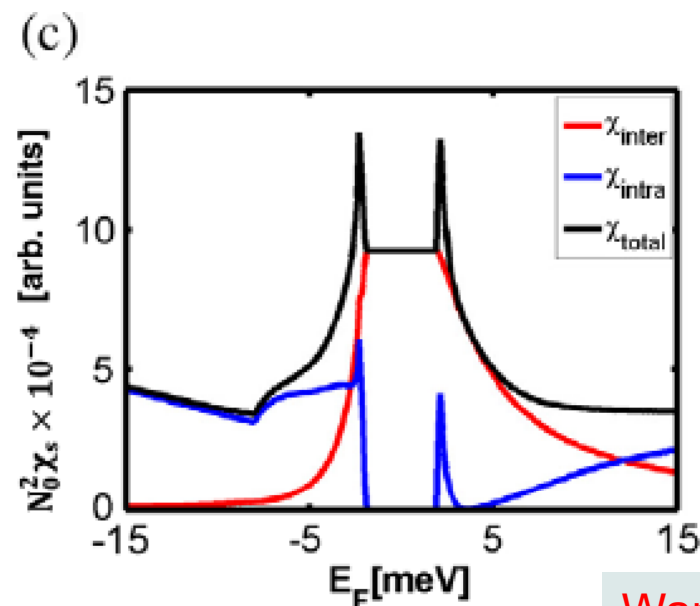
Be prepared

Magnetic doping in type-II quantum Wells



$$H = H_0 + H_{\text{BIA}} + H_{\text{SIA}} + H_{\text{ex}}. \quad H_{\text{ex}} = \sum_{\vec{R}_n} \mathbf{S}_M(\vec{R}_n) \cdot \tilde{\mathbf{s}},$$

$$\tilde{\chi}_s = \lim_{q \rightarrow 0} \text{Re} \left[\sum_{i,j,\sigma,\sigma',\vec{k}} \frac{|\langle u_{i\sigma,\vec{k}} | \tilde{\mathbf{s}} | u_{j\sigma',\vec{k}+\vec{q}} \rangle|^2 [f_{i\sigma}(\vec{k}) - f_{j\sigma'}(\vec{k}+\vec{q})]}{E_{j\sigma'}(\vec{k}+\vec{q}) - E_{i\sigma}(\vec{k}) + i\Gamma} \right]$$

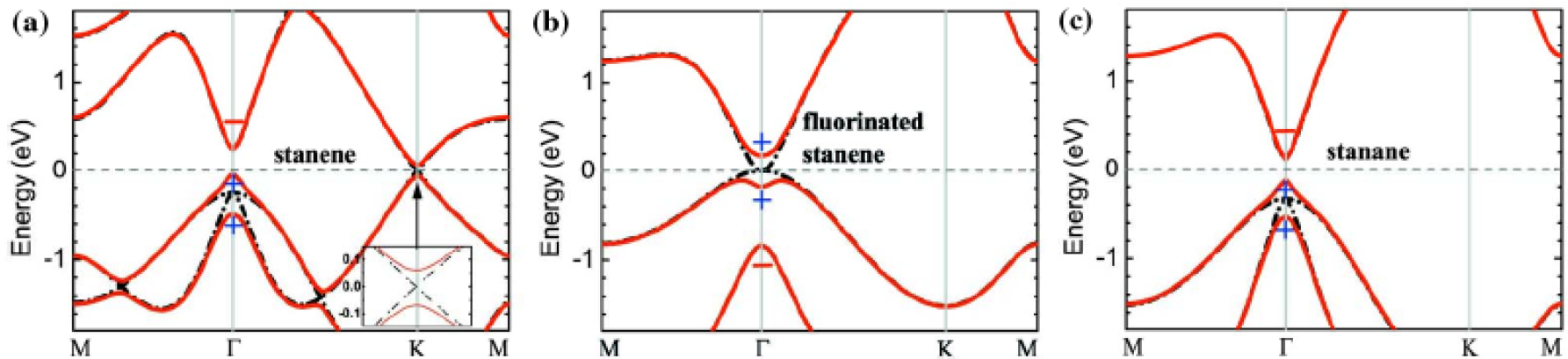
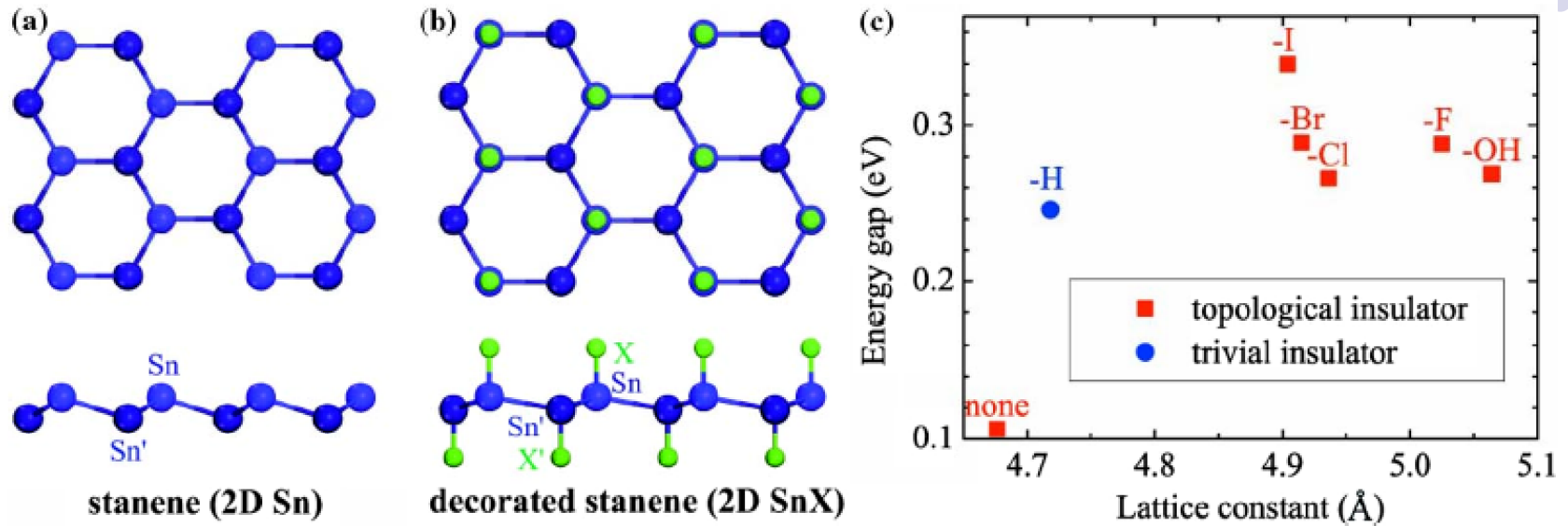


[Wang et al. PRL 2014](#)

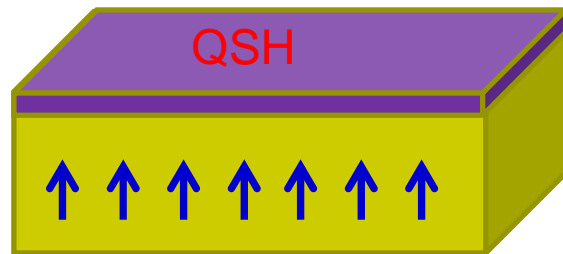


2D topological insulators in Graphene type compounds

Graphene-type QSH states

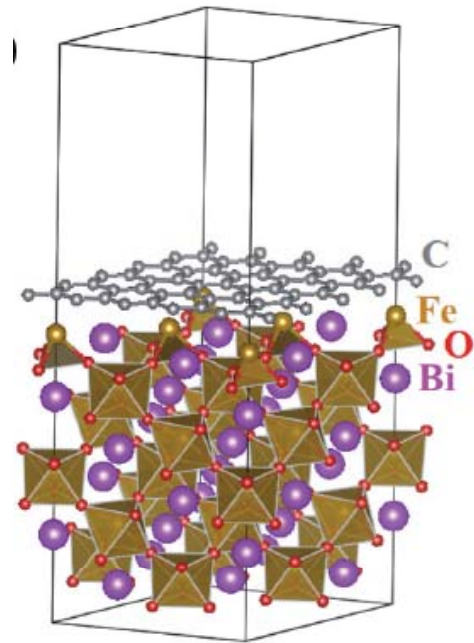


Magnetic proximity in Graphene-type QSH states



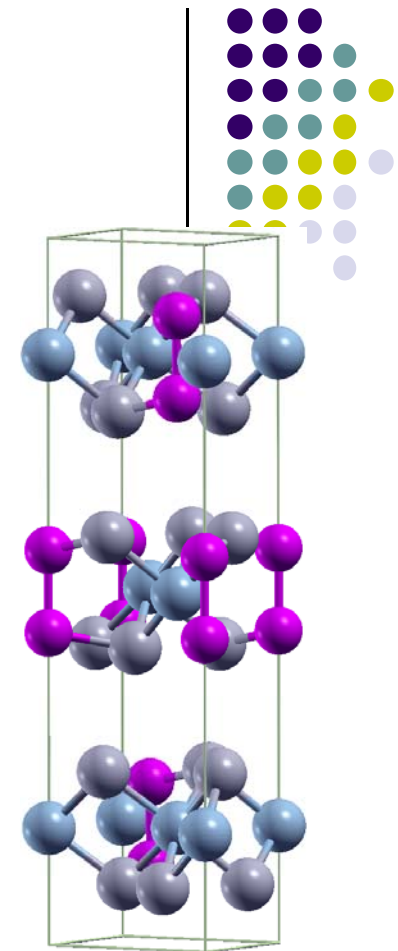
FM insulators
Or
AFM insulators

Difficulty:
1) Out of plane magnetic moment at the interface
2) Strong magnetic proximity
3) Charge transfer



BiFeO₃
Multiferroic materials
G-type antiferromagnet
T_N=653K

Qiao et al. PRL 112, 116404 (2014)



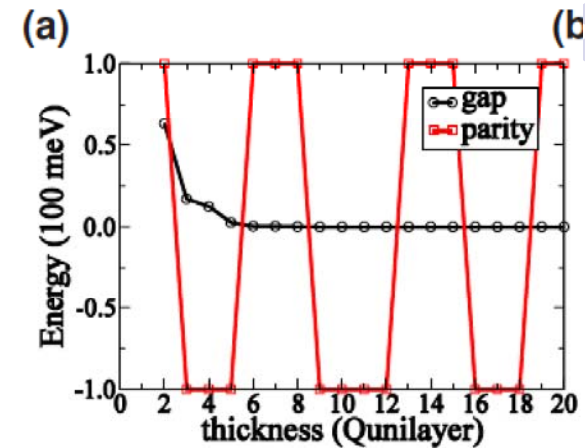
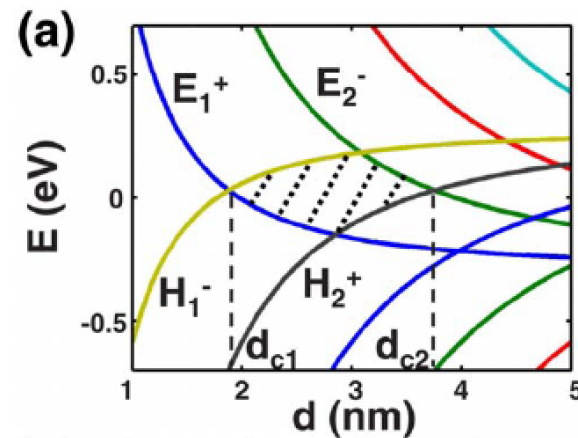
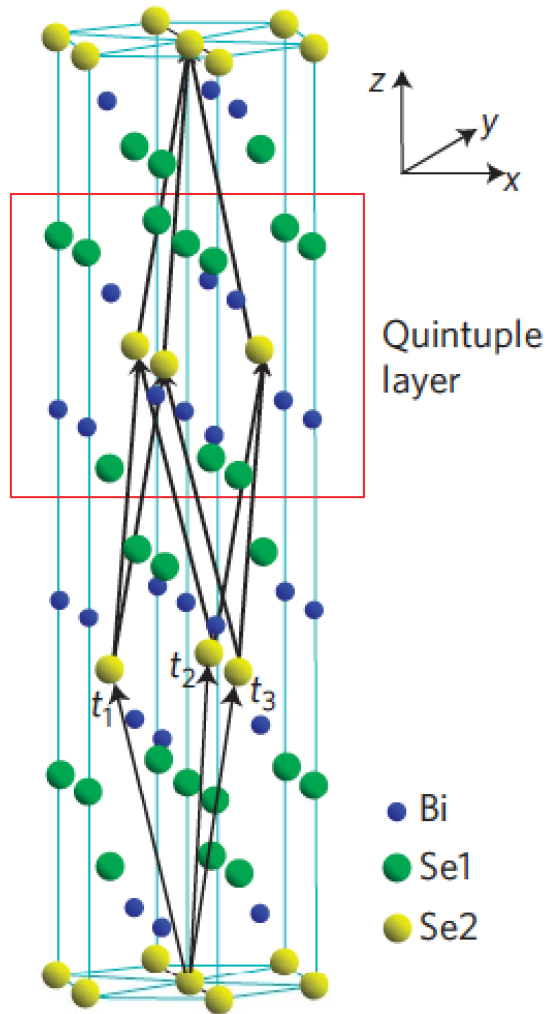
CrGeTe₃
Ferromagnetic Insulator.
Out of plane.

Be prepared

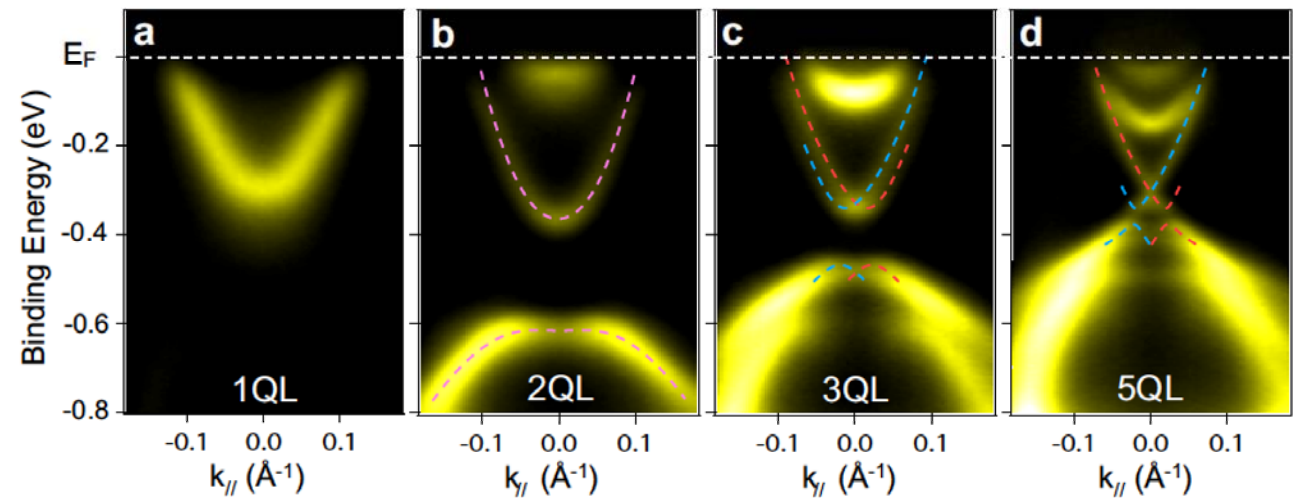


2D topological insulators from the 2D limit of 3D topological insulators

From 3D topological insulators



Liu et al. PRB 2010



Zhang et al. Nature Physics 2010

Magnetic impurity doping in QSH states

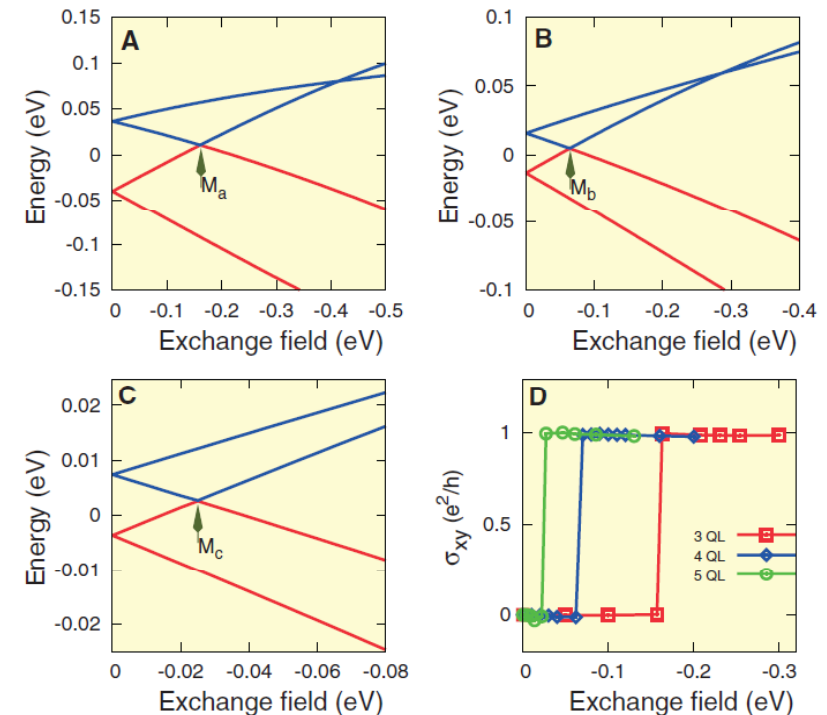
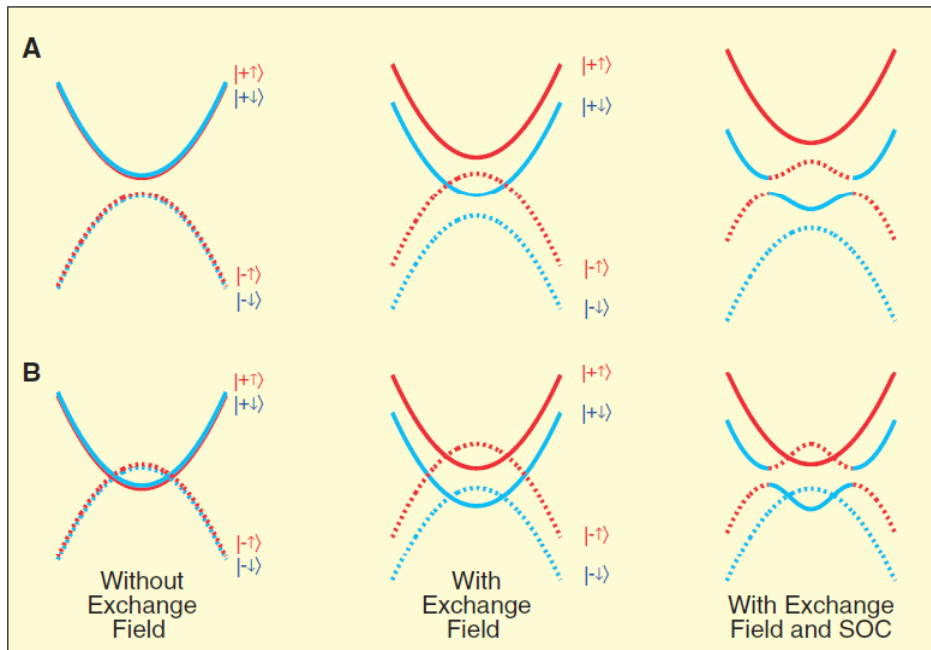


$$H = H_{sf} + H_{Zeeman}$$

$$= \begin{pmatrix} m_k + gM & iv_Fk_- & 0 & 0 \\ -iv_Fk_- & -m_k - gM & 0 & 0 \\ 0 & 0 & m_k - gM & -iv_Fk_- \\ 0 & 0 & iv_Fk_- & -m_k + gM \end{pmatrix}$$

m_k a finite mass
 g the effective factor
 M the exchange field along z direction

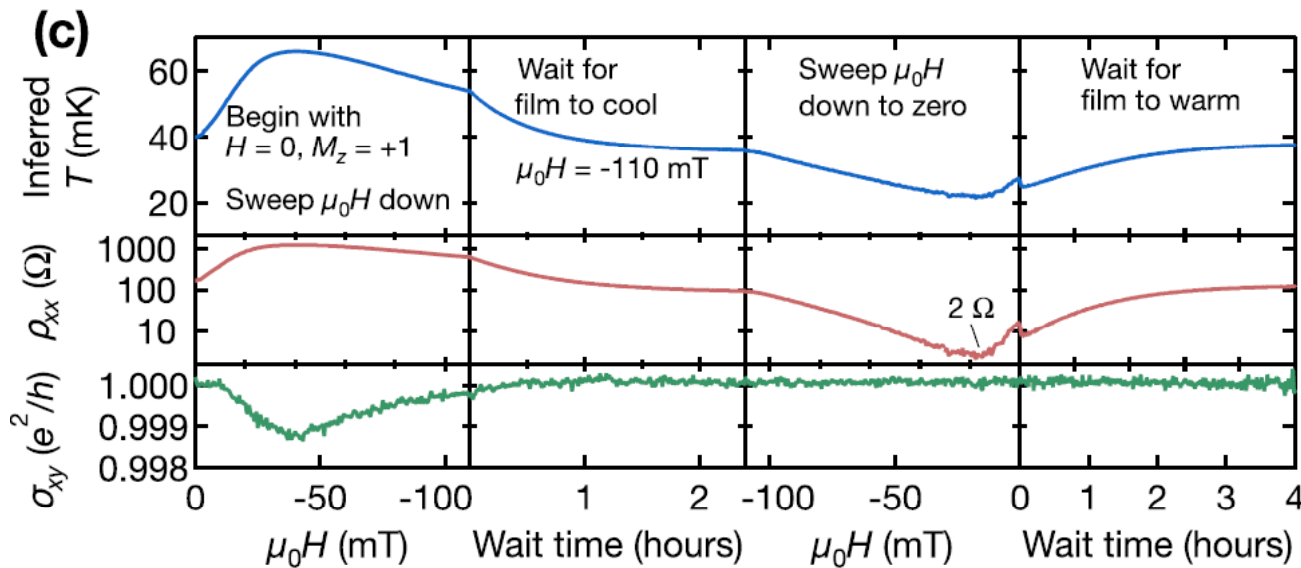
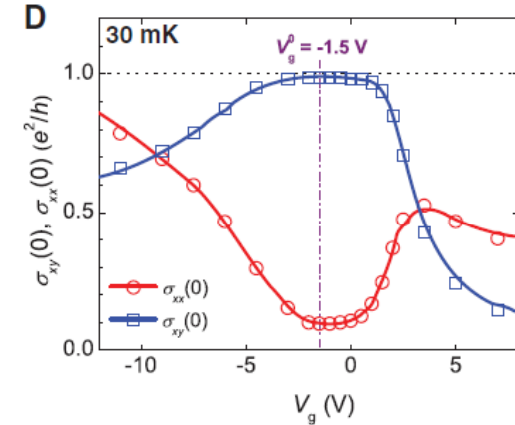
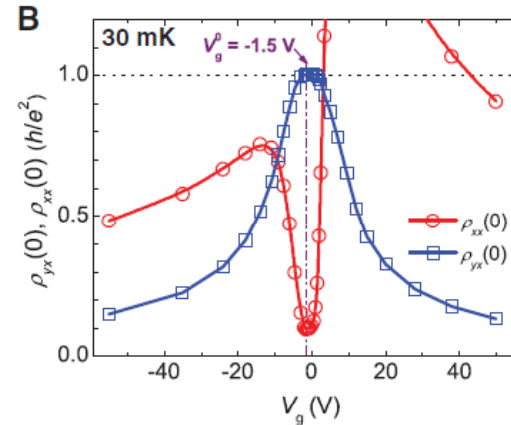
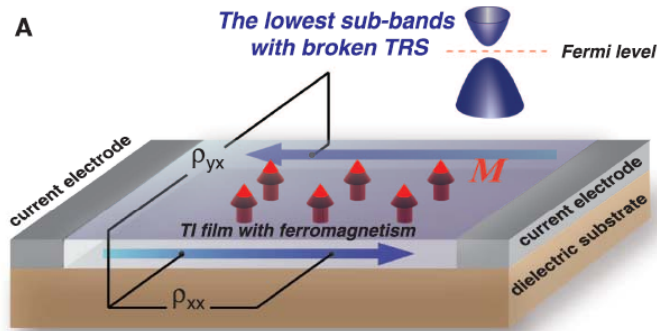
Yu et al. Science 2010
Liu et al. PRL 101, 146802 (2008)
Zhang et al. PRL 112, 216803 (2014)



Really observed QAH Effect in experiments



30mK



$$\sigma_{xy} \sim 1.0 e^2/h; \rho_{xx} \sim 2\Omega$$

Chang et al. Science 2013
Kou et al. PRL 2014
Checkelsky Nat. Phys. 2014

Further optimize the Position of the chemical position for 0 ρ_{xx} ?

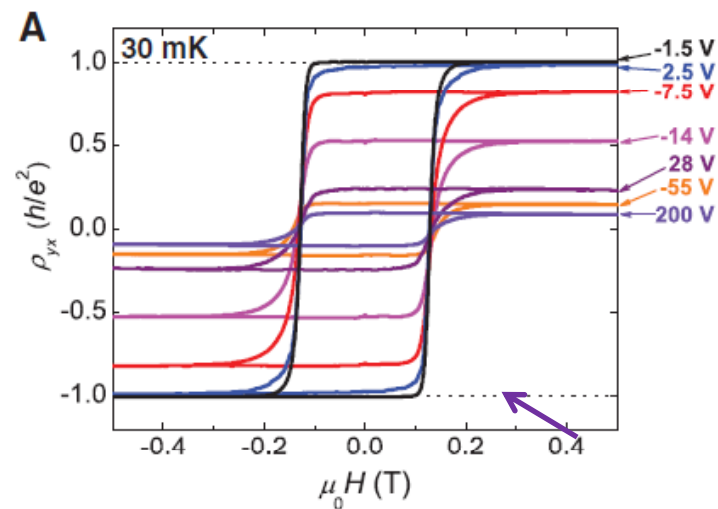
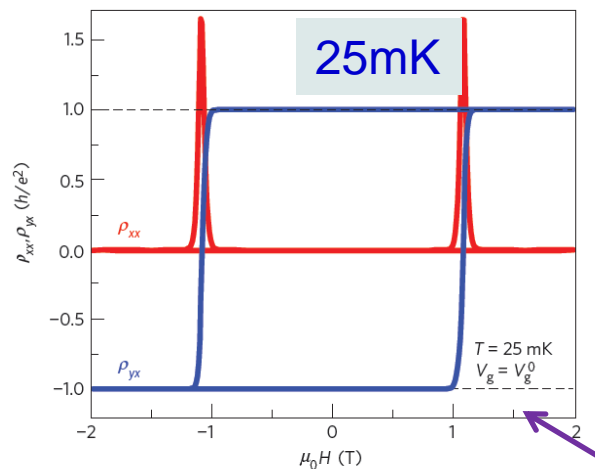
Bestwick et al
PRL 114, 187201 (2015)

Precise quantization of the QAH effect—V-doped case

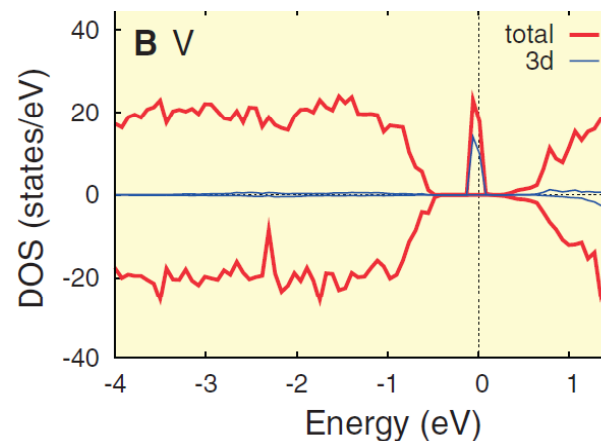
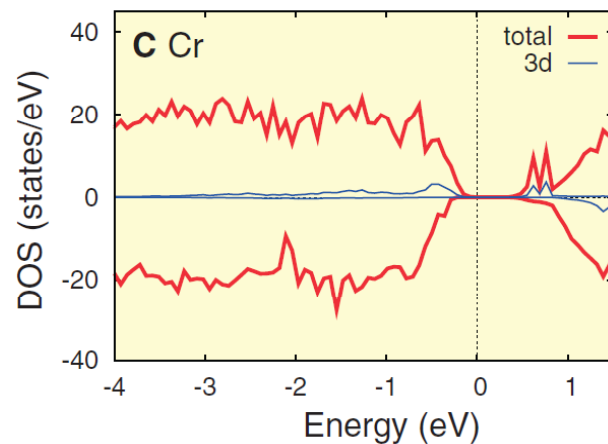


4 QL $V_{0.11}(\text{Bi}_{0.29}\text{Sb}_{0.71})_2\text{Te}_3$

$$\sigma_{xy} \sim 0.9998e^2/h; \rho_{xx} \sim 3.3\Omega$$



- 1) Self-driven, without magnetic training
- 2) A large coercive field H_c



- 1) Impurity bands and the QAH effect
- 2) Exchange mechanism

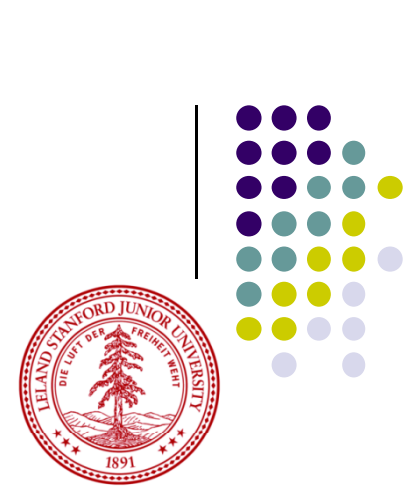
Chang et al. Nat. Materials 2015



Outlook

- How to understand the ‘surprised’ observation?
- How to improve the working temperature to K or ten K level?
- Can we realize the QAH effect in much cheaper systems?
- Superconductor + the QAH effect?

Acknowledgement



✓ Yong Xu, Gang Xu, Jing Wang, Shou-Cheng Zhang

✓ Chaoxing Liu



✓ Cui-zu Chang



✓ Hong Yao, Zhong Wang

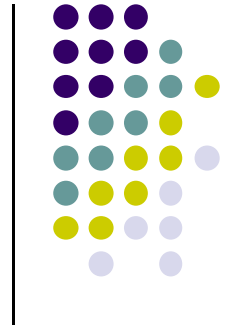




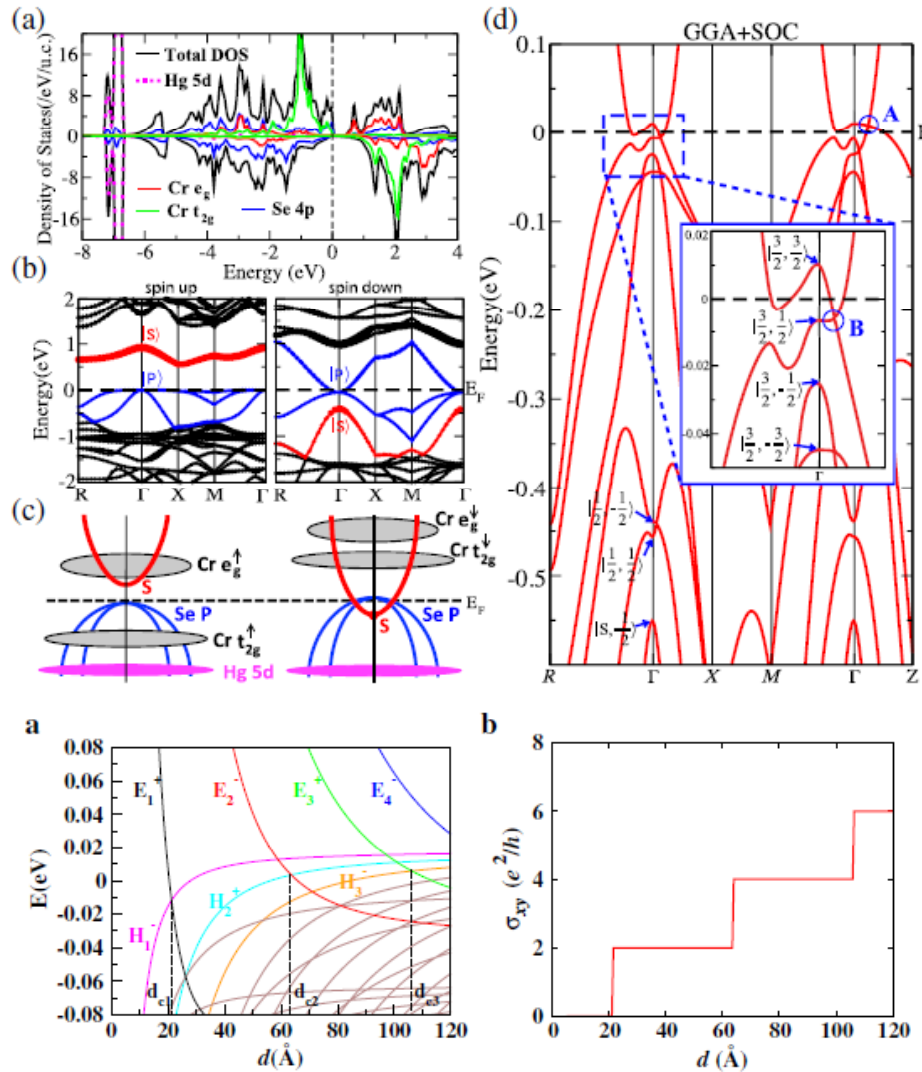
Summary

- Background
- Design 2D topological insulators
 - Type-II Quantum Wells, polar Quantum Wells
 - Graphene, Silience, Stanene et al.
 - 2D limit from 3D topological insulators
- Outlook

Thank you for your attention!

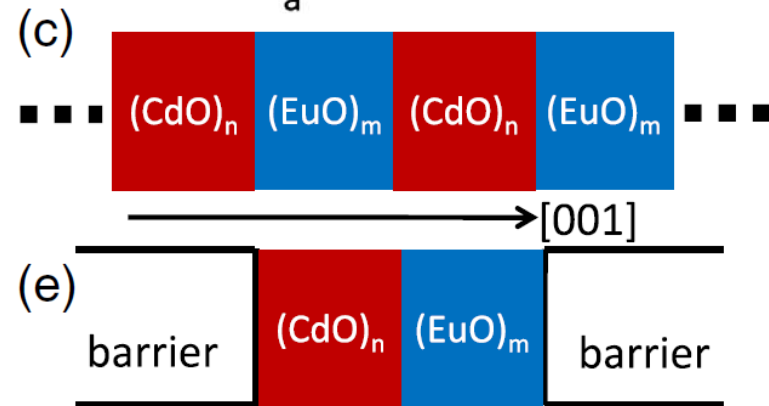
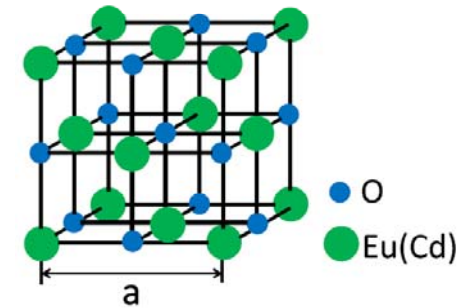


Intrinsic magnetic insulators



Xu et al. PRL 107, 186806 (2011)

Intrinsic FM insulators
+
Type-II Quantum Wells structure



Zhang et al. PRL 112, 096804 (2014)

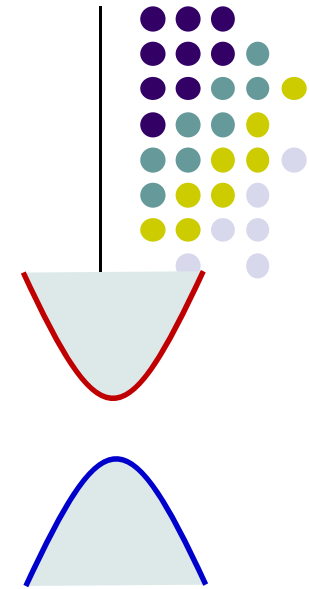
The toy model of the QAH effect

$$H = \sum_{\vec{k}} H(\vec{k})$$

$$H(\vec{k}) = \epsilon(\vec{k}) + V d_\alpha(\vec{k}) \sigma^\alpha$$

$$E_{\pm}(\vec{k}) = \epsilon(\vec{k}) \pm V \sqrt{\sum_{\alpha} d_{\alpha}^2(\vec{k})}$$

$$\min_{k \in BZ} E_{+}(\vec{k}) > \max_{k \in BZ} E_{-}(\vec{k})$$

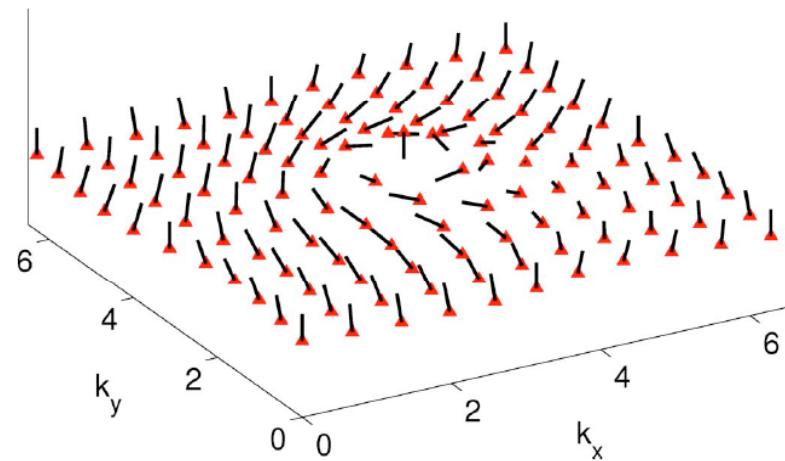


The Hall conductivity:

$$\begin{aligned} \sigma_{xy} &= -\frac{1}{8\pi^2} \int \int_{FBZ} dk_x dk_y \hat{d} \cdot \partial_x \hat{d} \times \partial_y \hat{d} \\ &= -\frac{n}{2\pi} \end{aligned}$$

with

$$\hat{d}_\alpha(\vec{k}) = d_\alpha(\vec{k}) / \sqrt{\sum_{\alpha} d_{\alpha}^2(\vec{k})}$$



Quantized conductivity without the Landau levels!