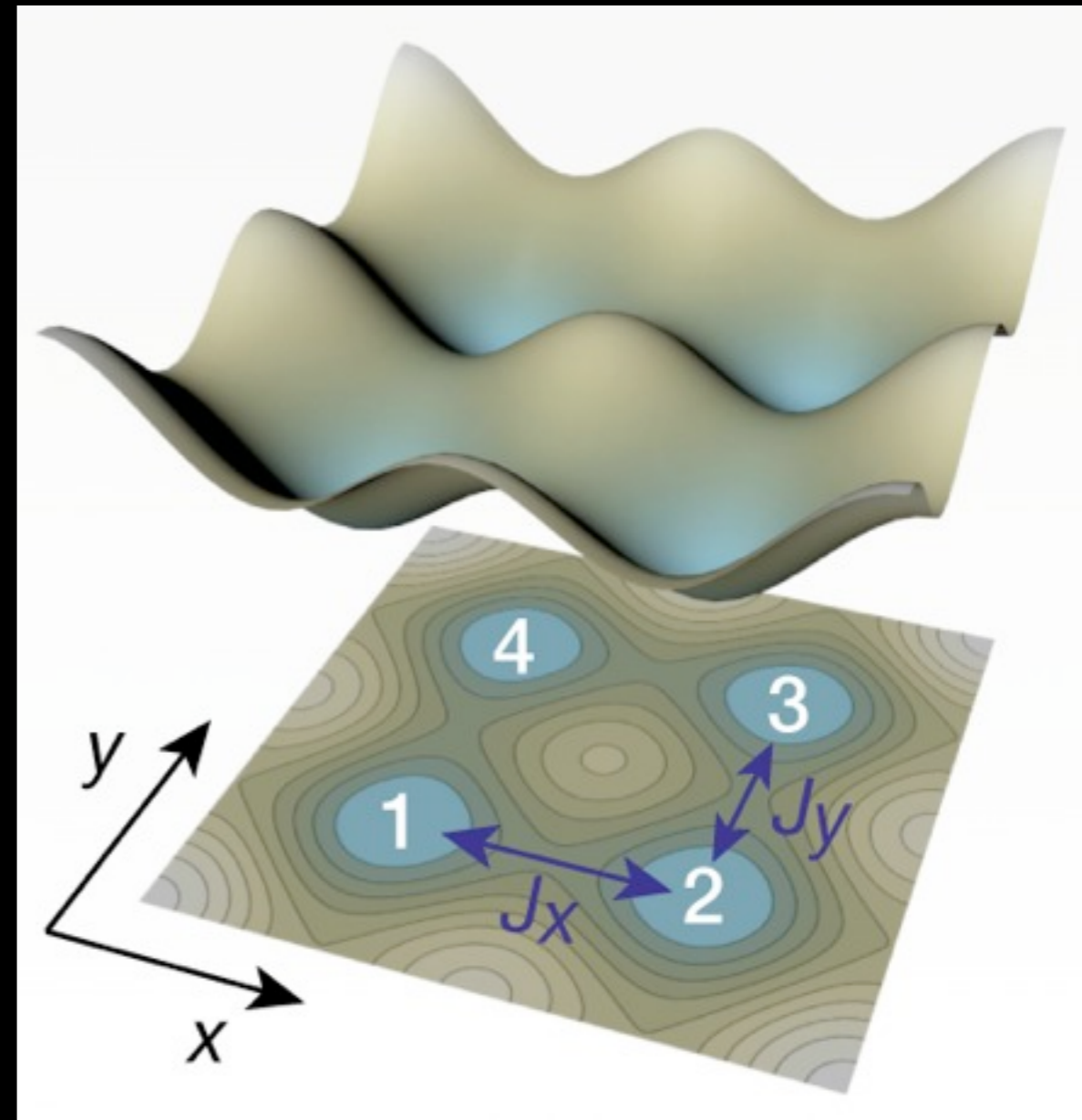


QUANTUM SIMULATION IN OPTICAL SUPERLATTICE

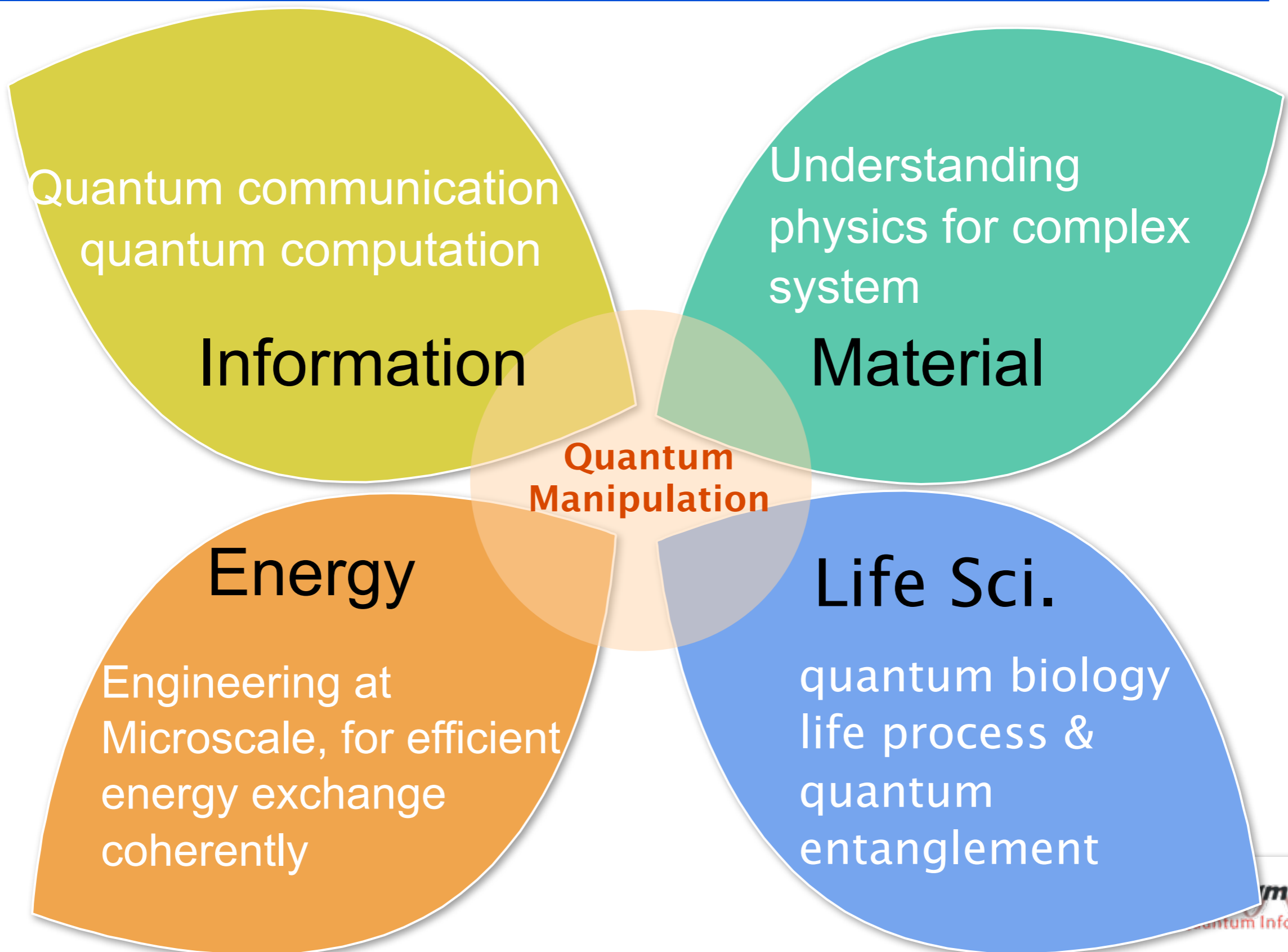
Yu-Ao Chen

Center for Quantum Engineering
Shanghai Division of Quantum Physics and Quantum Information,
National Lab for Physical Sciences at the Microscale,
University of Science and Technology of China





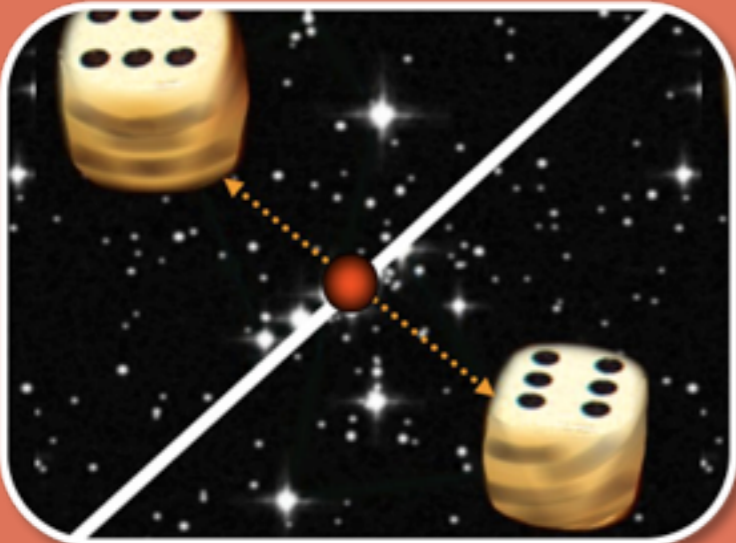
National Laboratory for Physical Sciences at Microscale



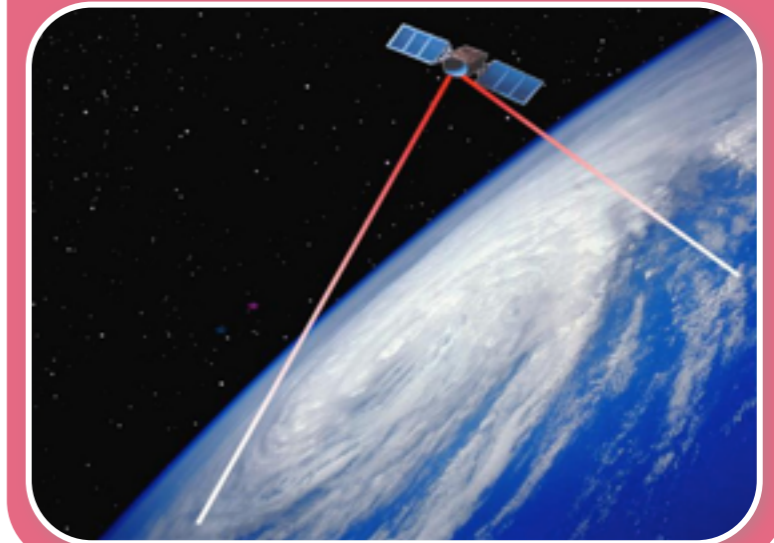


Division of Quantum Physics & Quantum Information

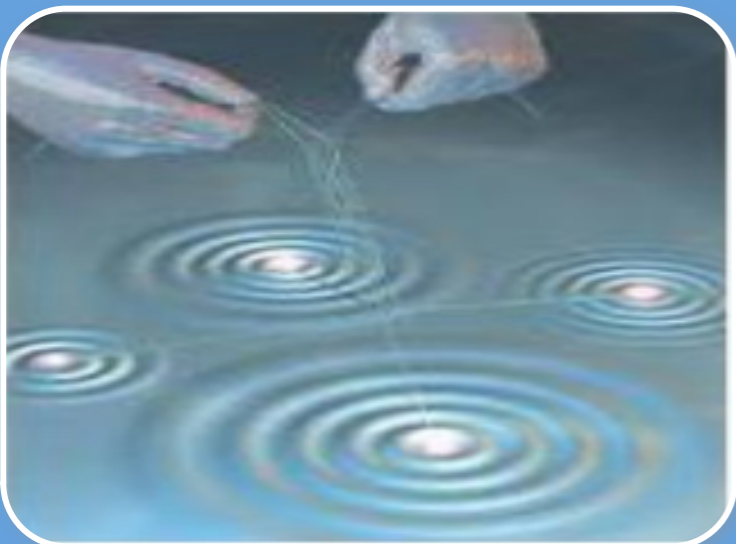
Fundermantal test of quantum machanics



Quantum Communication



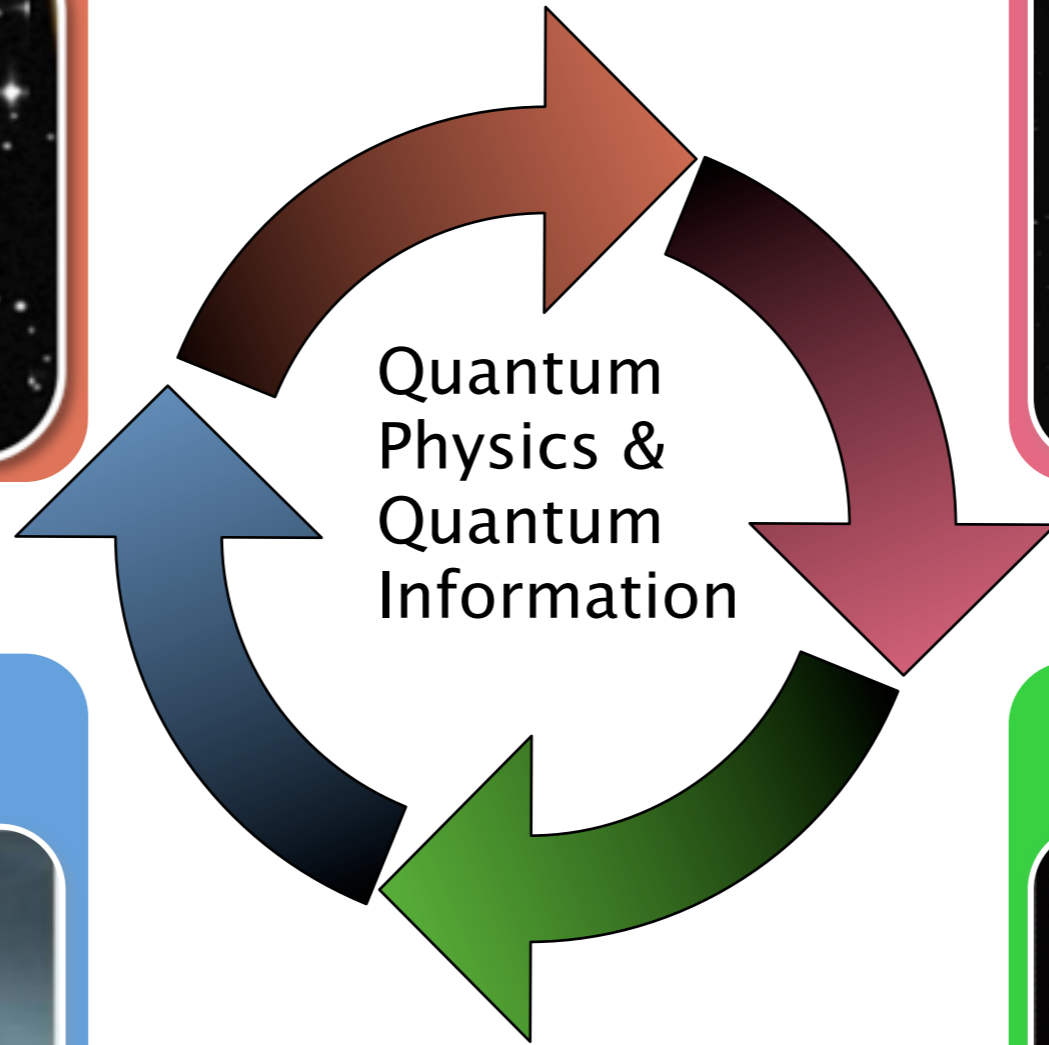
Quantum Simulation



Quantum Computation



Quantum Physics & Quantum Information









200 ft
50 m



K. Chen
陈凯



C.-Y. Lu
陆朝阳



S. Chen
陈帅



Z.-S. Yuan
苑震生



Z.-B. Chen
陈增兵



B. Zhao
赵博



J.-W. Pan
潘建伟



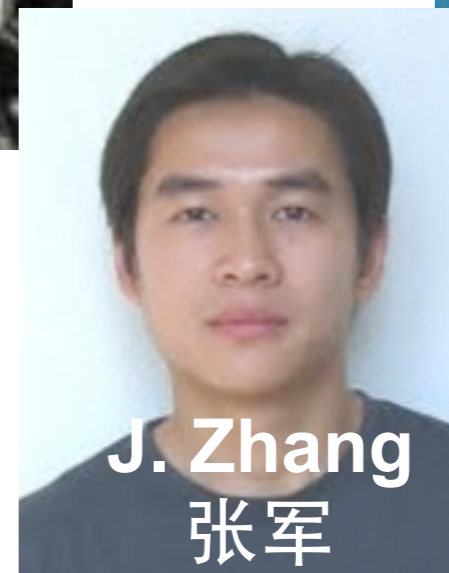
Y.-J. Deng
邓友金



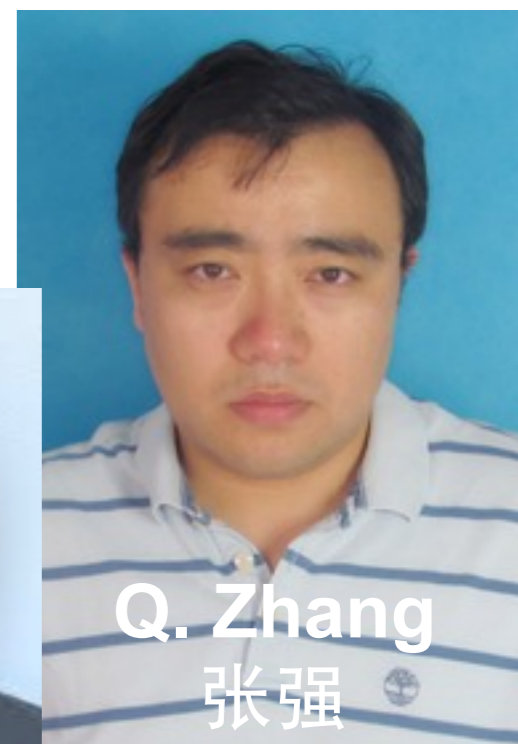
X.-H. Bao
包小辉



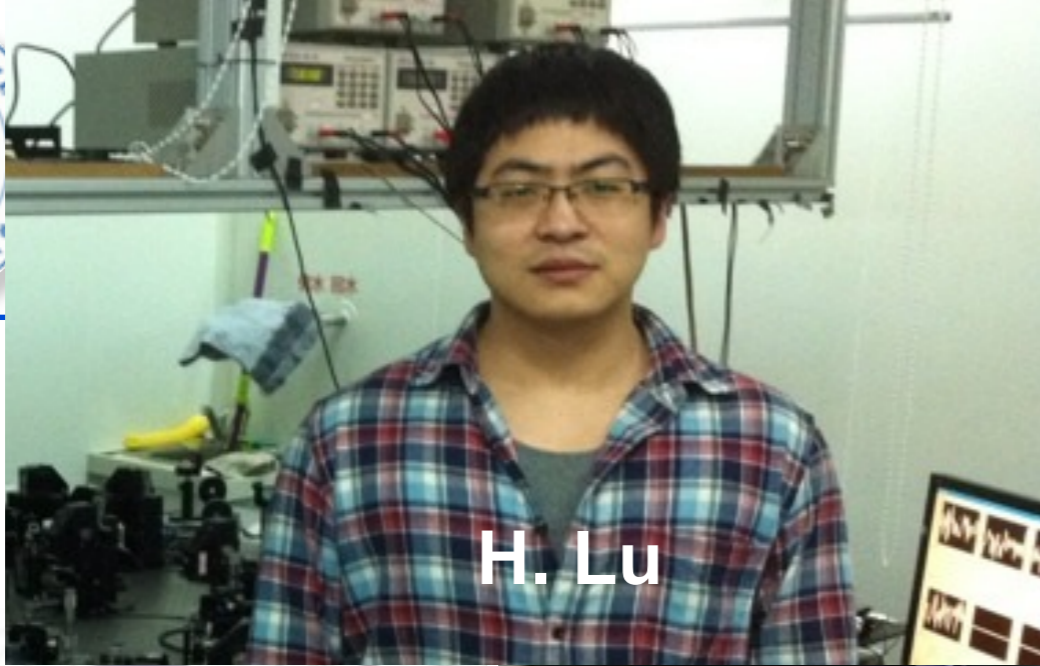
C.-Z. Peng
彭承志



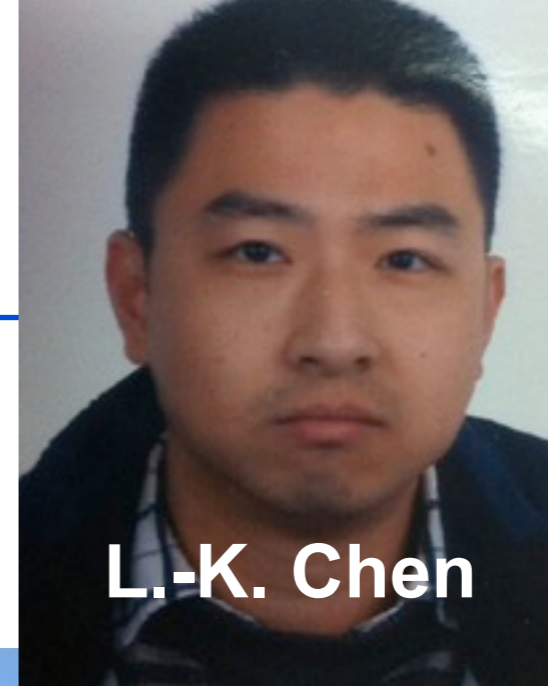
J. Zhang
张军



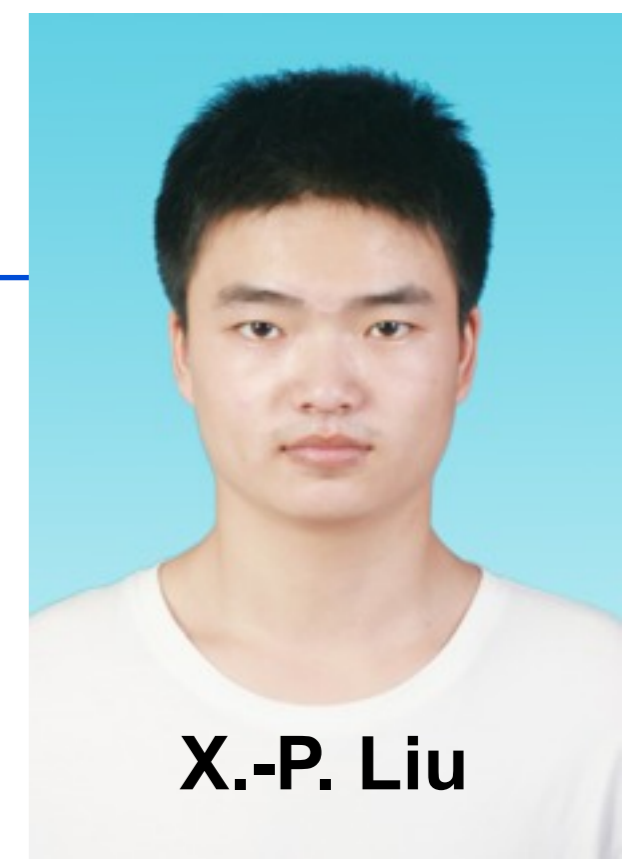
Q. Zhang
张强



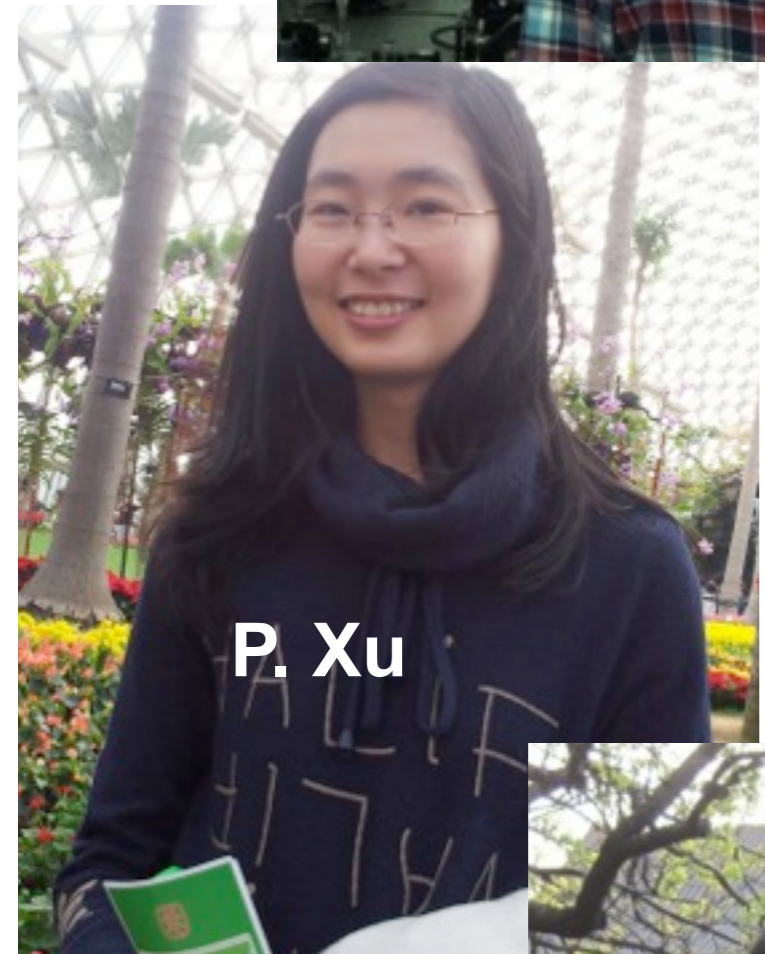
H. Lu



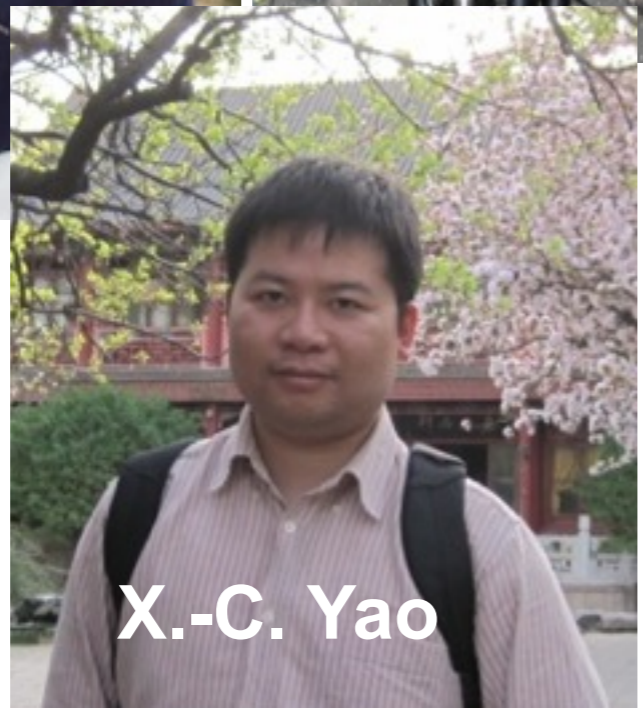
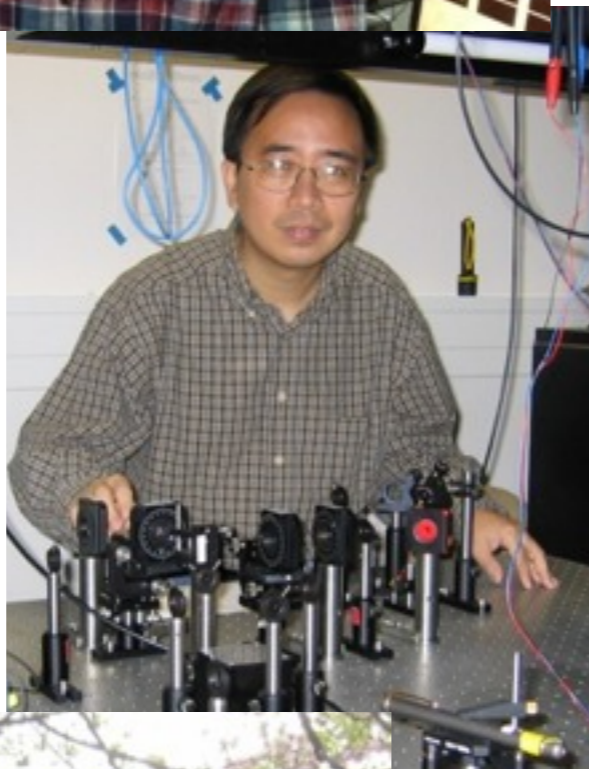
L.-K. Chen



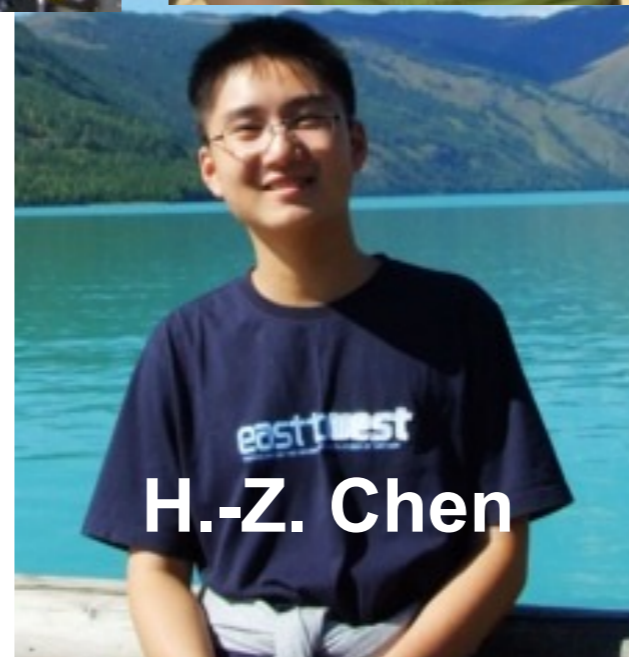
X.-P. Liu



P. Xu



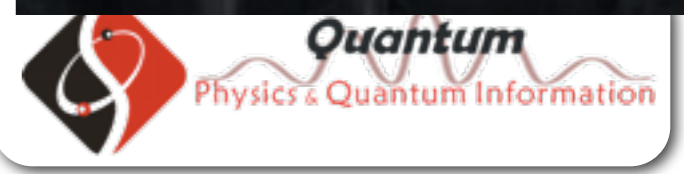
X.-C. Yao

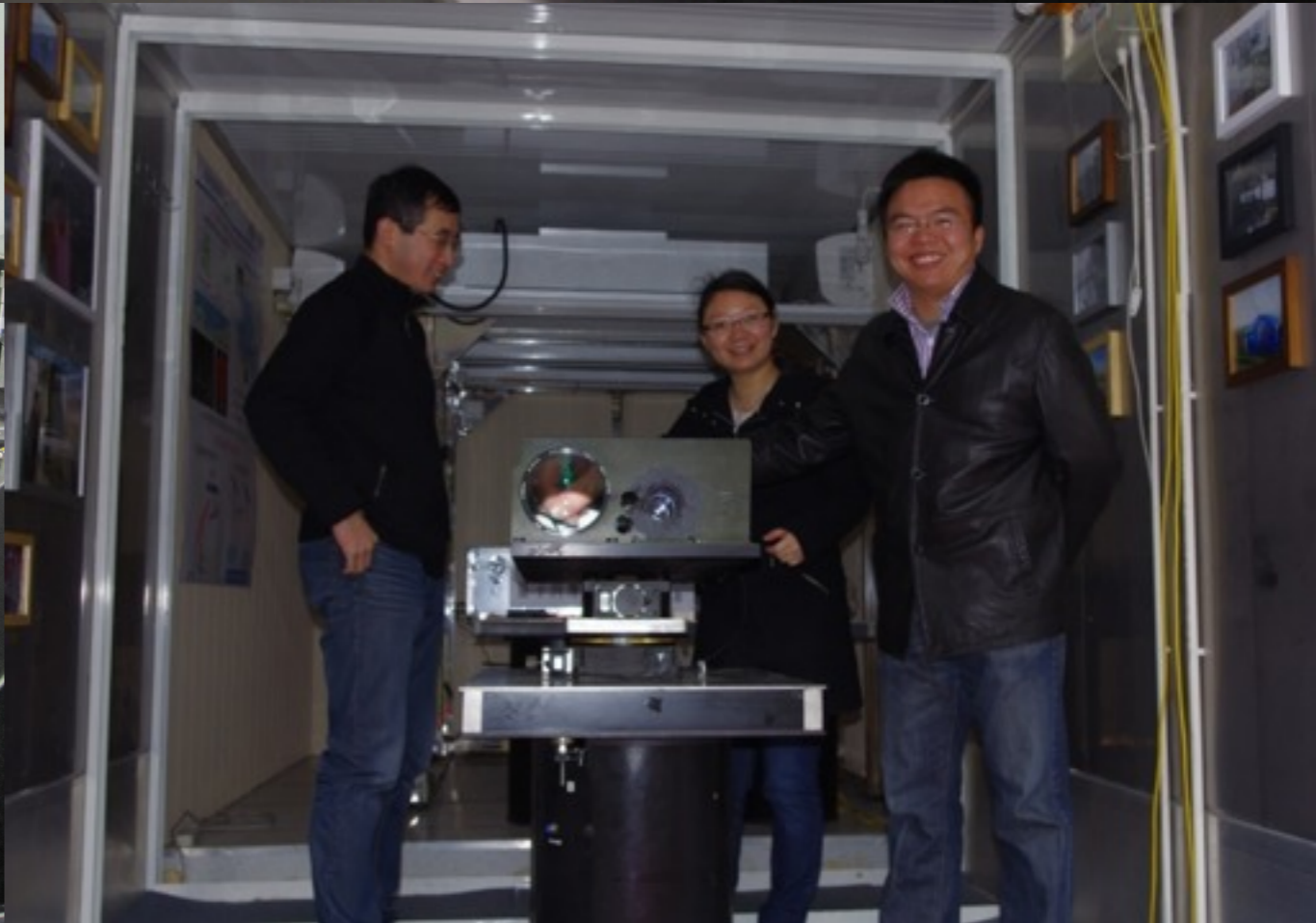
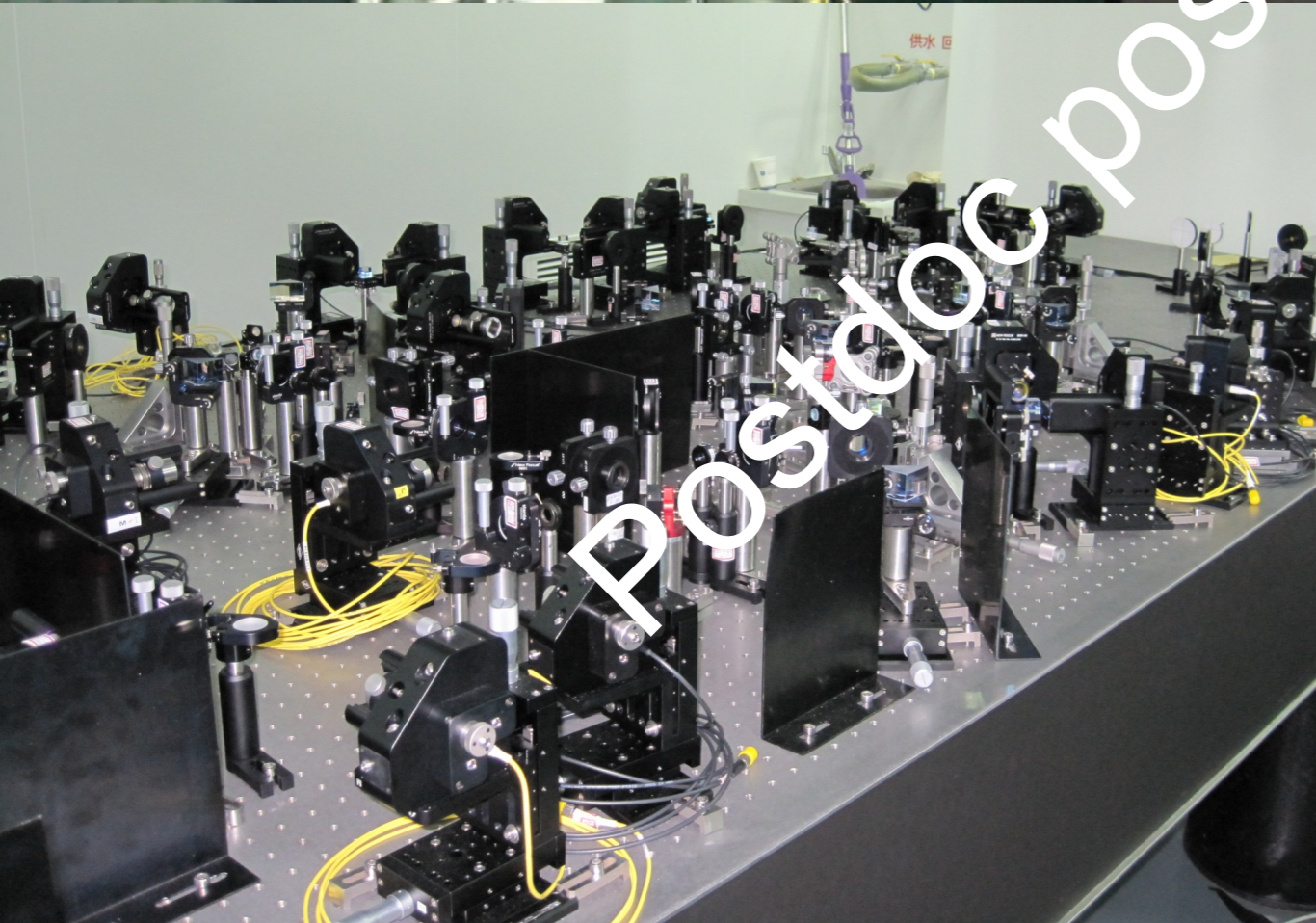
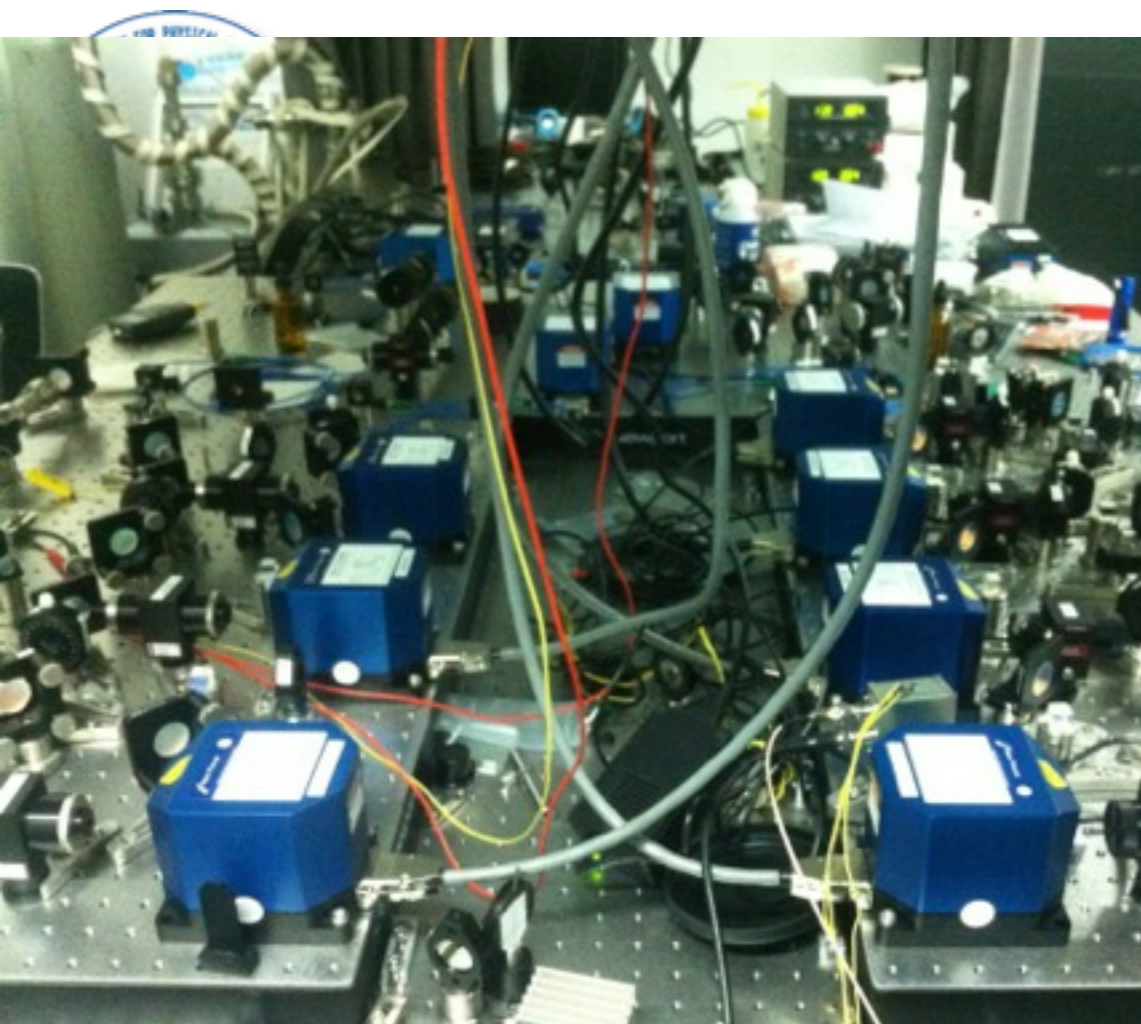


H.-Z. Chen



Y.-P. Wu

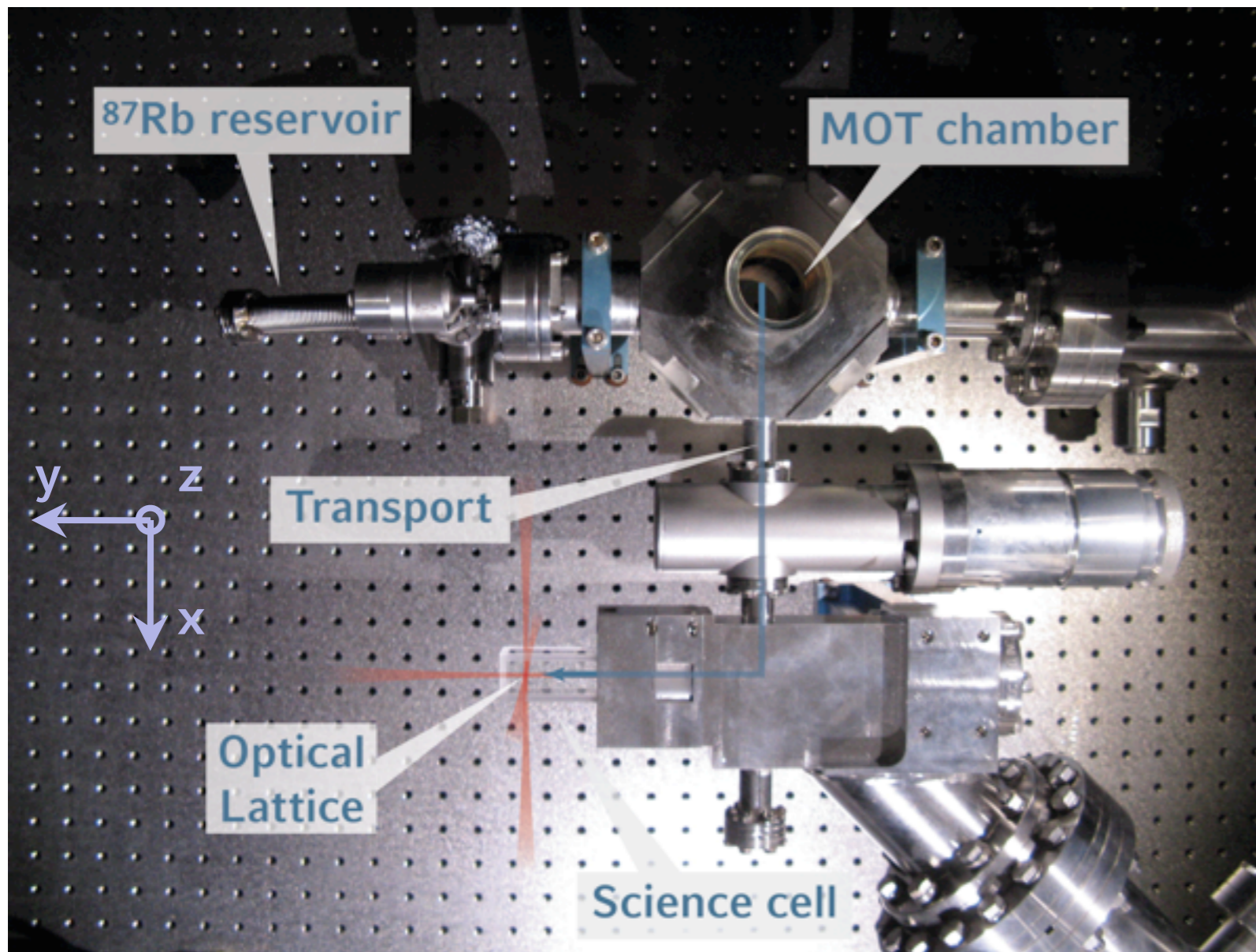




Postdoc positions available

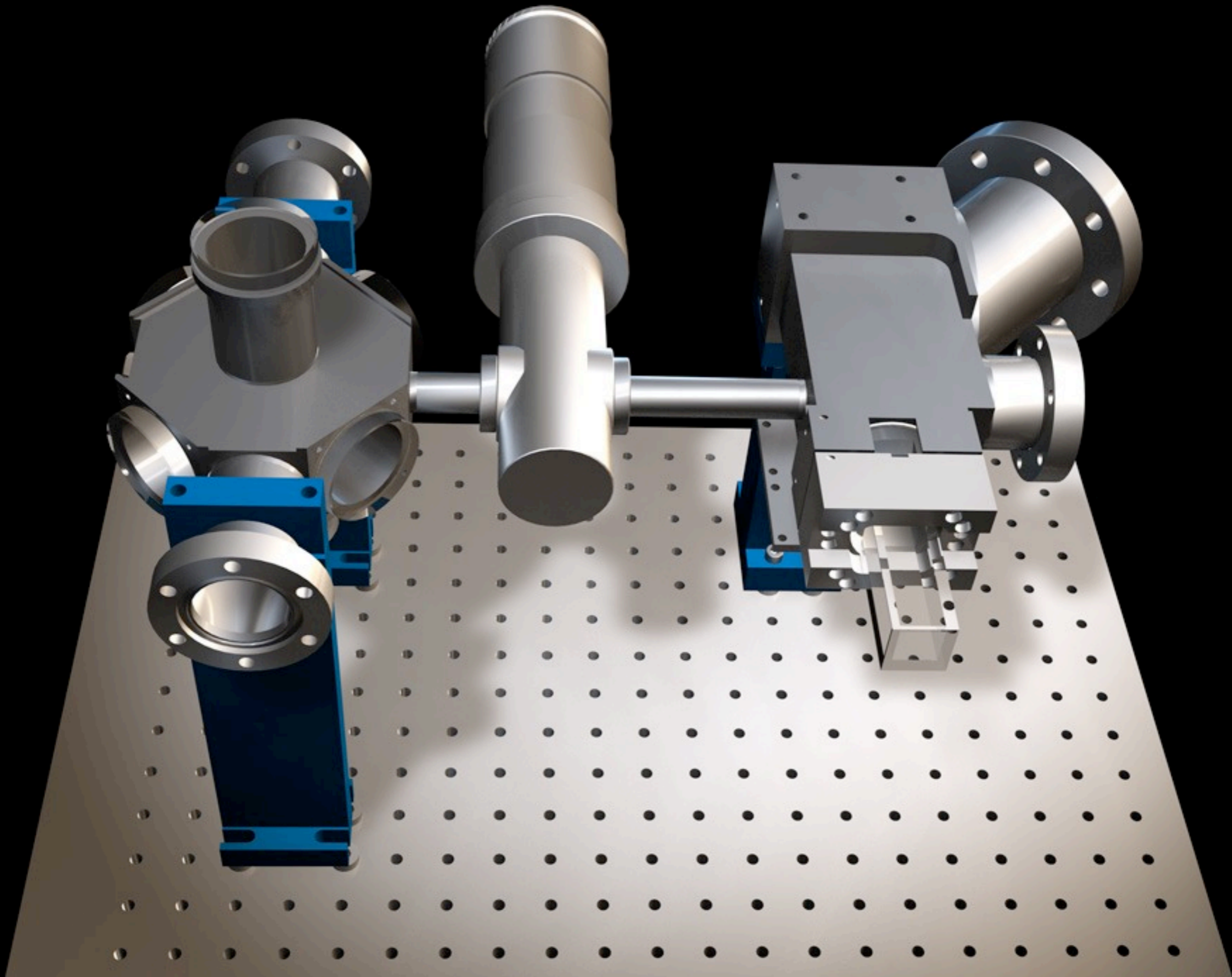


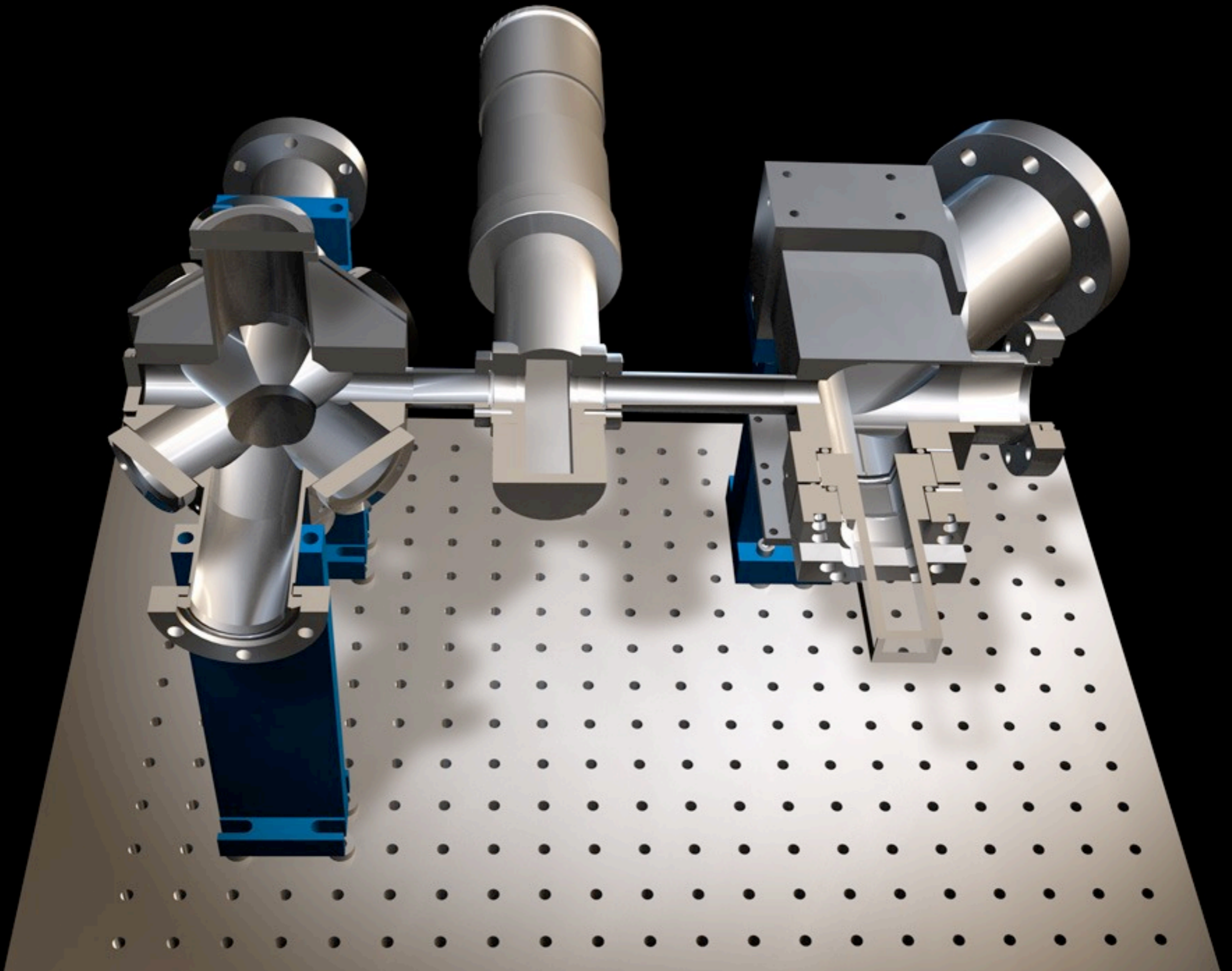
The Boson Experiment in Munich

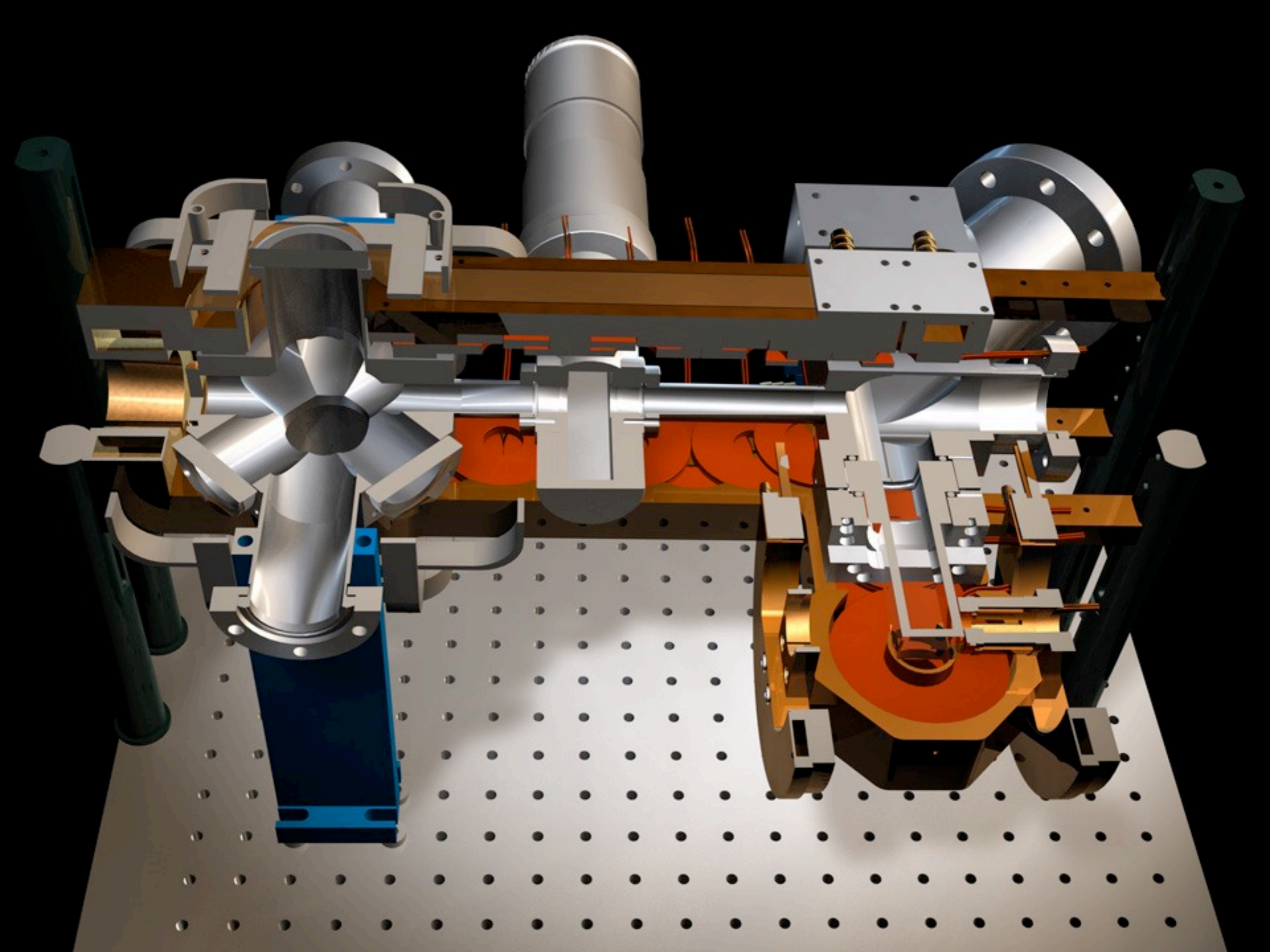


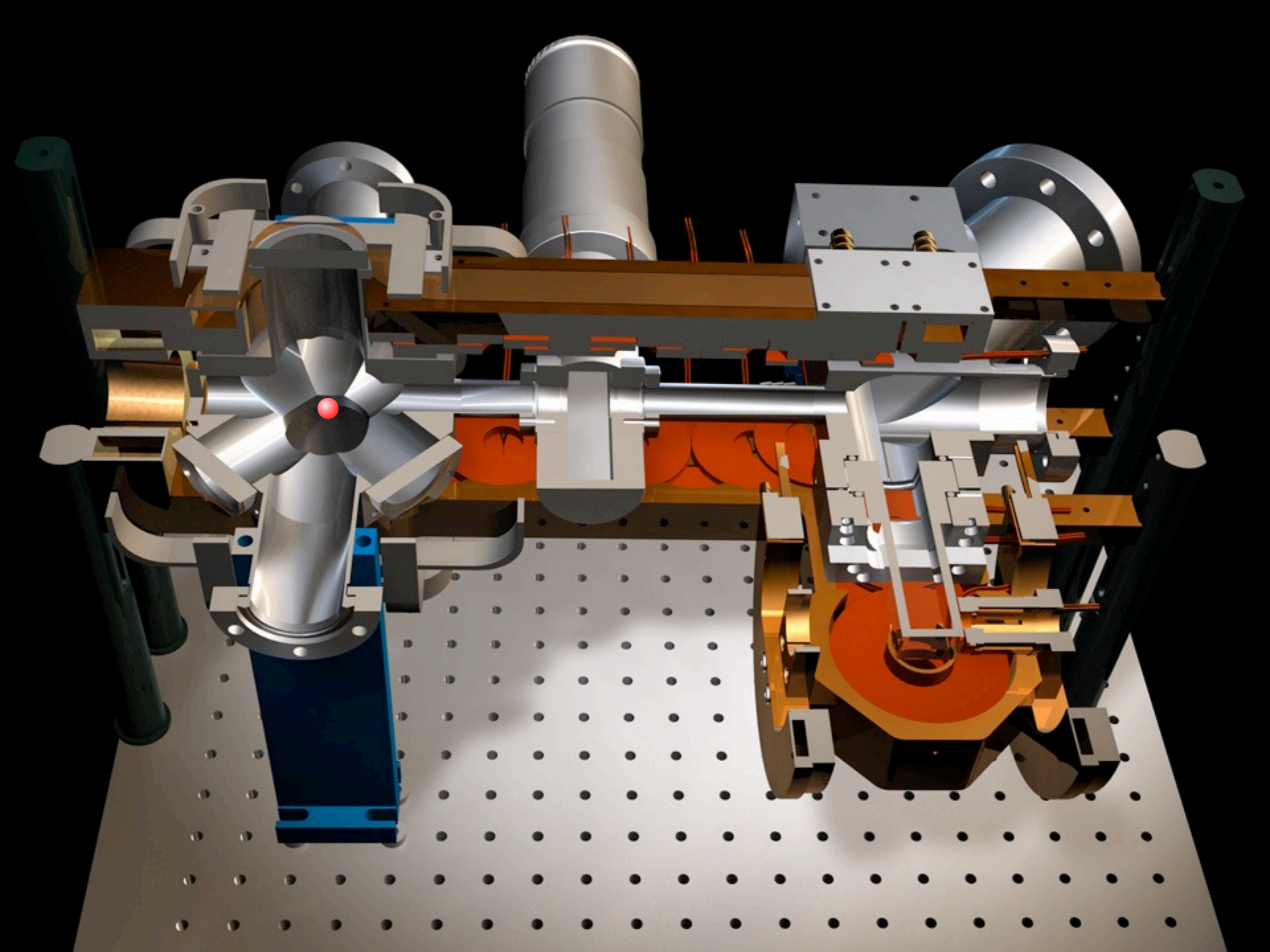
Bare chamber (no magnetic coils, no optics, etc.)

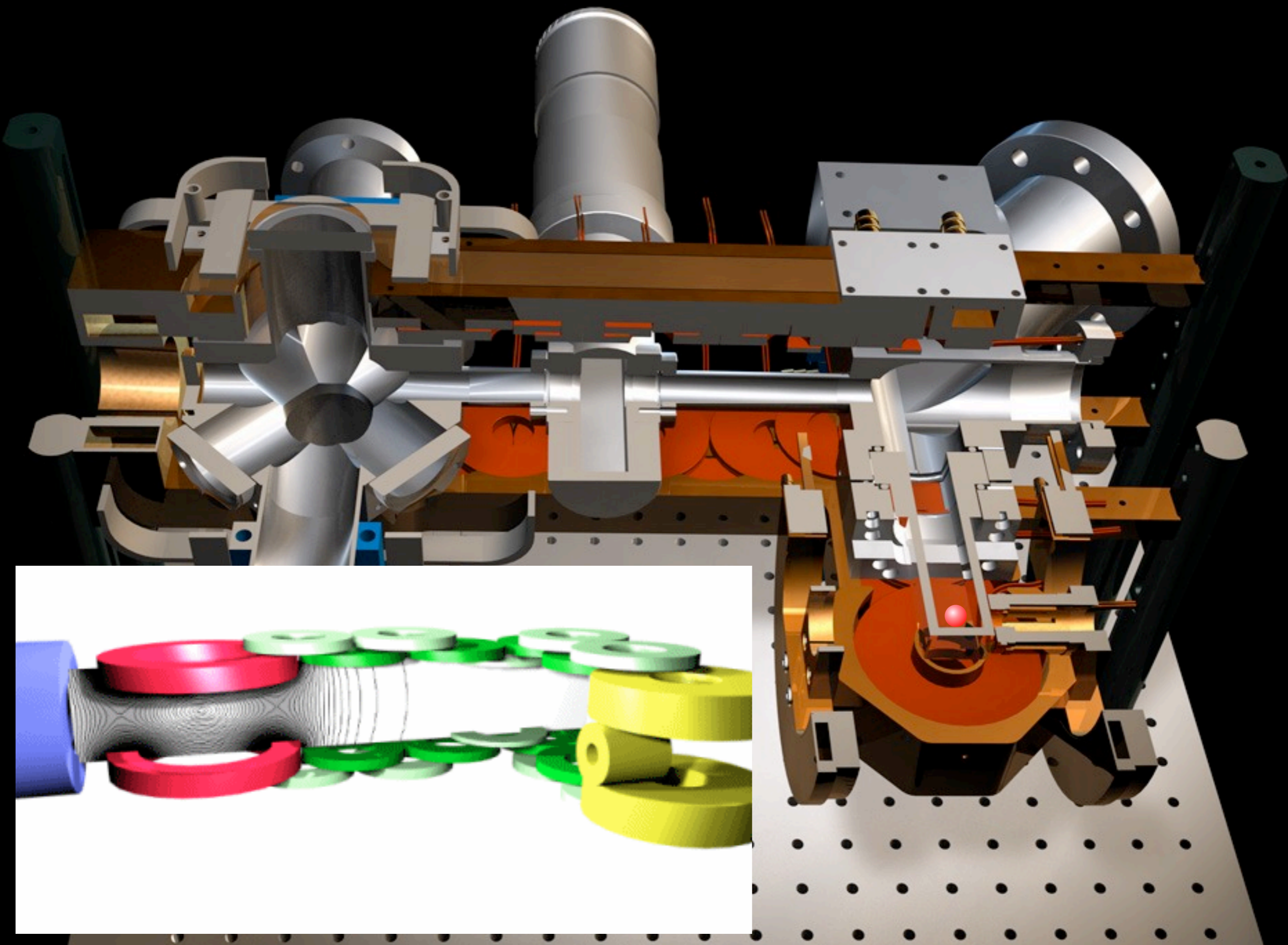
- ^{87}Rb (bosons)
- $N_{BEC} \sim 10^5$
- $T \sim 10\text{nK}$
- 3D optical lattice
 $\lambda_z = 843\text{nm}$
 $\lambda_{xs,ys} = 767\text{nm}$
- **1D Superlattice**
+ $\lambda_{xl} = 1534\text{nm}$
- **Another Superlattice**
+ $\lambda_{yl} = 1534\text{nm}$

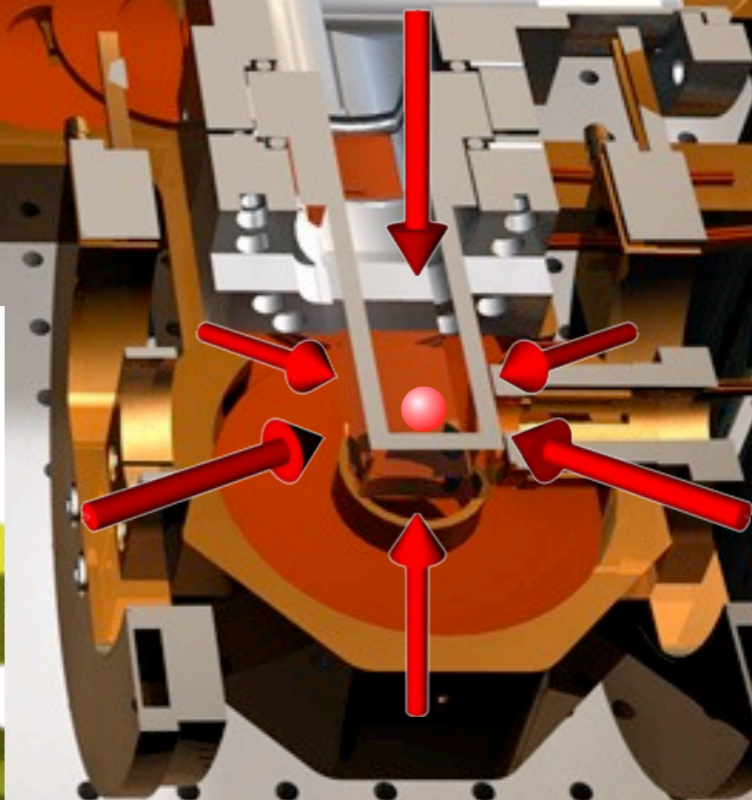
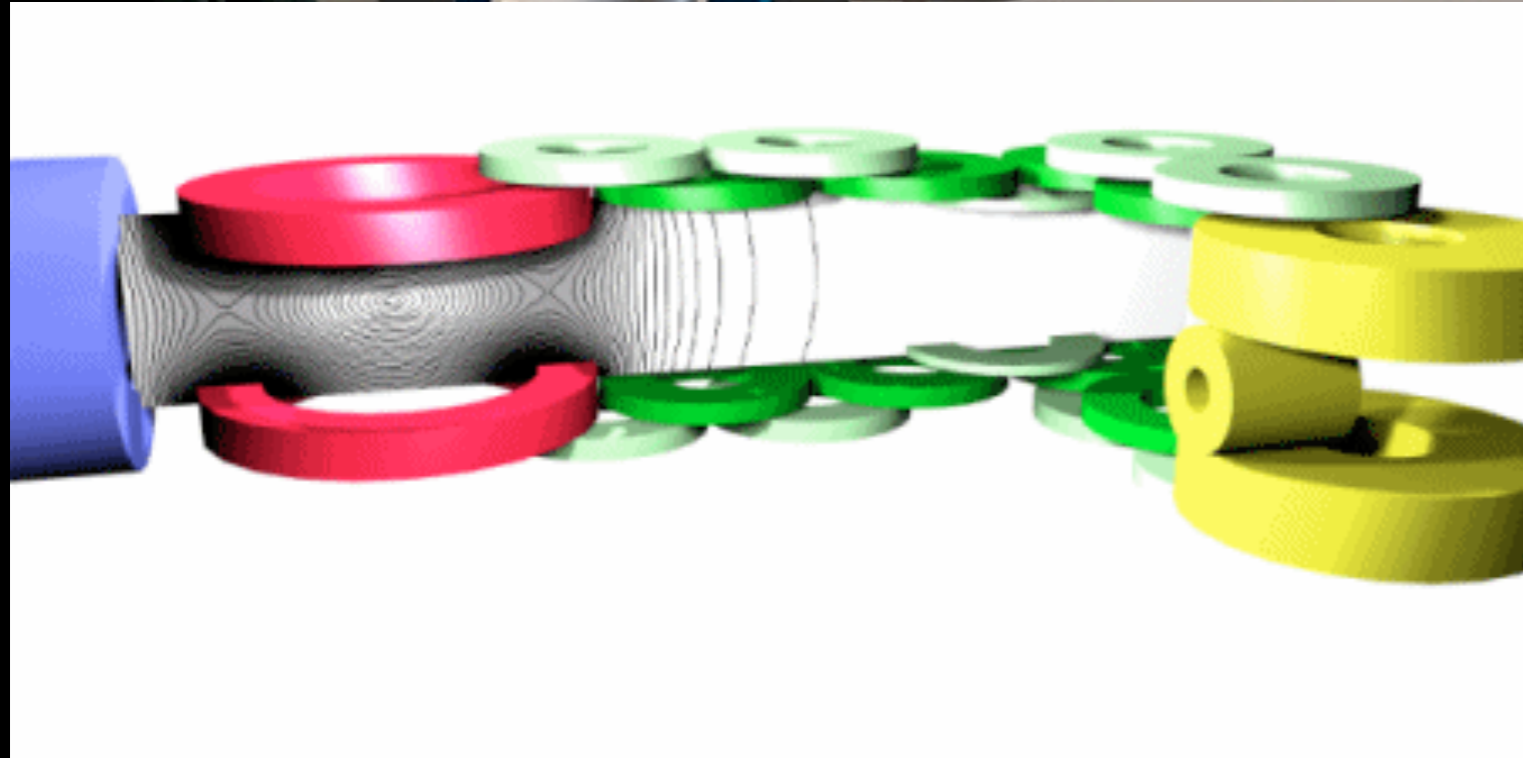
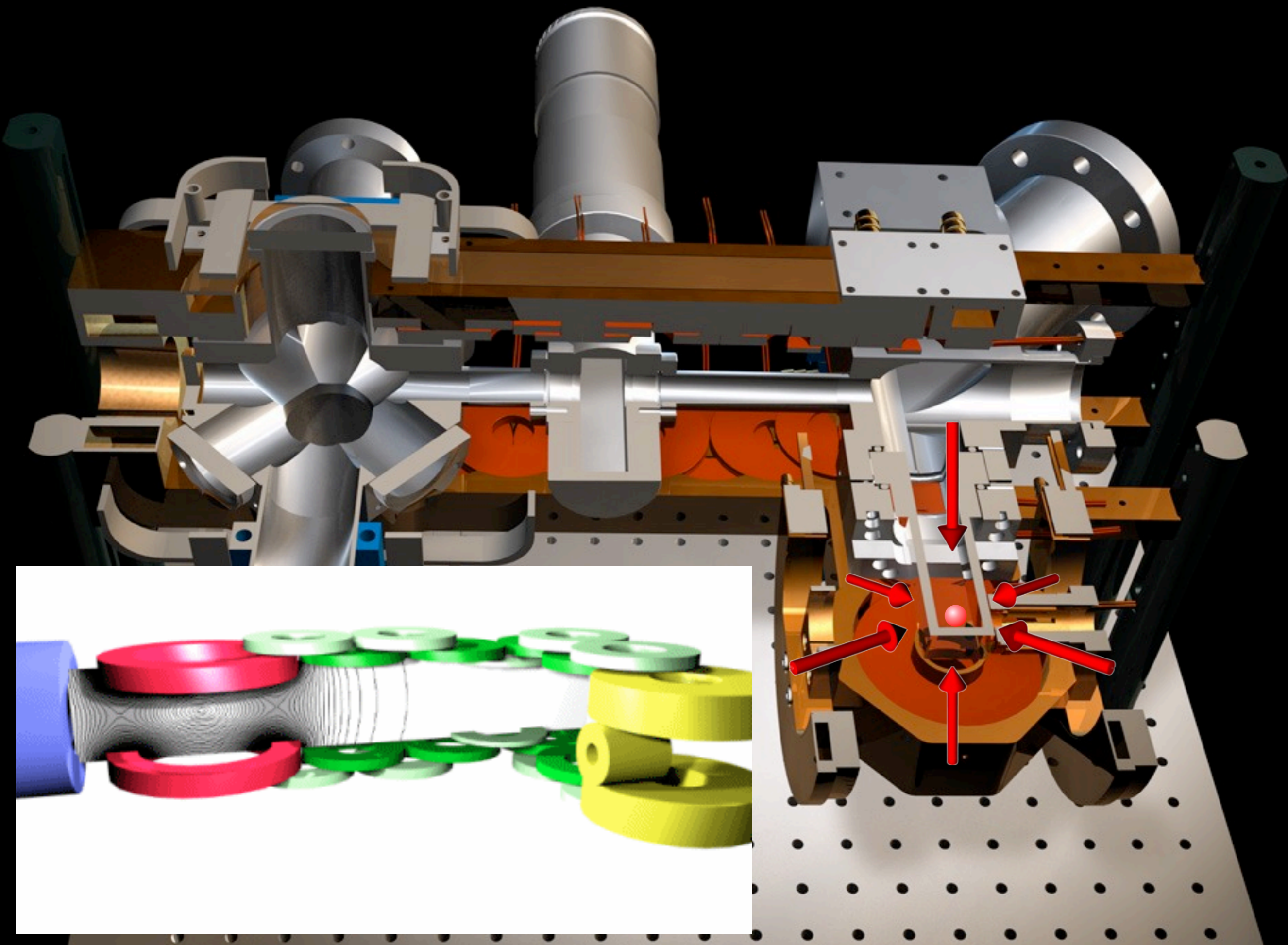


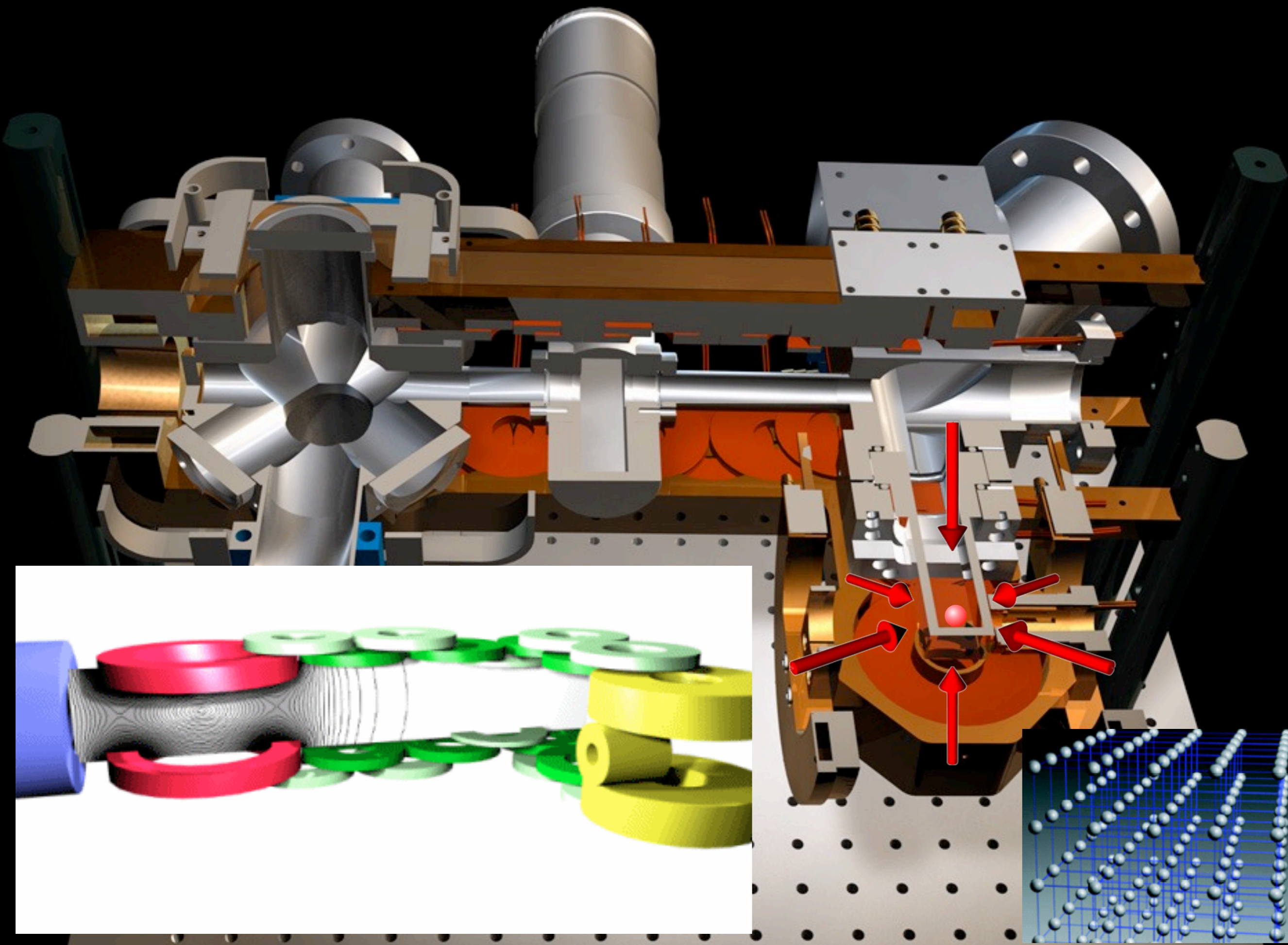












And a lot of optics and electronics !

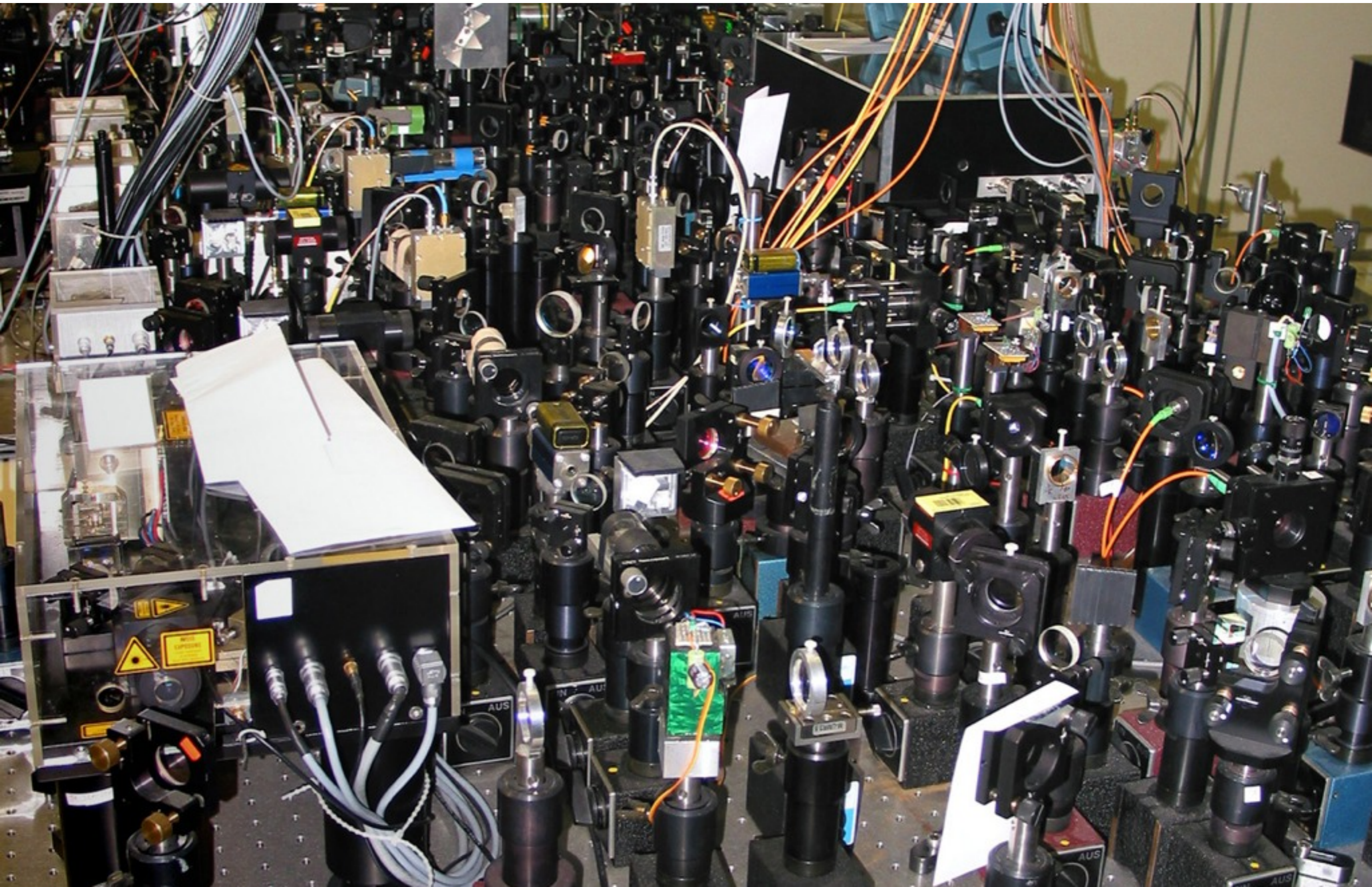
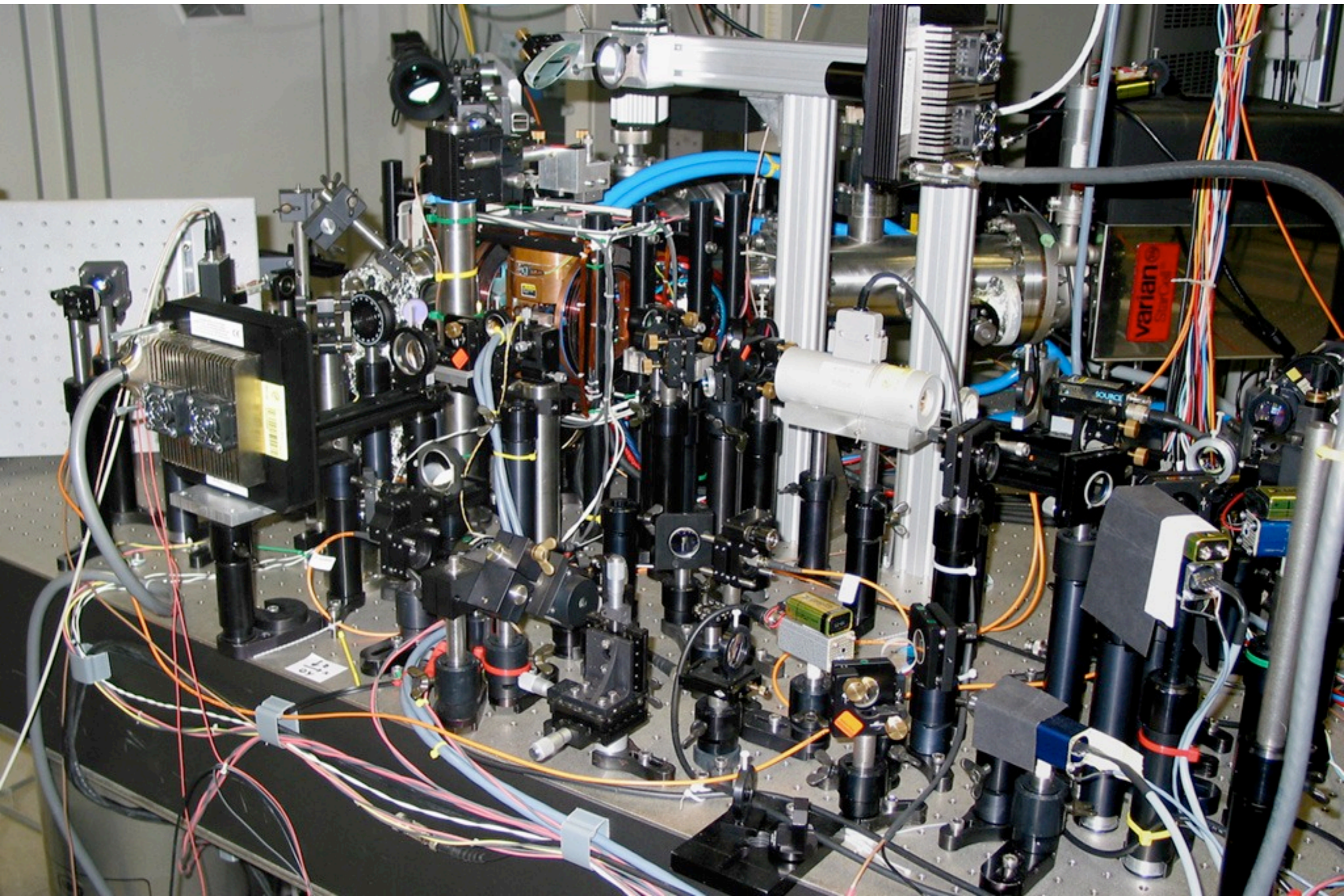


Table 2

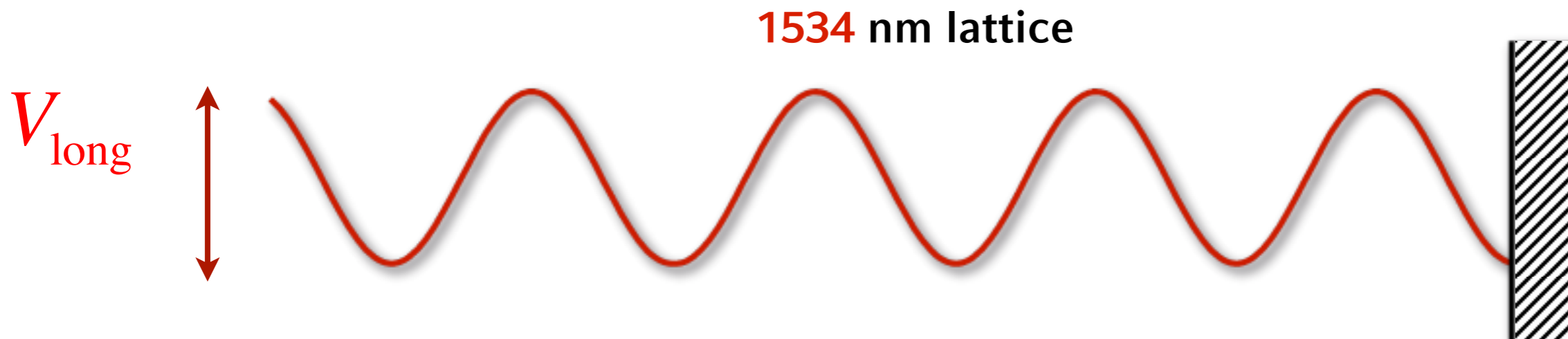




The Bichromatic Superlattice

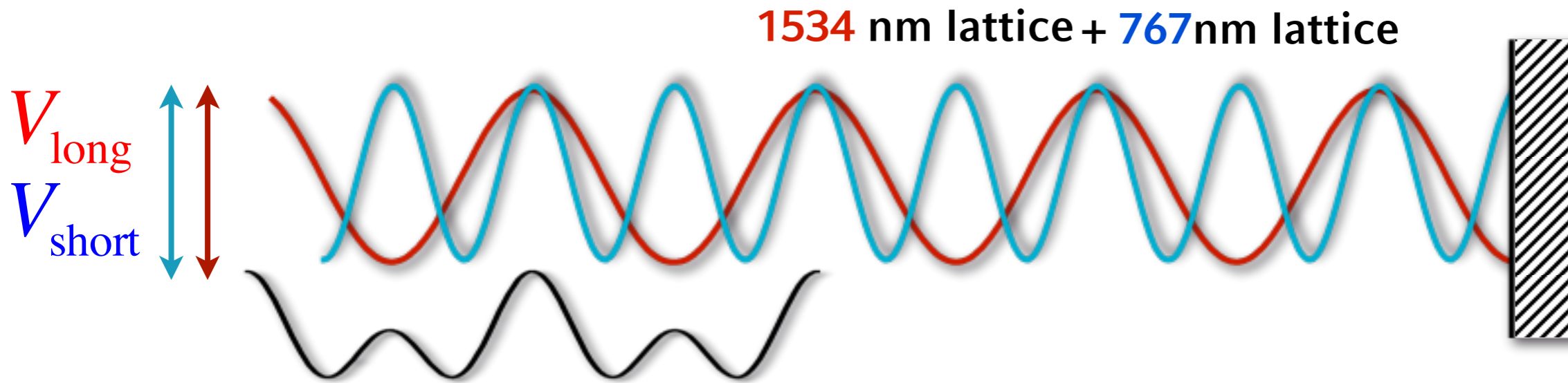


The Bichromatic Superlattice



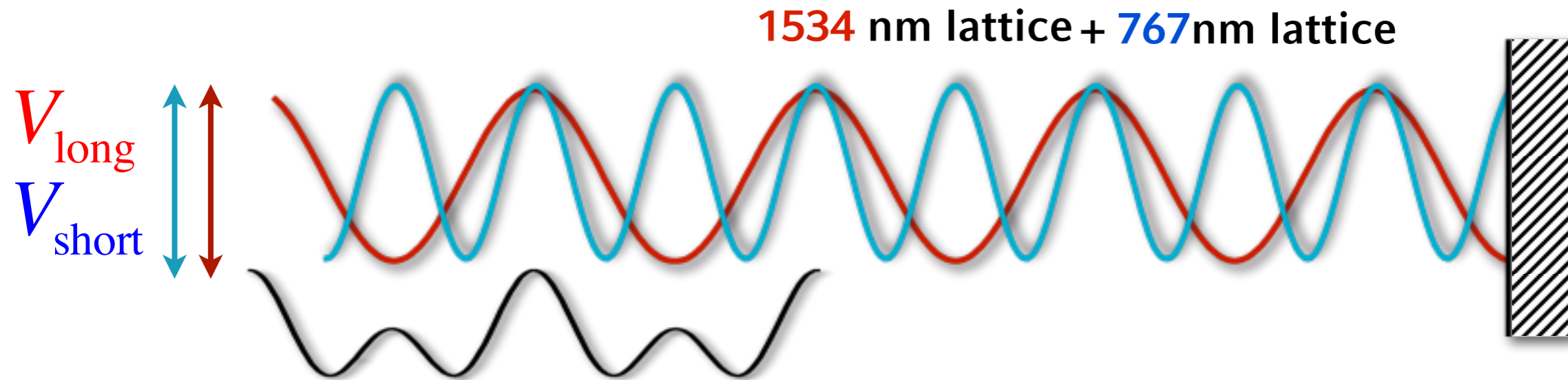


The Bichromatic Superlattice

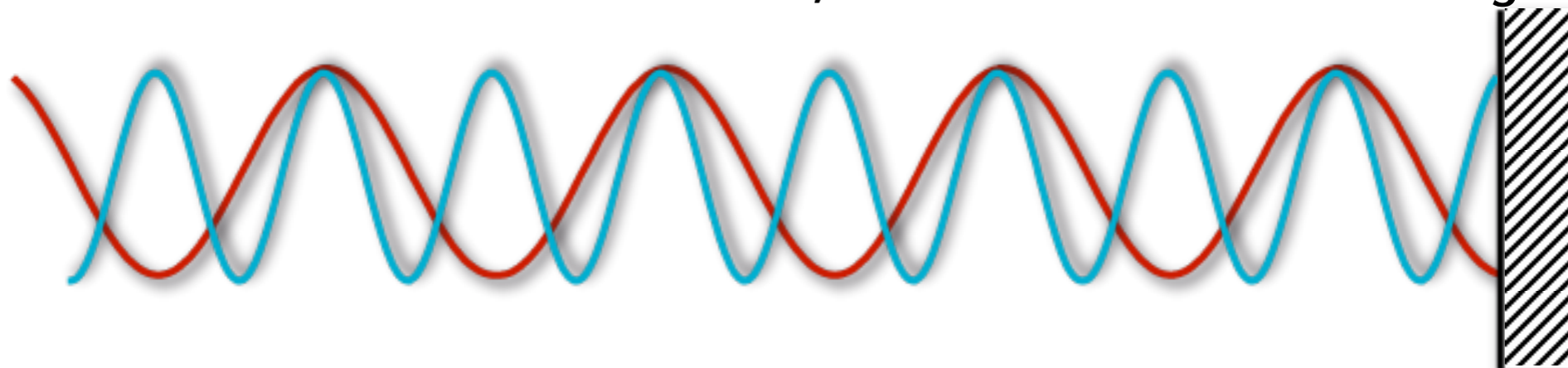




The Bichromatic Superlattice

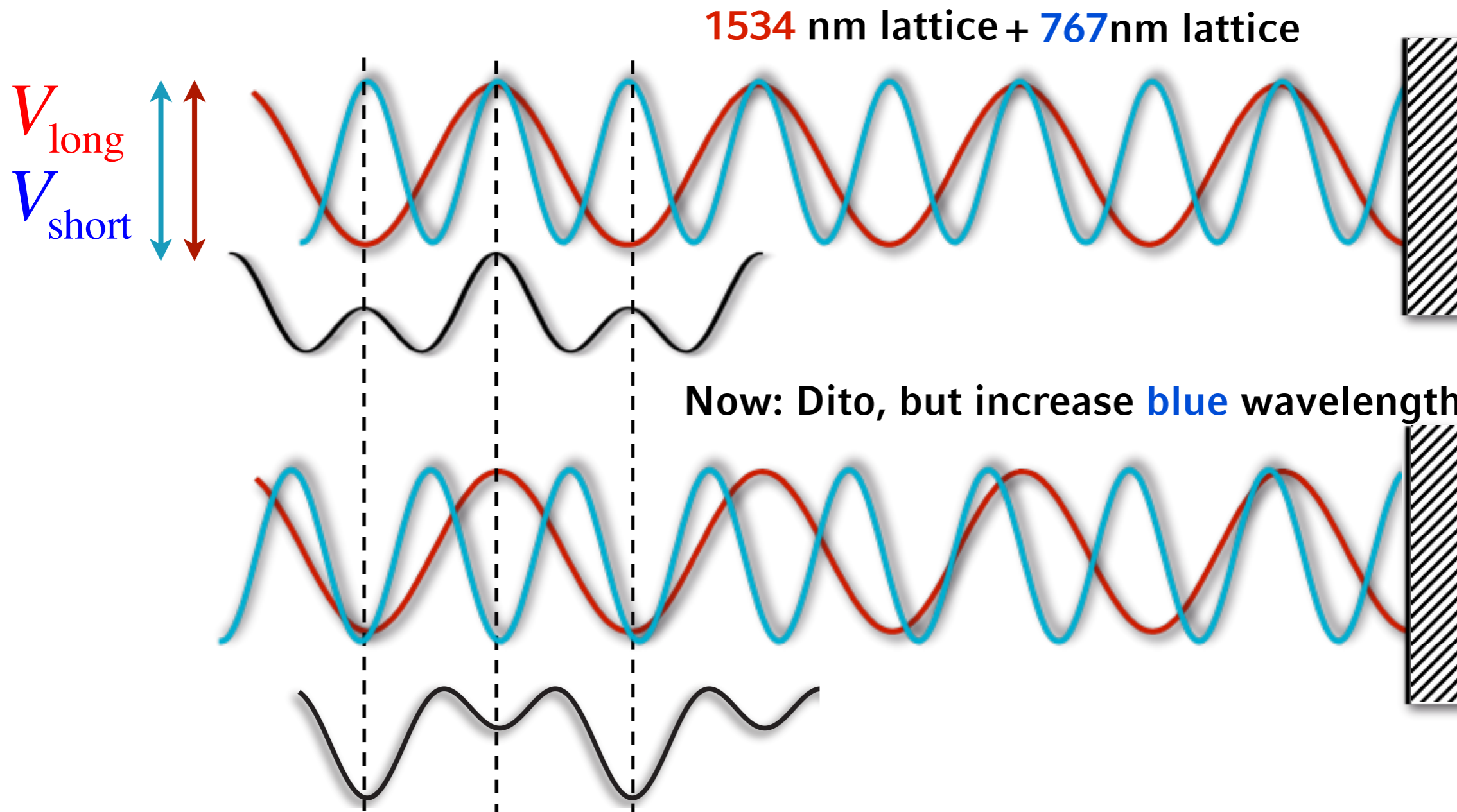


Now: Dito, but increase blue wavelength





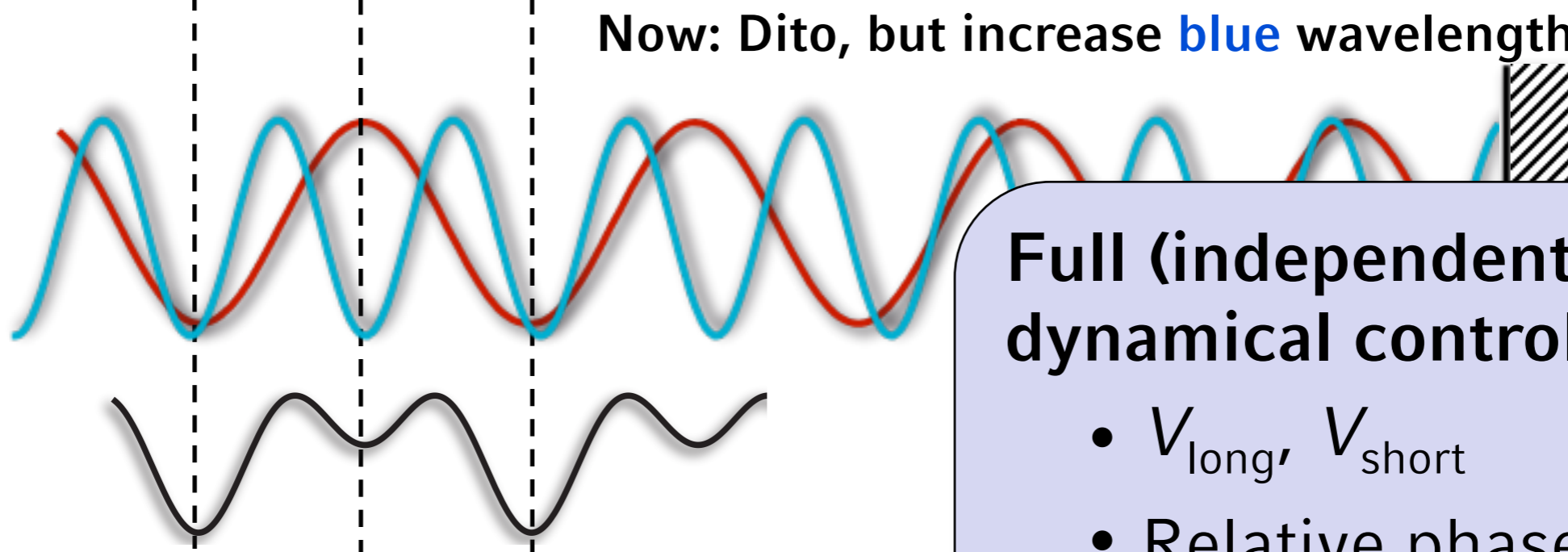
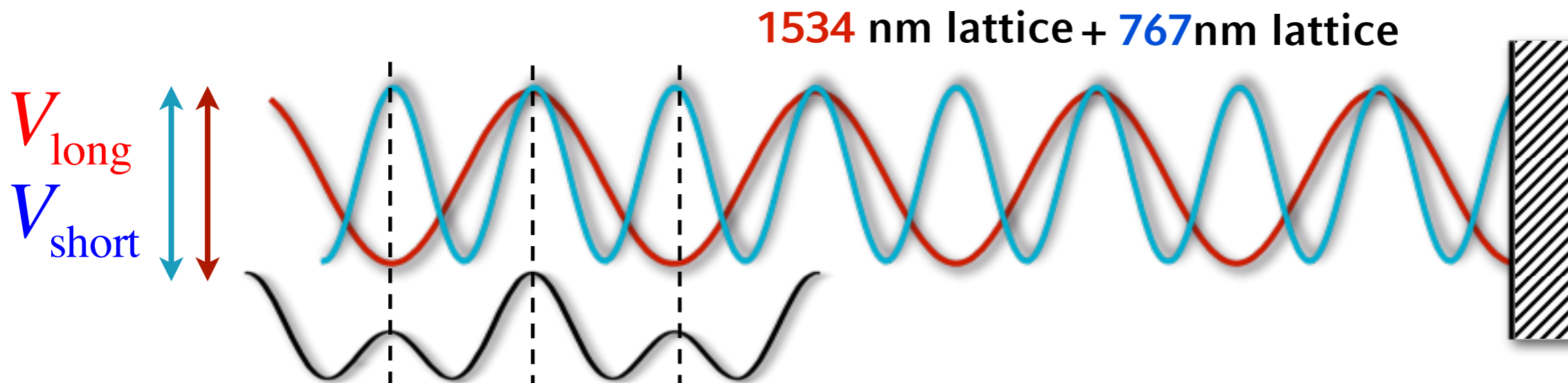
The Bichromatic Superlattice



$$V(x) = V_l \cos(2k_l x) + V_s \cos(4k_l x + \varphi)$$



The Bichromatic Superlattice



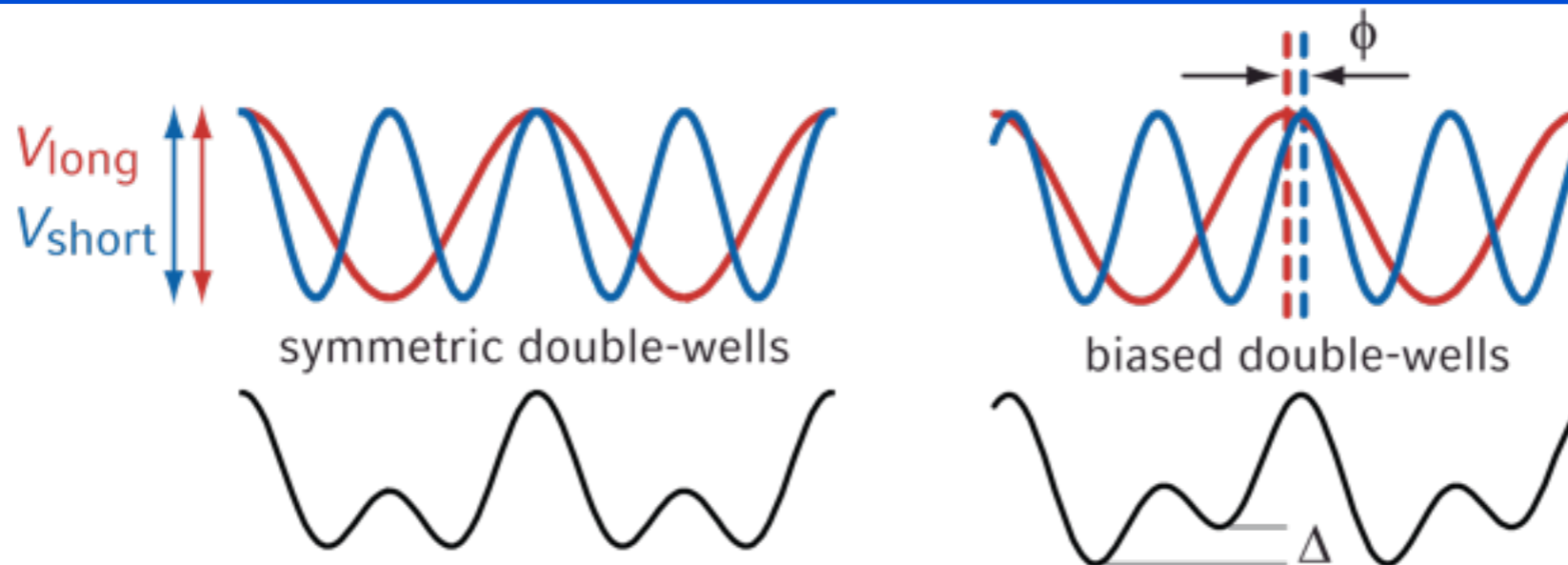
$$V(x) = V_l \cos(2k_l x) + V_s \cos(4k_l x + \varphi)$$

**Full (independent)
dynamical control over:**

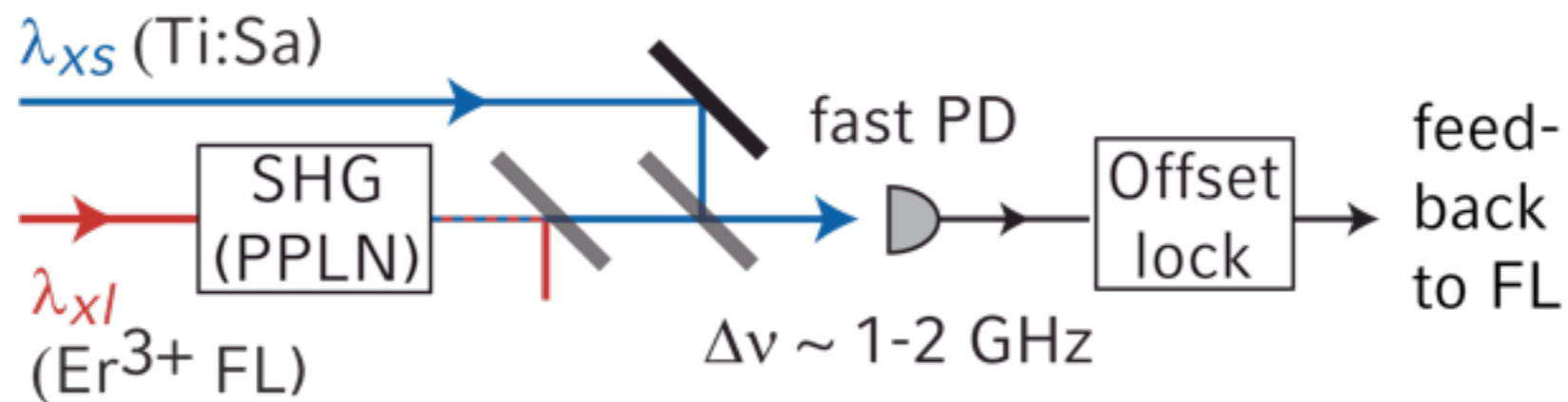
- $V_{\text{long}}, V_{\text{short}}$
- Relative phase φ
- Transverse lattices



The Bichromatic Superlattice



- Adjusting φ by fine-tuning λ_l
- Offset-locking frequency-doubled fiber-laser to Ti:Sa-laser

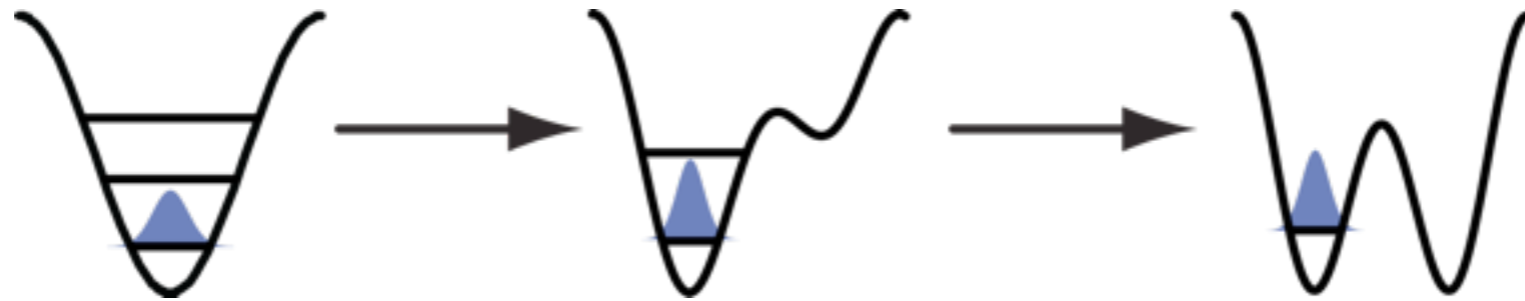


- Offset-frequency $\Delta\nu = 1 - 2\text{GHz}$
→ Tuning range $\varphi = 0 \dots 2\pi$



The Bichromatic Superlattice

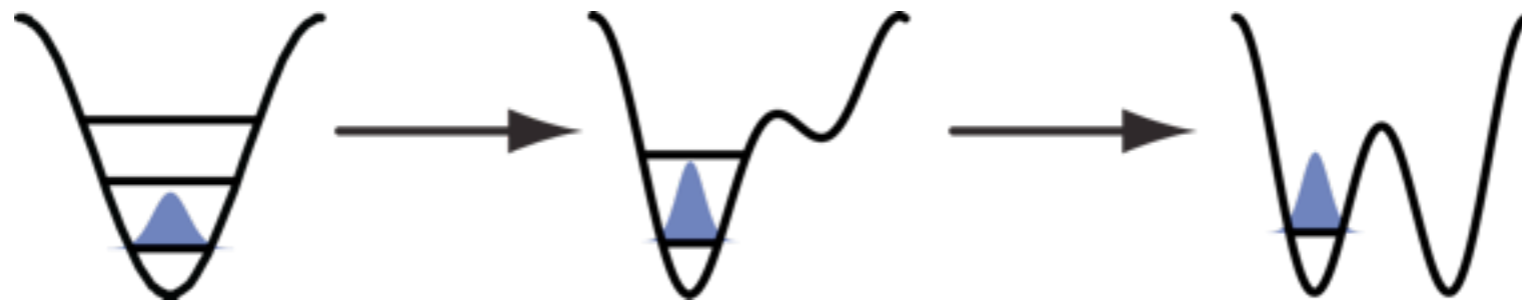
- Novel *state preparation* techniques, e.g. **patterned loading**



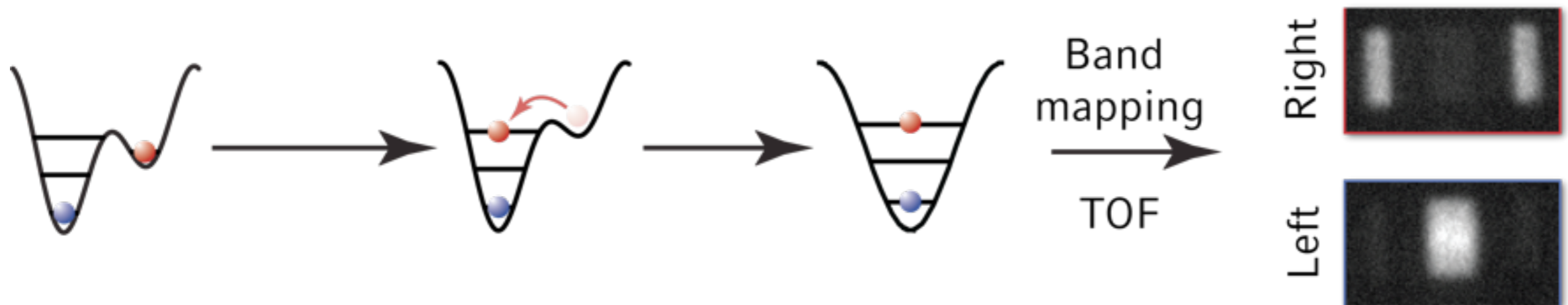


The Bichromatic Superlattice

- Novel *state preparation* techniques, e.g. **patterned loading**



- Novel *read-out* methods, e.g. **sub-lattice resolved detection**

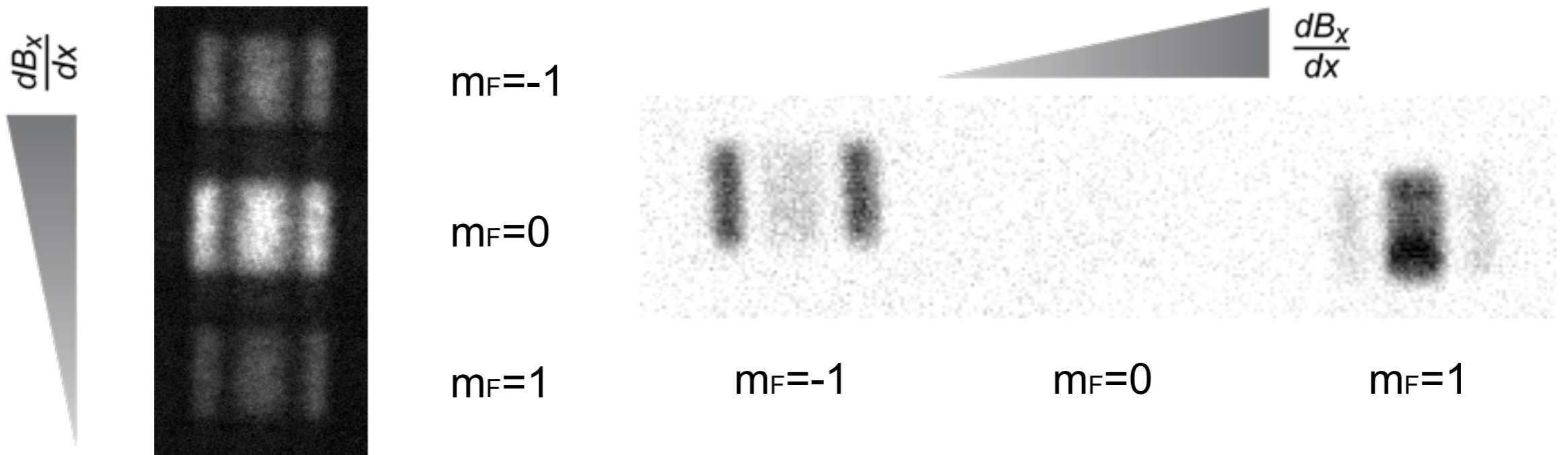


(Alternative: Raman-spectroscopy)



The Bichromatic Superlattice

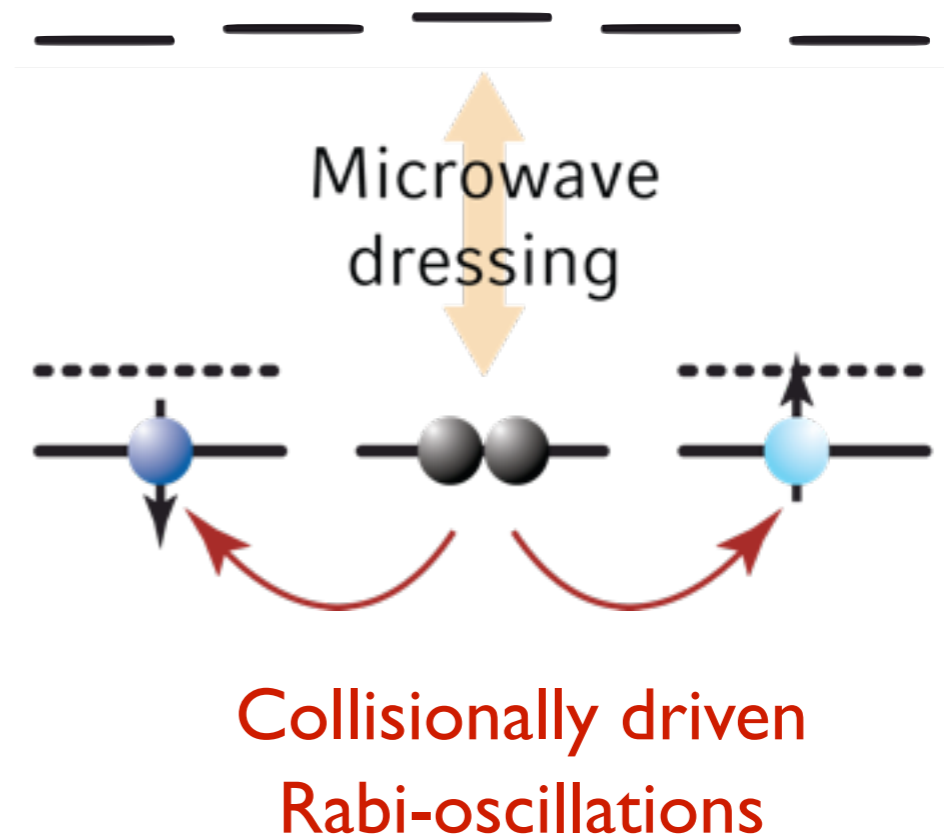
- Novel *state read out* techniques, e.g. **Stern-Gerlach**





Preparing Spin Singlets

- Atom pairs in long-lattice wells $|F = -1, m_F = 0\rangle$
- Initialize in $|F = 1, m_F = 0\rangle$
- Microwave-dressed spin-changing collisions
→ **Spin-pairs** in $|F = 1, m_F = \pm 1\rangle$

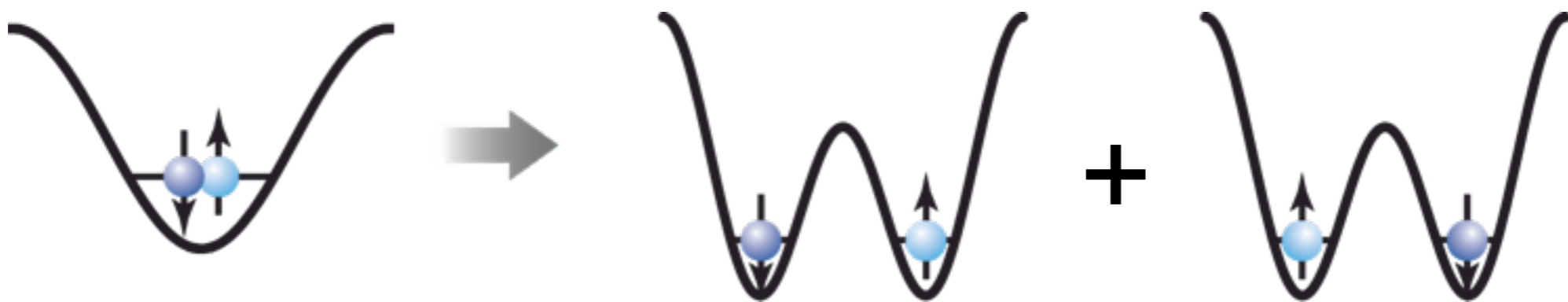




Preparing Entangled Spin Singlets

- **Spin pairs** in $|F = 1, m_F = \pm 1\rangle \equiv |\uparrow\rangle, |\downarrow\rangle$
- Barrier raised *slowly* to split
→ Crossing a miniature Mott-transition: $n_{\text{Left}} = n_{\text{Right}}$

J. Sebby-Strabley et al., PRL **98** (2007)



- **Bosons:** Symmetric wavefunction → Triplet $|t_0\rangle$
(Fermions: Antisymmetric wavefunction → Singlet $|s_0\rangle$)

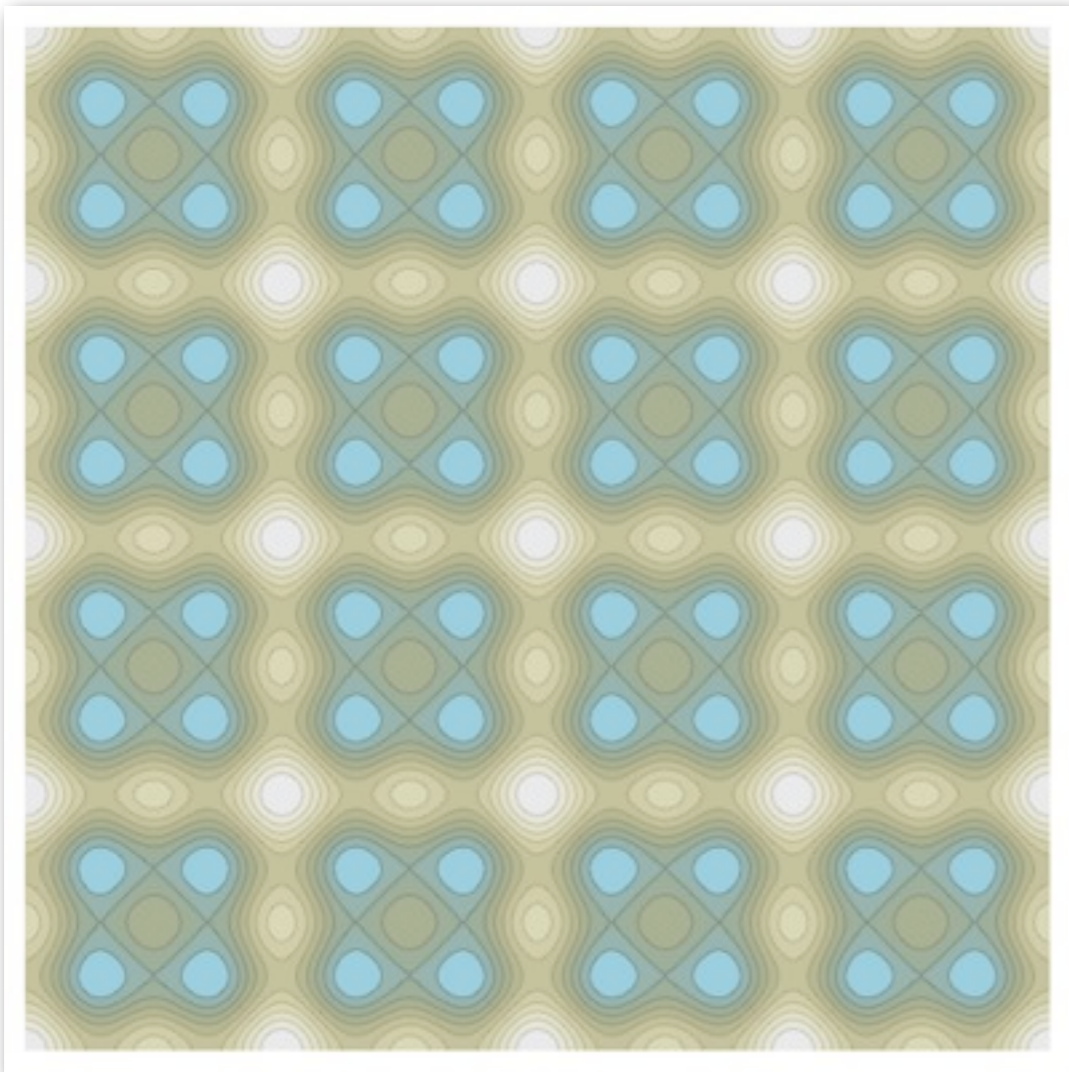
Details on the loading of the Spin-pairs:

S. Trotzky et al., Science **319** (2008), A.-M. Rey et al., PRL **99** (2007)



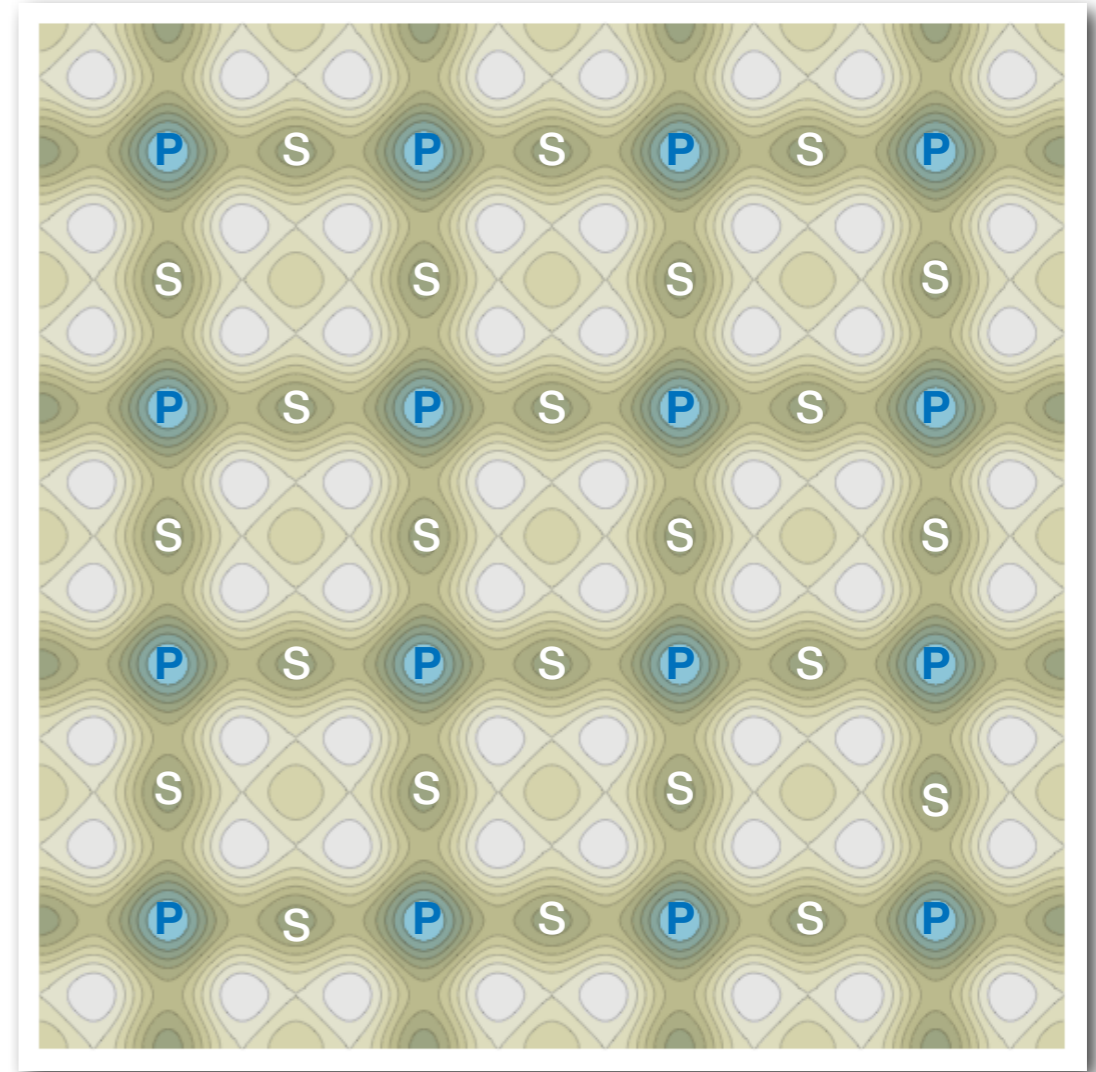


2D Superlattice Geometries (2 SL)



Coupled Plaquette Systems

see B. Paredes & I. Bloch, PRA **77**, 23603 (2008)
S. Trebst et al., PRL **96**, 250402 (2006)



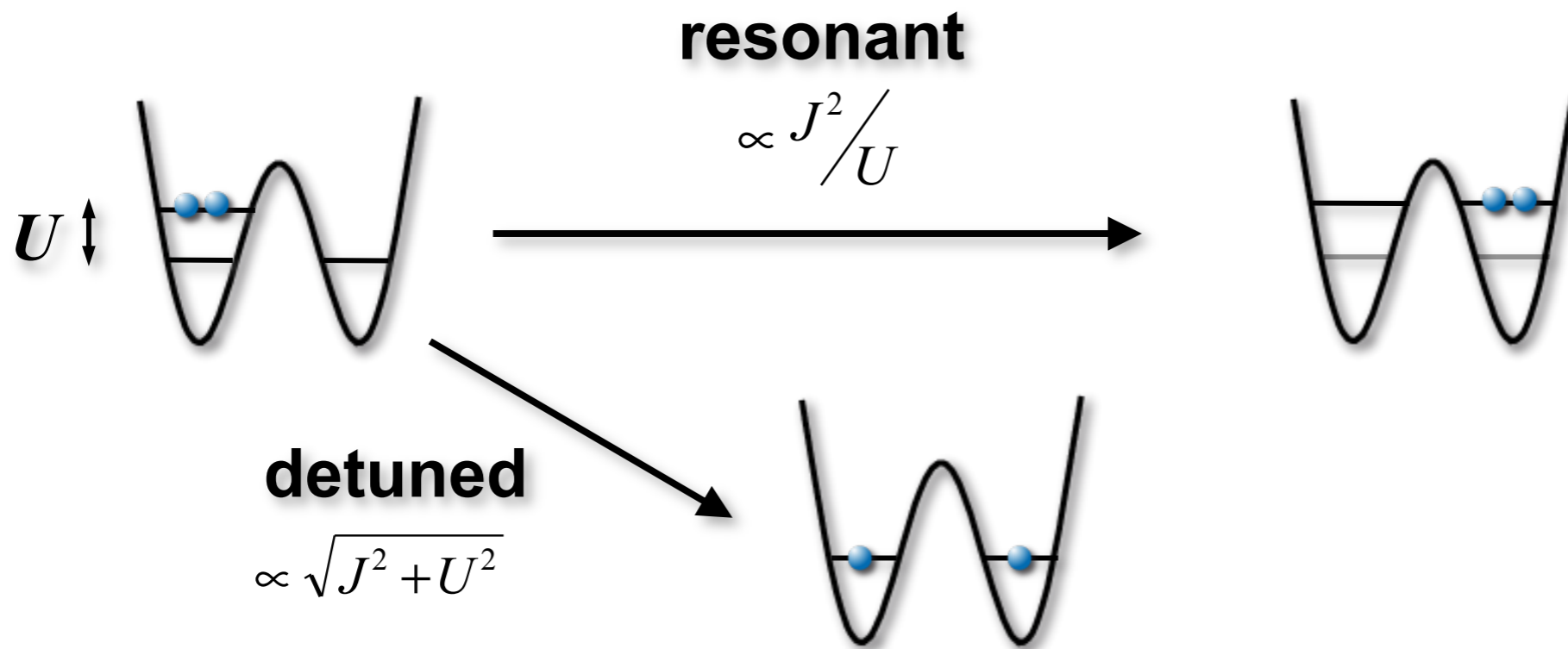
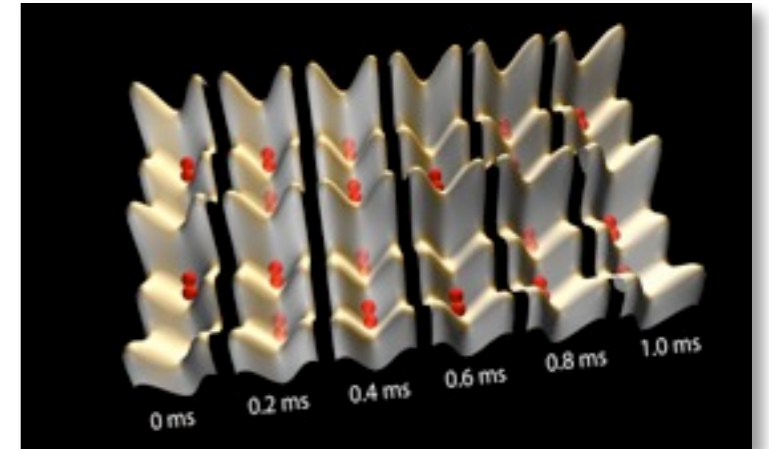
Higher Lattice Orbital Physics

see V. Liu, A. Ho, C. Wu and others work
exp: related to A. Hemmerich's exp.



Experiments in the Superlattice

- **Isolated double-wells:**
 - Correlated tunneling

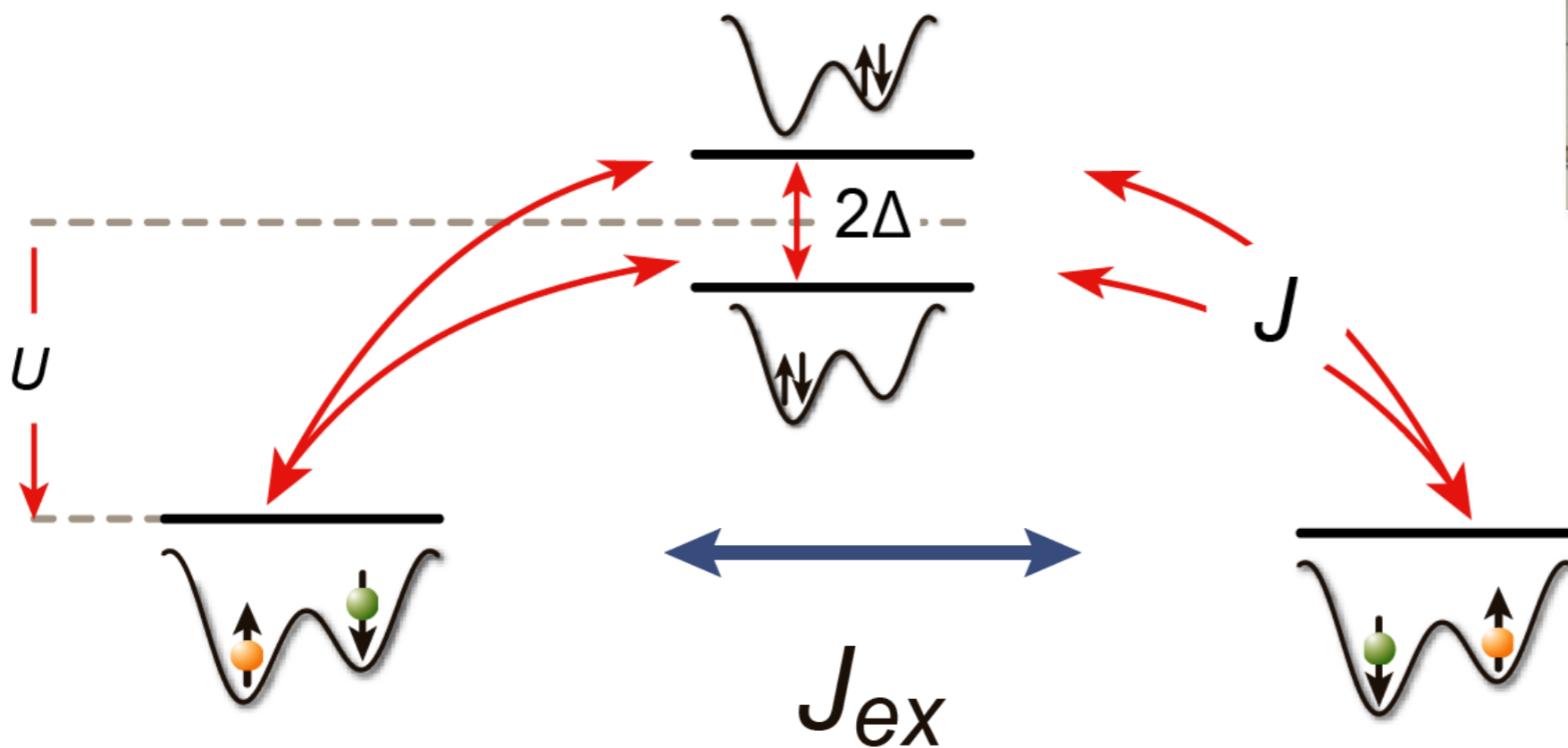
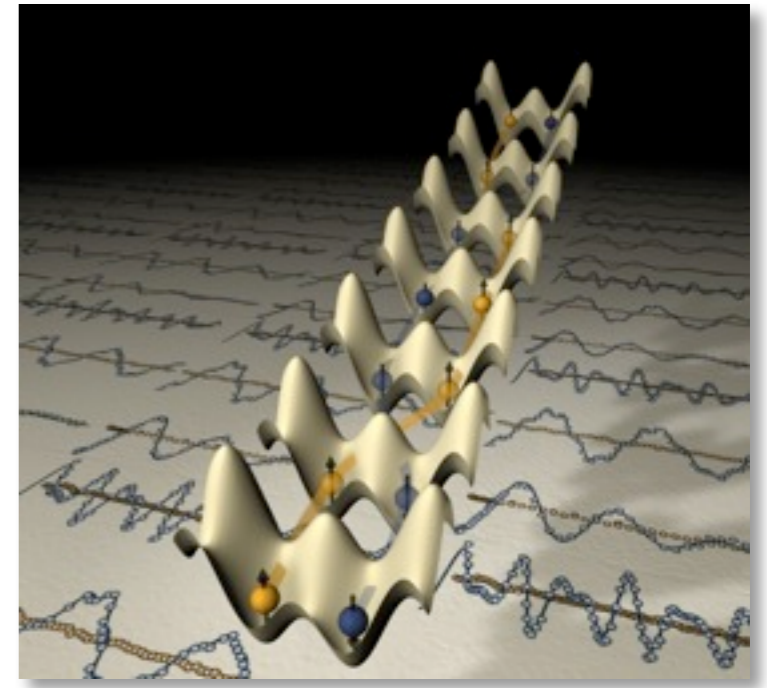




Experiments in the Superlattice

- **Isolated double-wells:**

- Correlated tunneling, Superexchange interactions

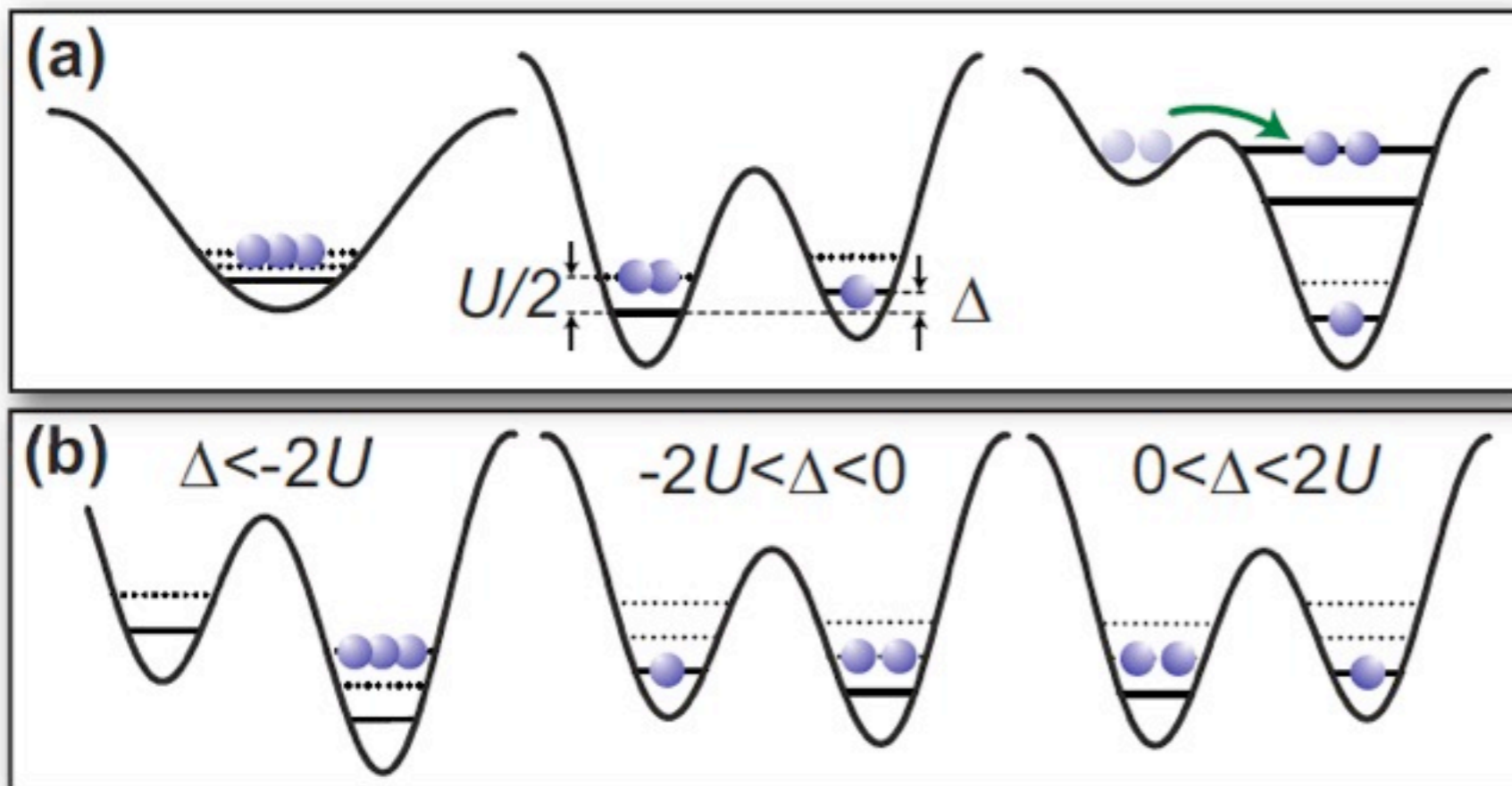
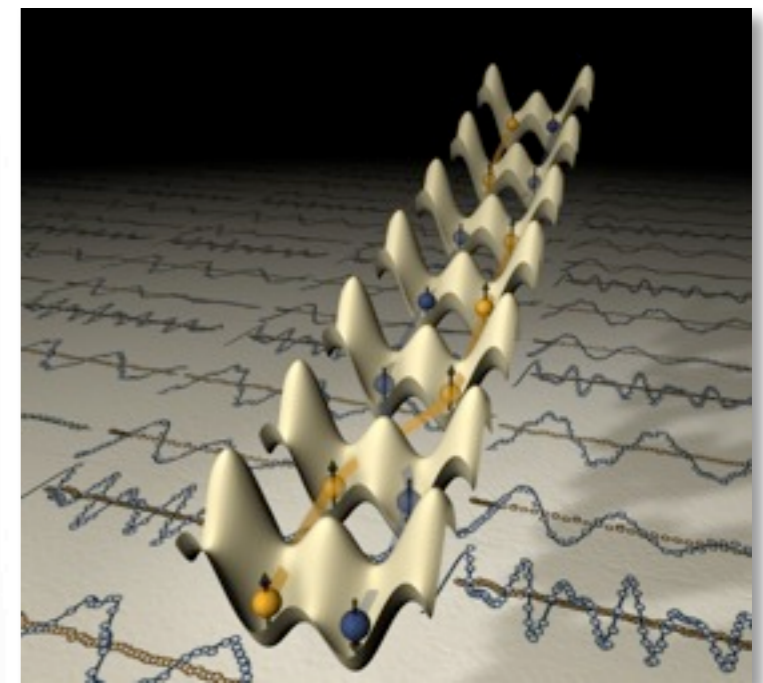
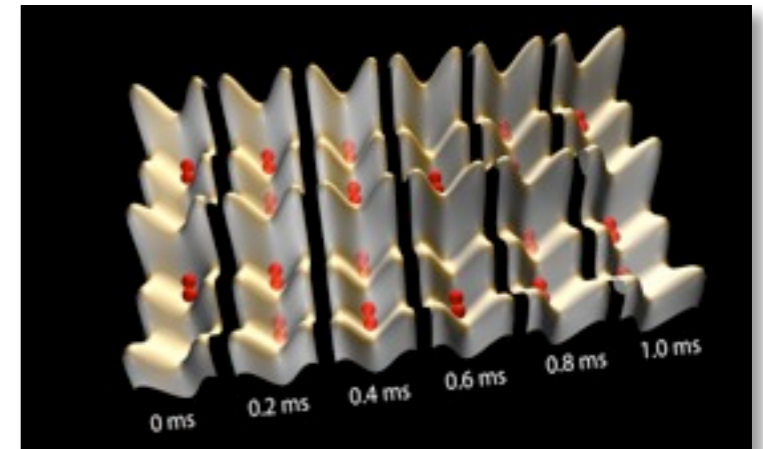




Experiments in the Superlattice

- **Isolated double-wells:**

- Correlated tunneling, Superexchange interactions
- Counting atoms via interaction blockade

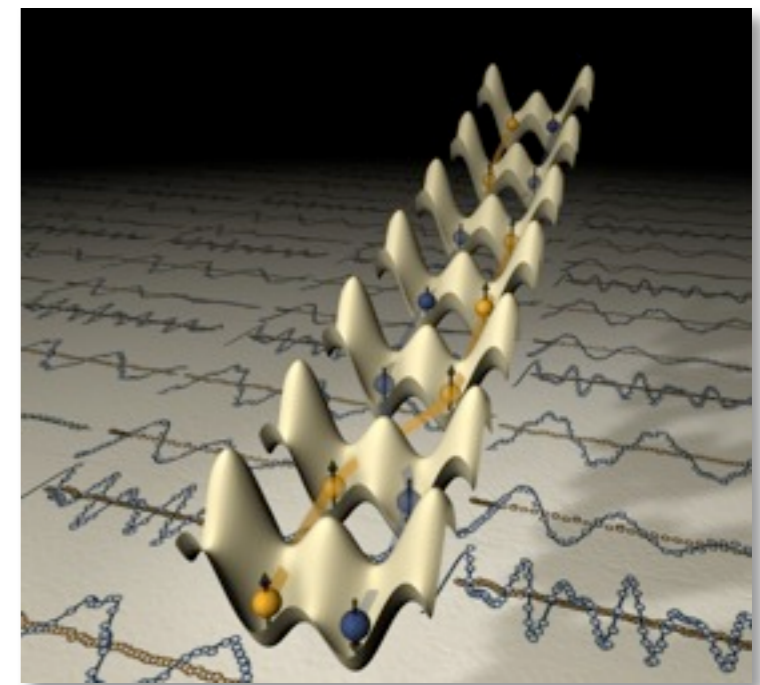
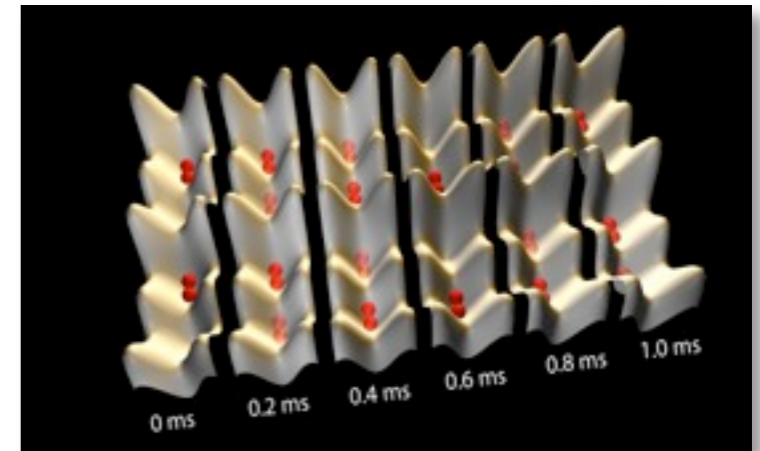
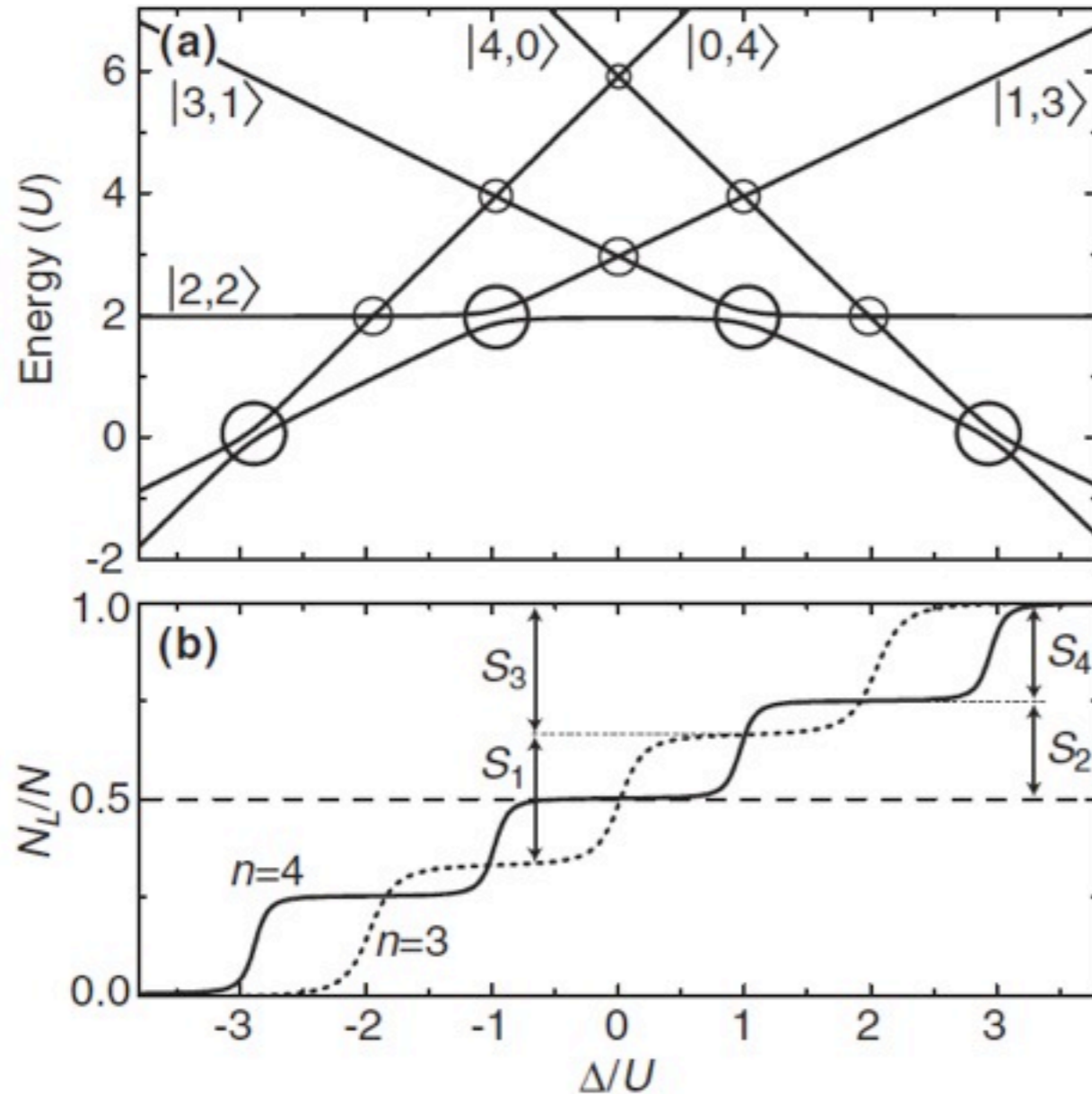




Experiments in the Superlattice

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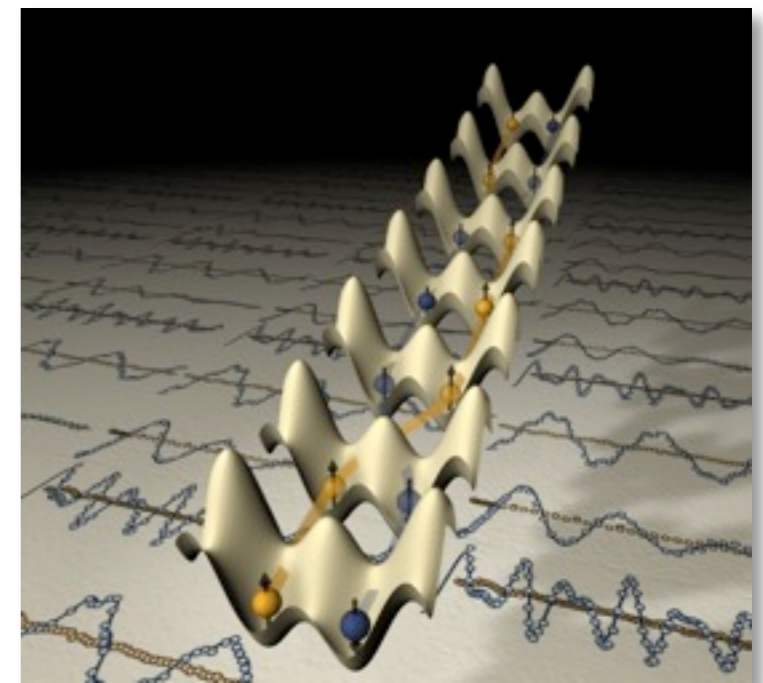
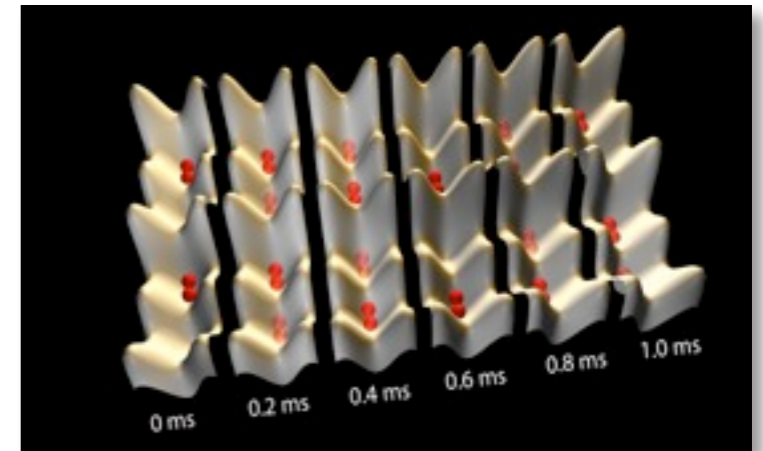
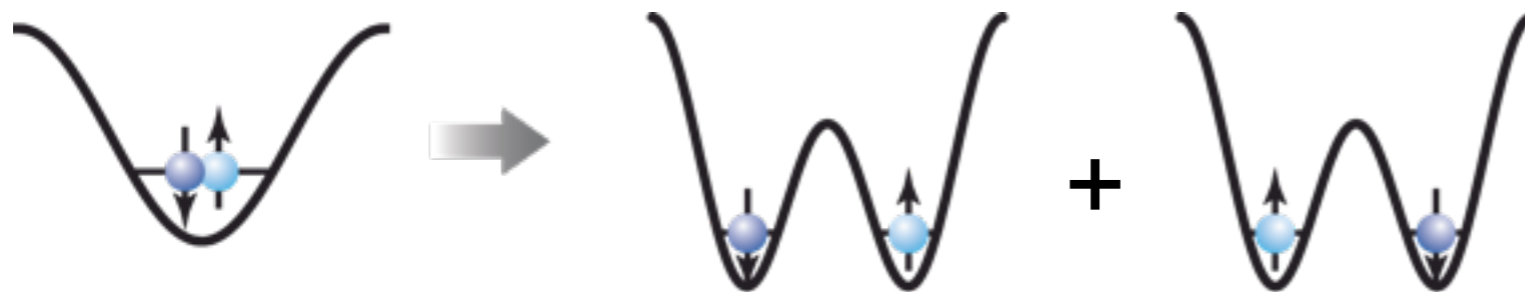




Experiments in the Superlattice

- **Isolated double-wells:**

- Correlated tunneling, Superexchange interactions
- Counting atoms via interaction blockade
- Control of n.n. spin correlations

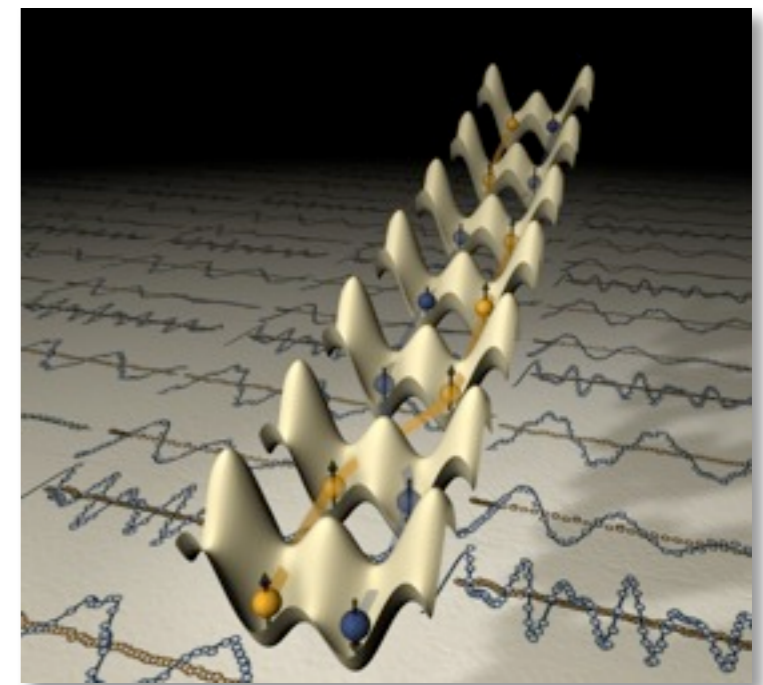
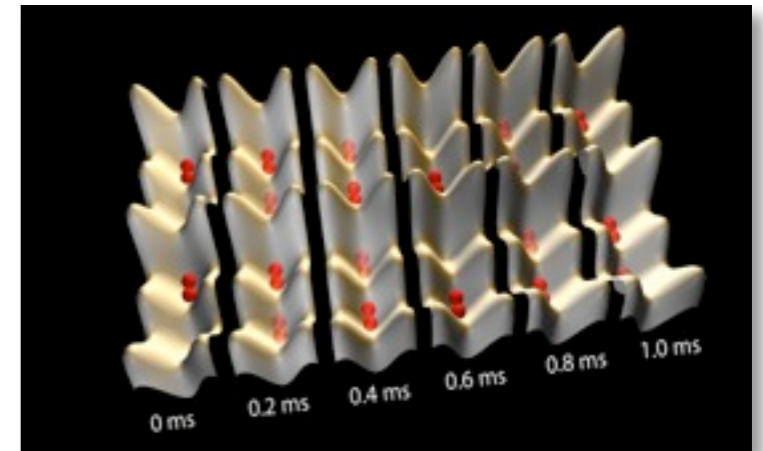
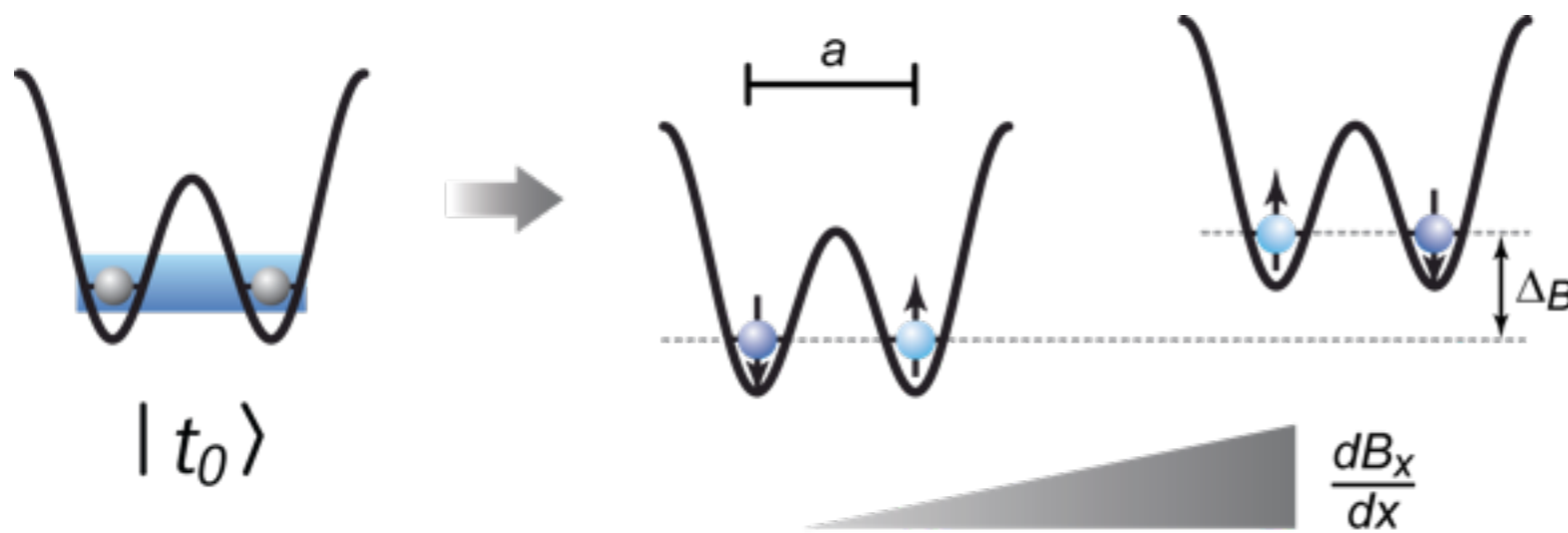




Experiments in the Superlattice

- **Isolated double-wells:**

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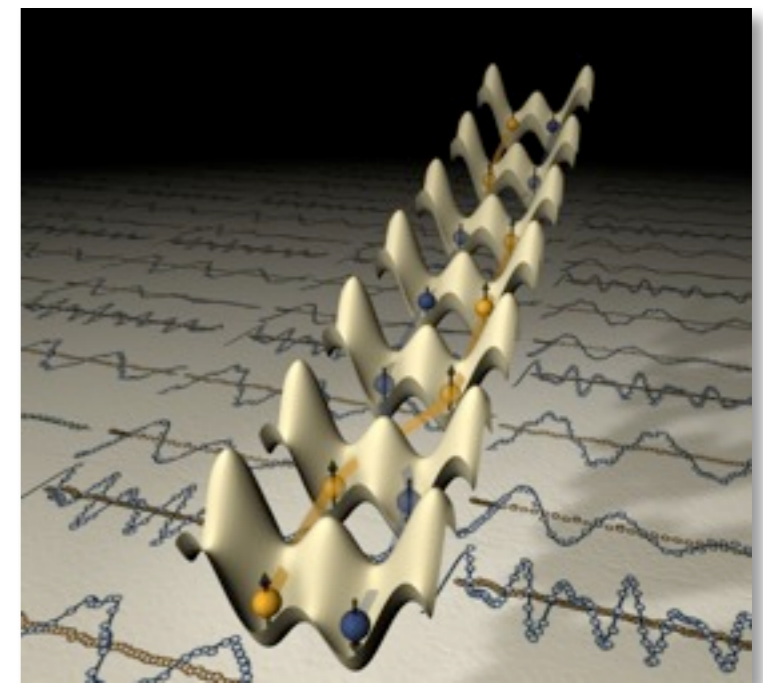
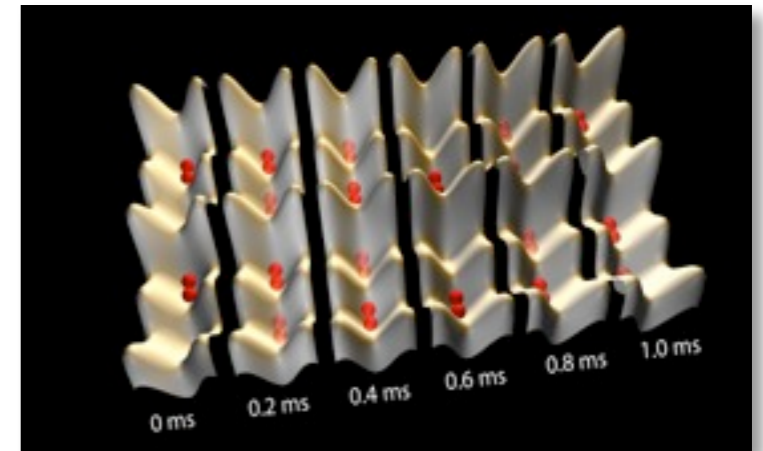
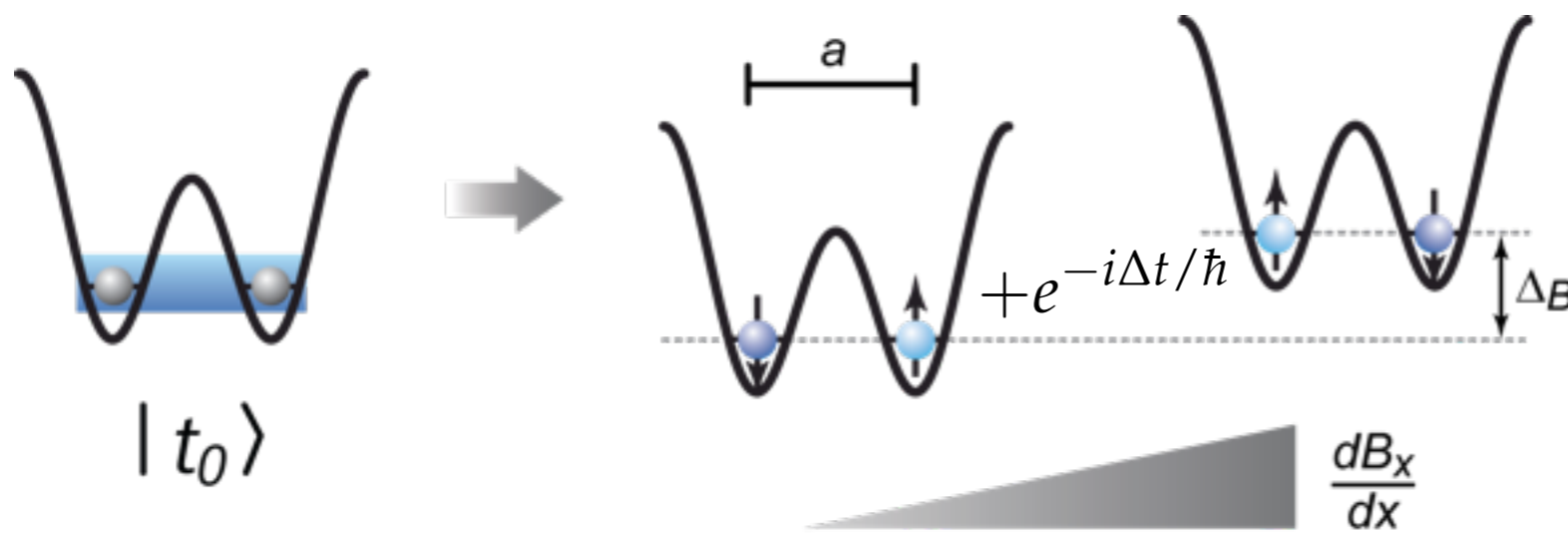




Experiments in the Superlattice

- **Isolated double-wells:**

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- Counting atoms via interaction blockade
- Control of n.n. spin correlations

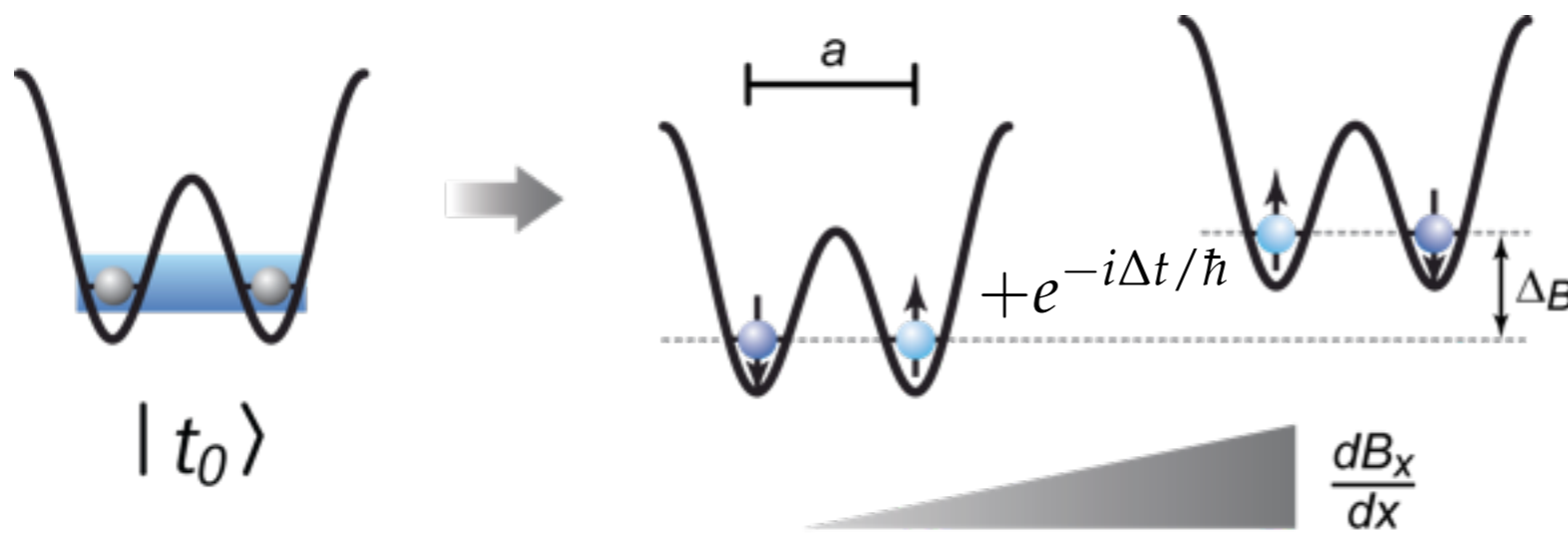
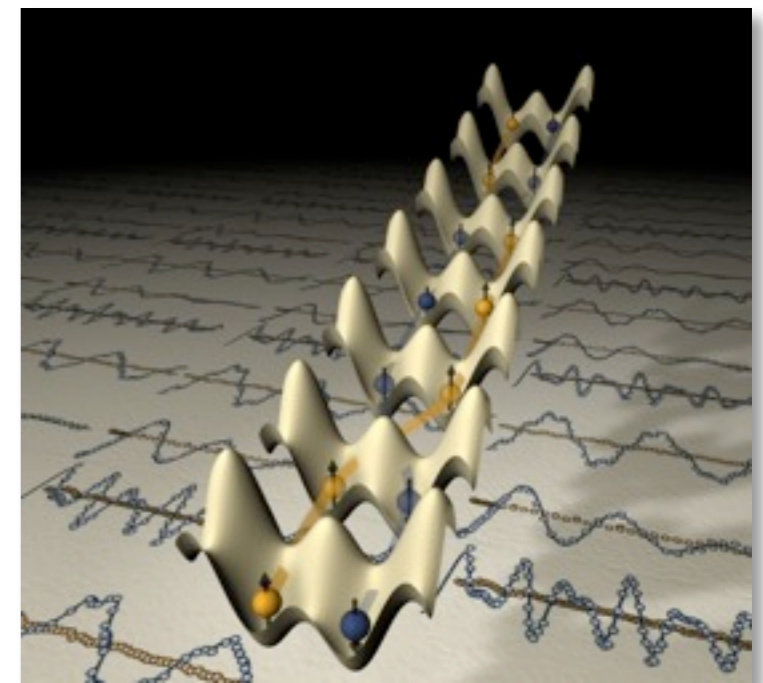
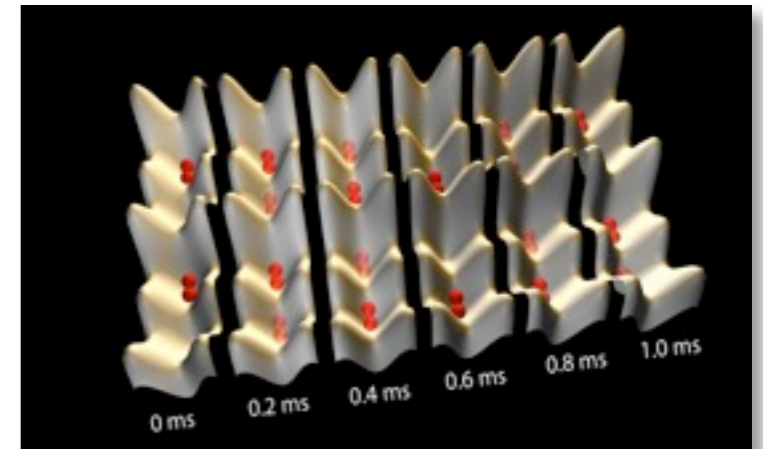




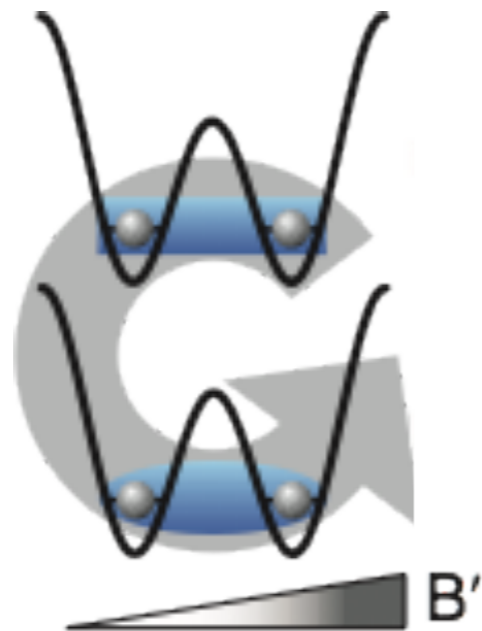
Experiments in the Superlattice

• Isolated double-wells:

- Correlated tunneling, Superexchange interactions
- Counting atoms via interaction blockade
- Control of n.n. spin correlations



→ **Triplet:**



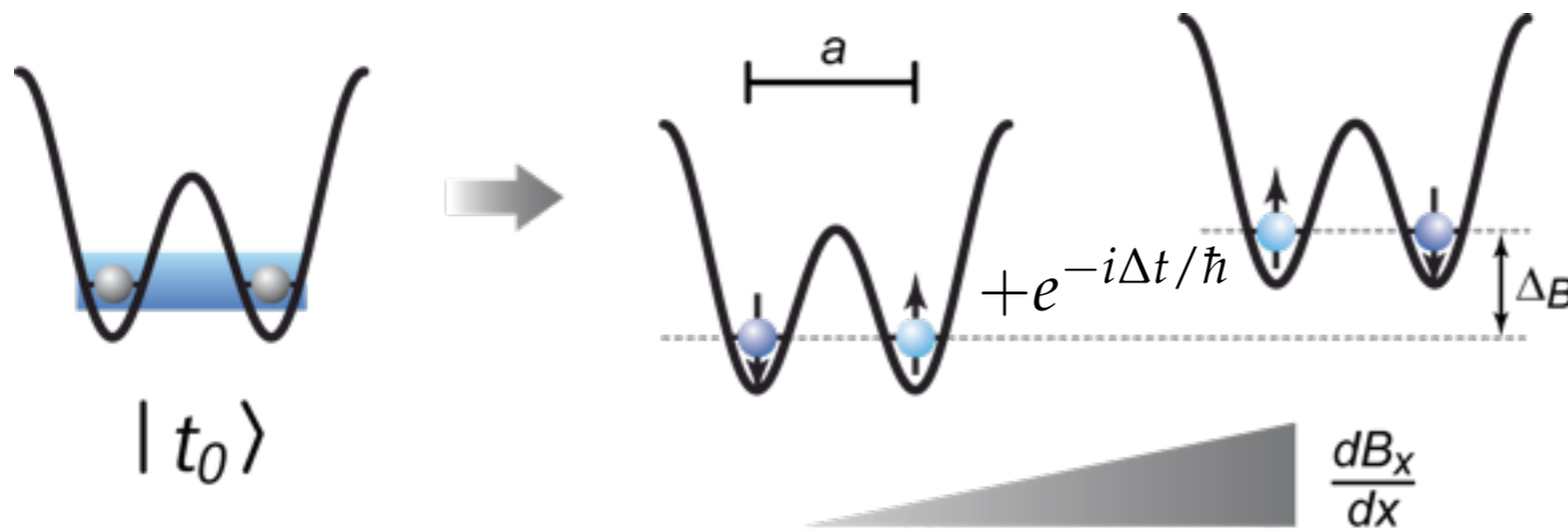
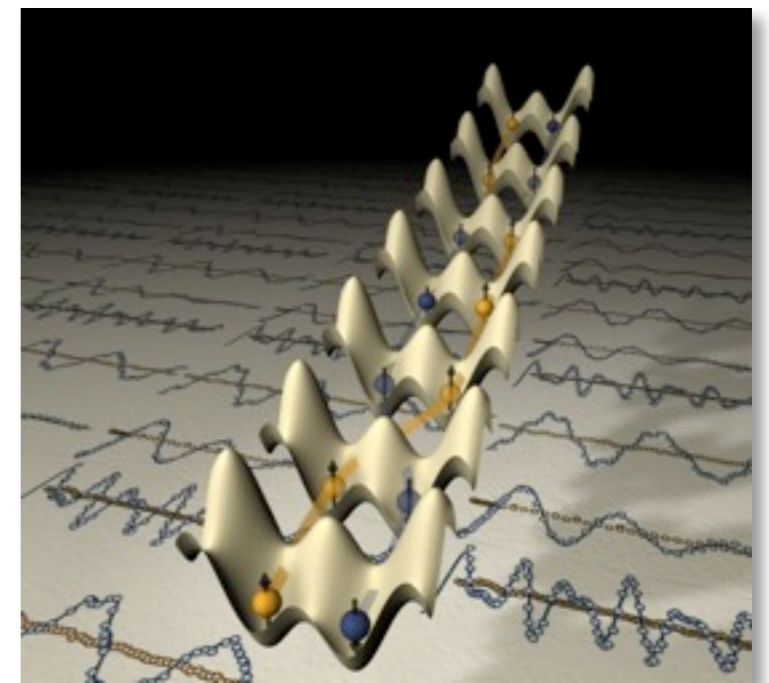
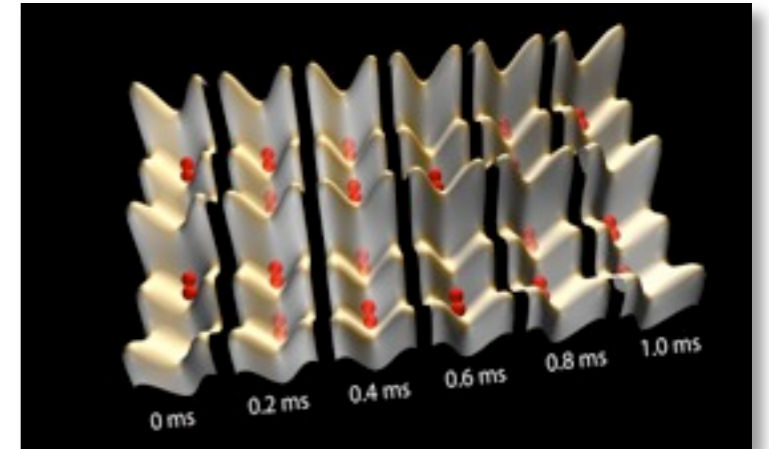
→ **Singlet:**



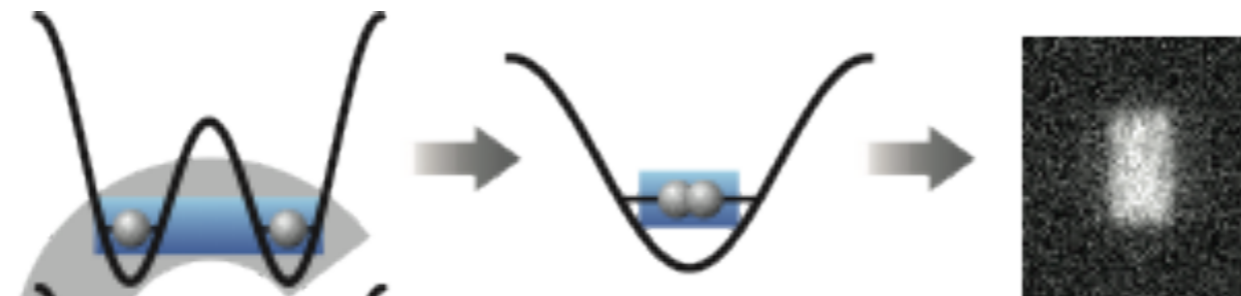
Experiments in the Superlattice

• Isolated double-wells:

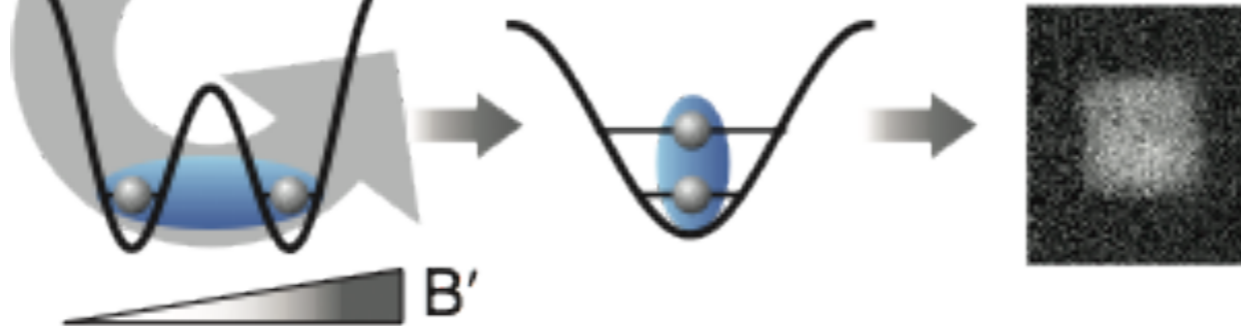
- Correlated tunneling, Superexchange interactions
- Counting atoms via interaction blockade
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→ **Triplet:**



→ **Singlet:**





Experiments in the Superlattice

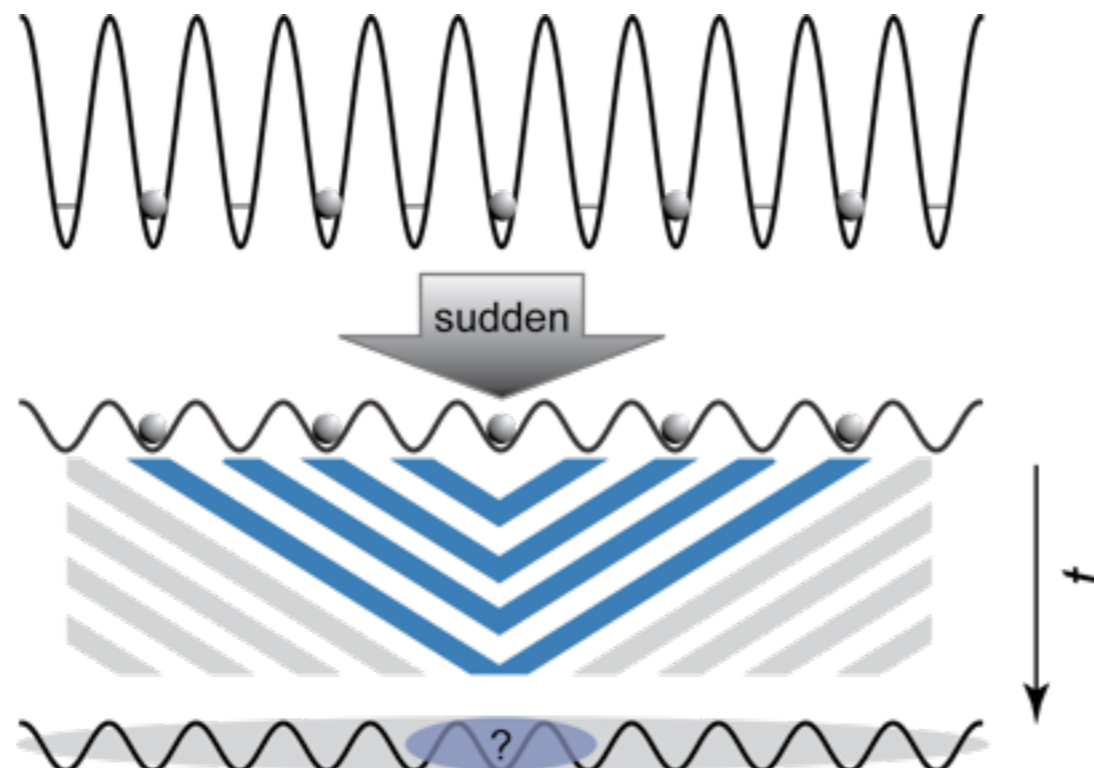
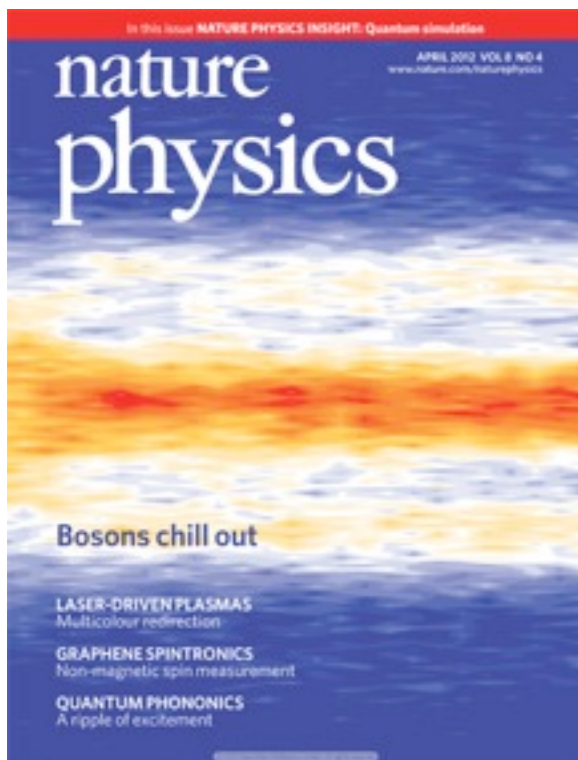
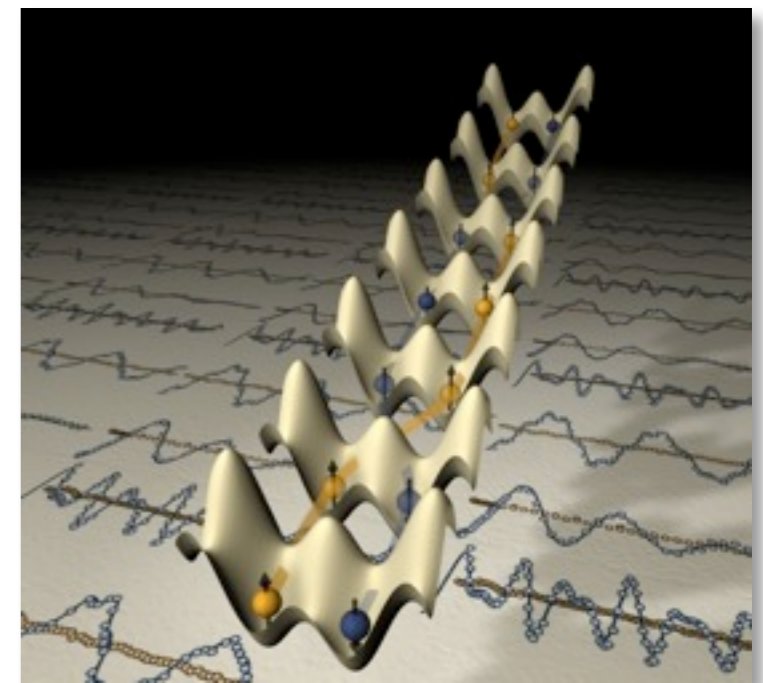
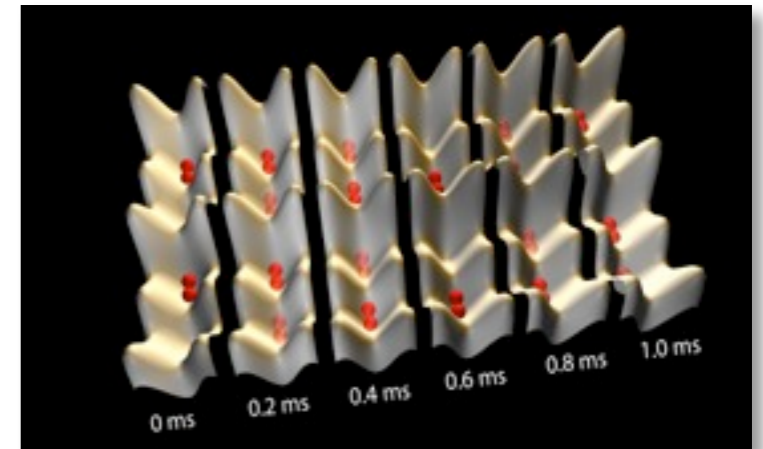
- **Isolated double-wells:**

- Correlated tunneling, Superexchange interactions
- Counting atoms via interaction blockade
- Control of n.n. spin correlations

...

- **Non-equilibrium & adiabatic dynamics:**

- Decay of patterned states (spin, density) after quantum quenches





Experiments in the Superlattice

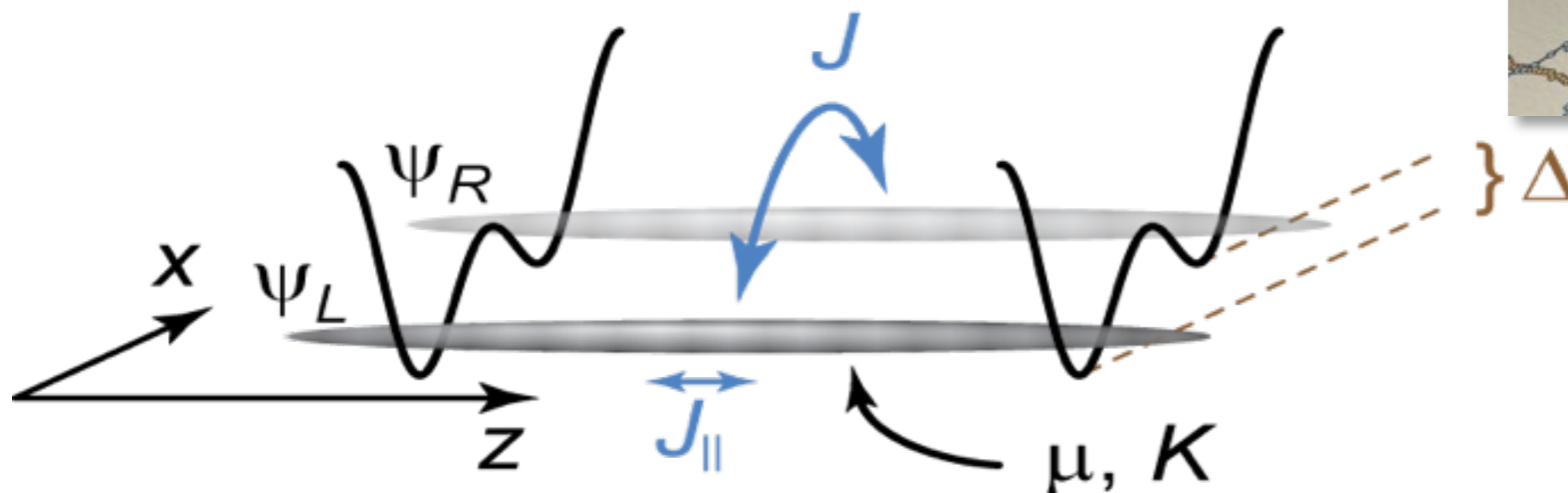
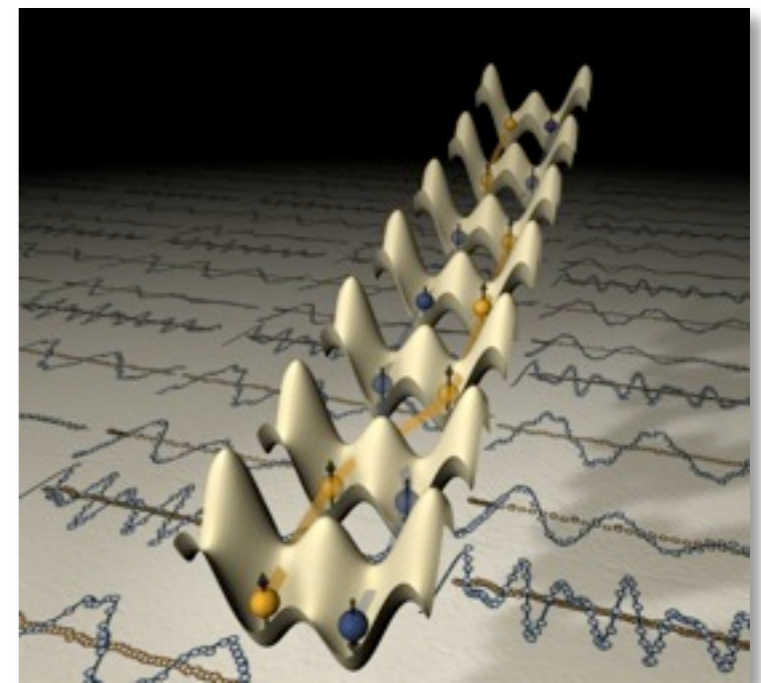
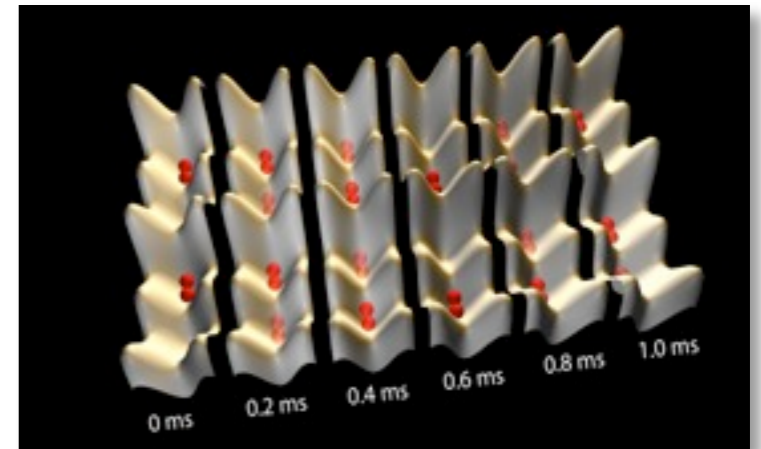
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Experiments in the Superlattice

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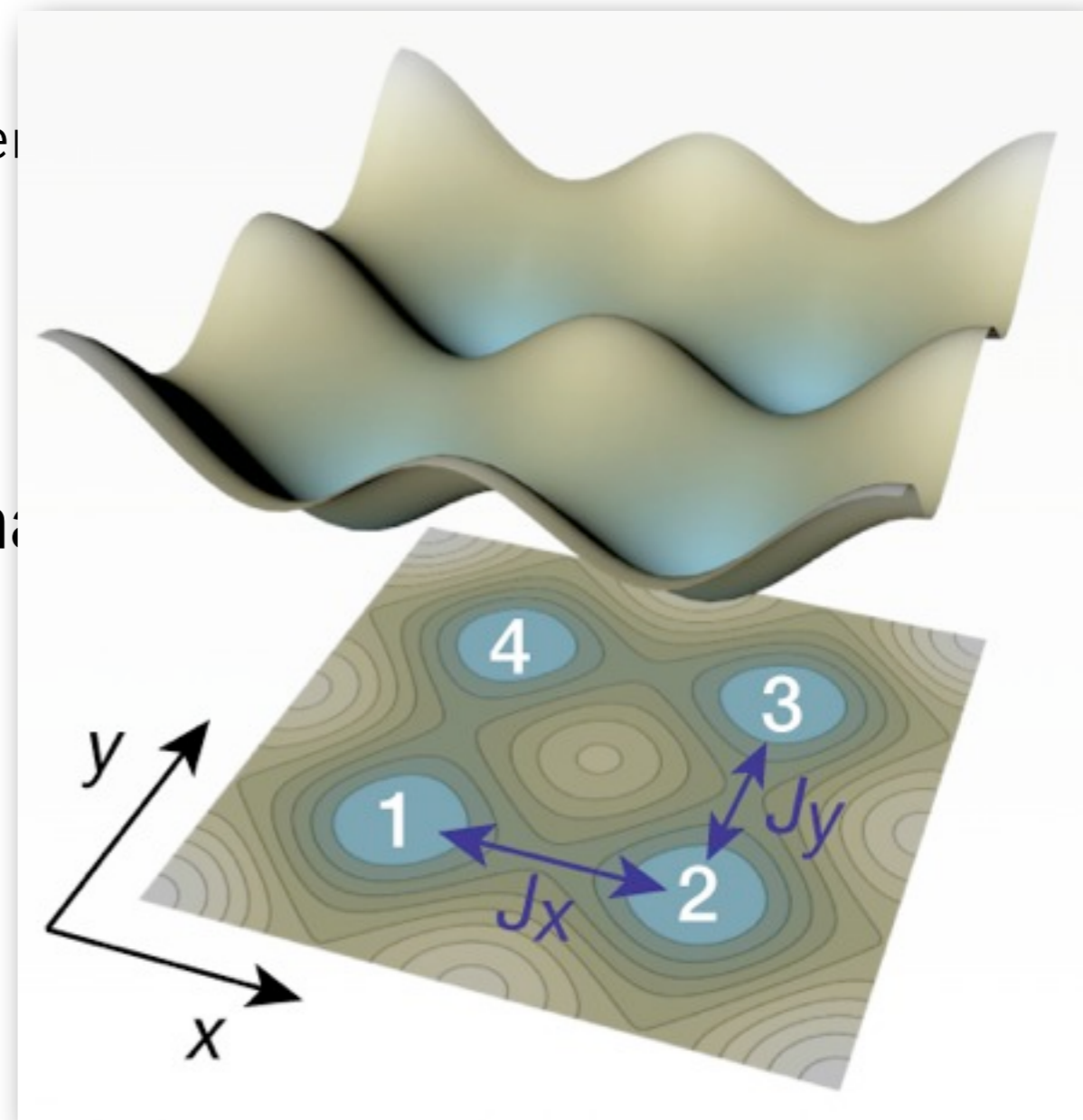
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...

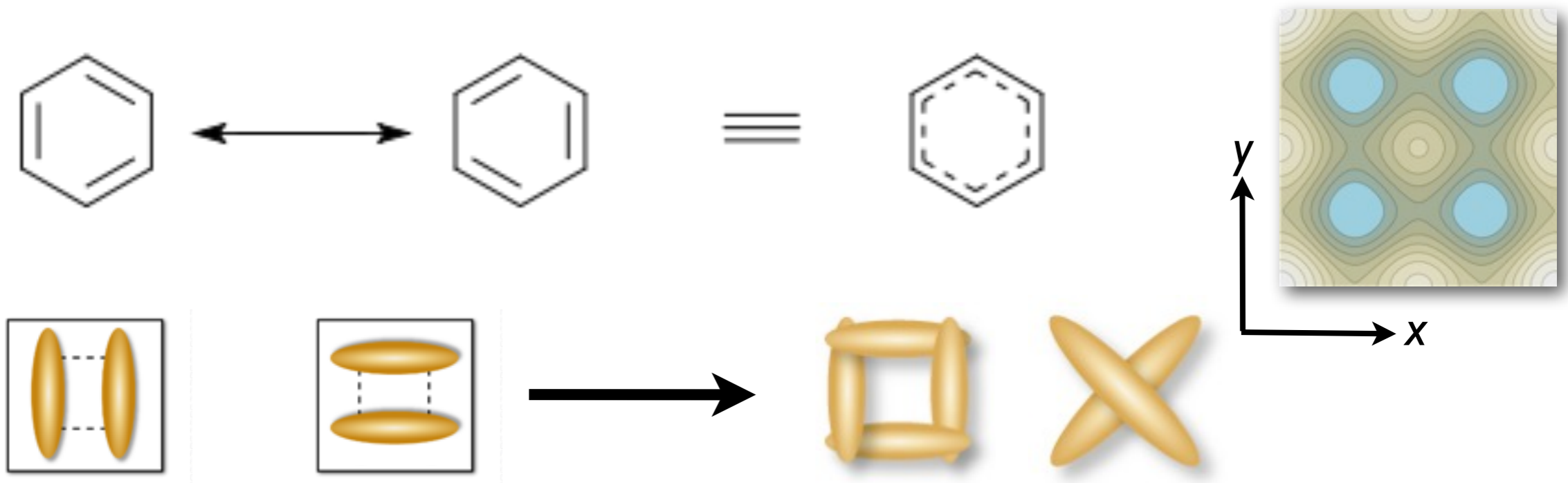
- **Isolated plaquettes:**

- Resonating Valence Bond State





Experiments in the Superlattice



- **Isolated plaquettes:**

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Experiments in the Superlattice

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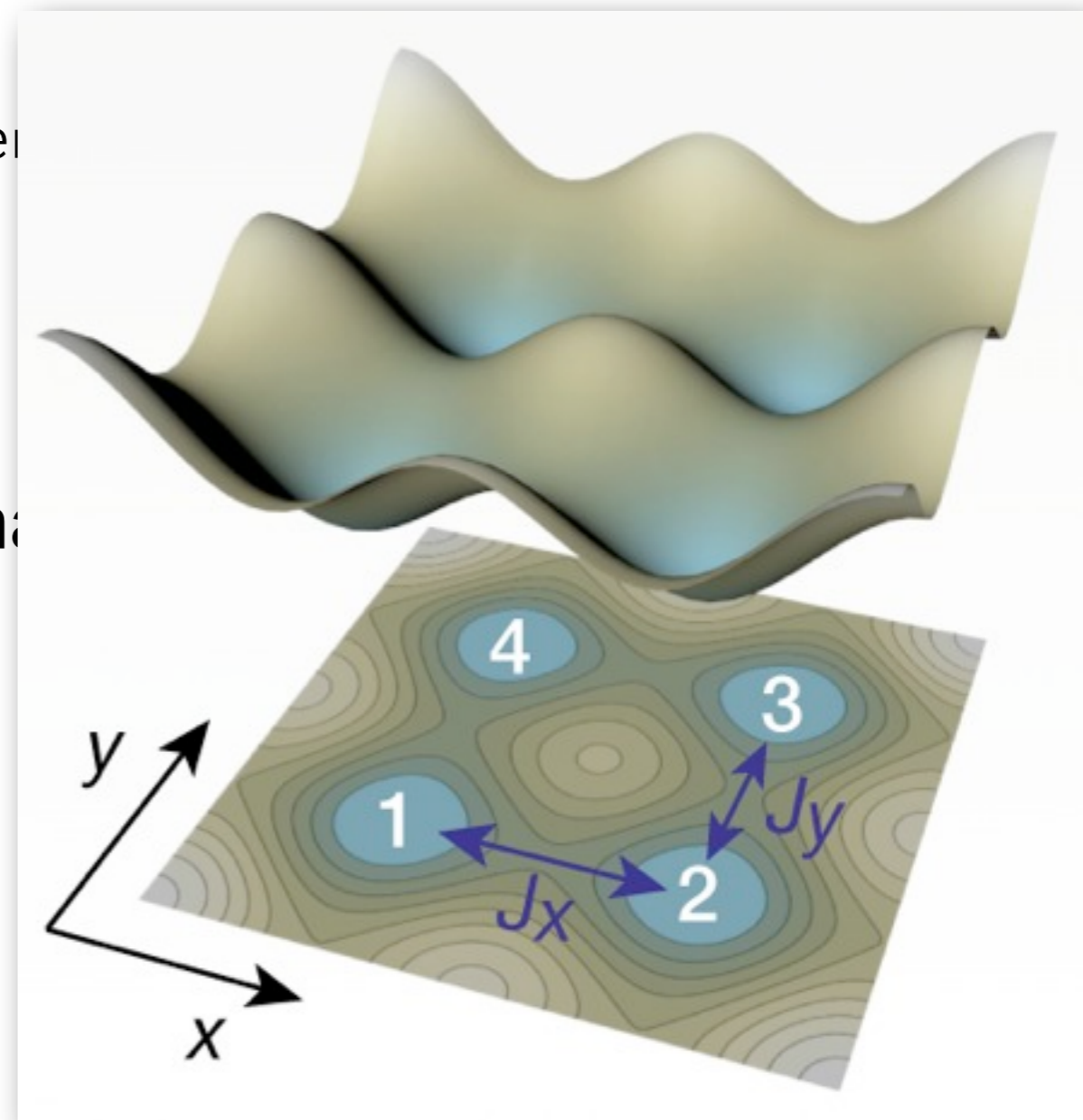
...

- **Isolated plaquettes:**

- Resonating Valence Bond State
- **Artificial Gauge Field**
- **Zak Phase in Topological Bloch Bands**

...

- **Many-body phases in the superlattice**



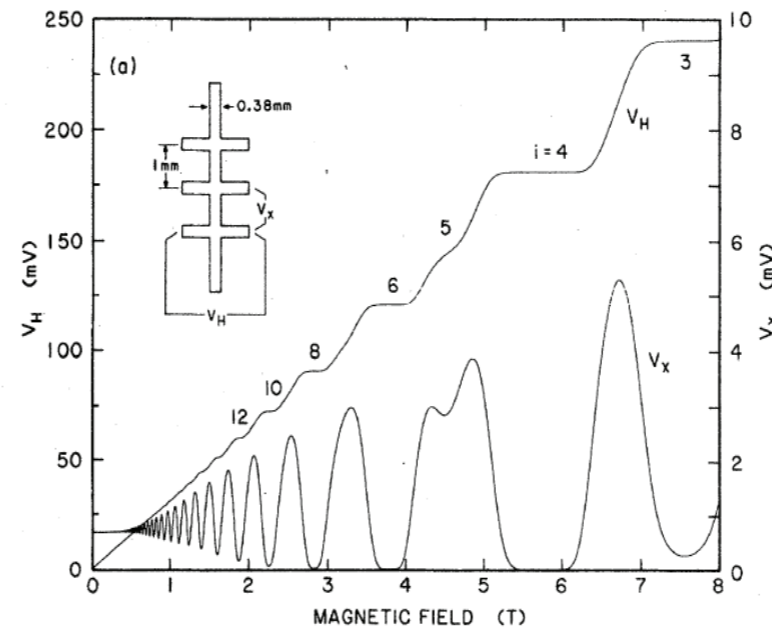
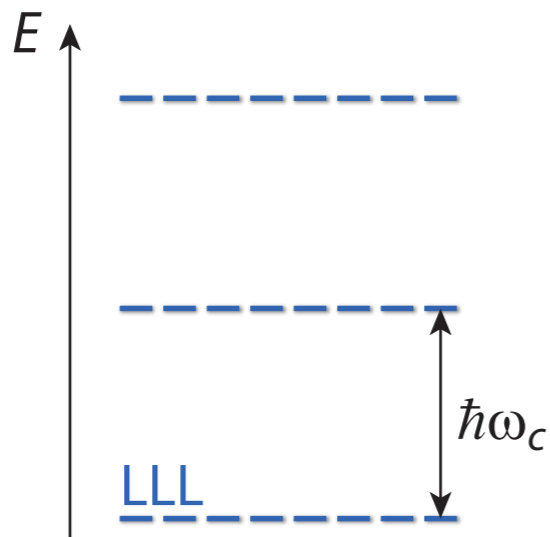
CREATION OF STRONG EFFECTIVE MAGNETIC FIELDS

**M. Aidelsburger *et al.*,
PRL 107, 255301 (2011)**



Quantum Hall effect in 2D electron gases

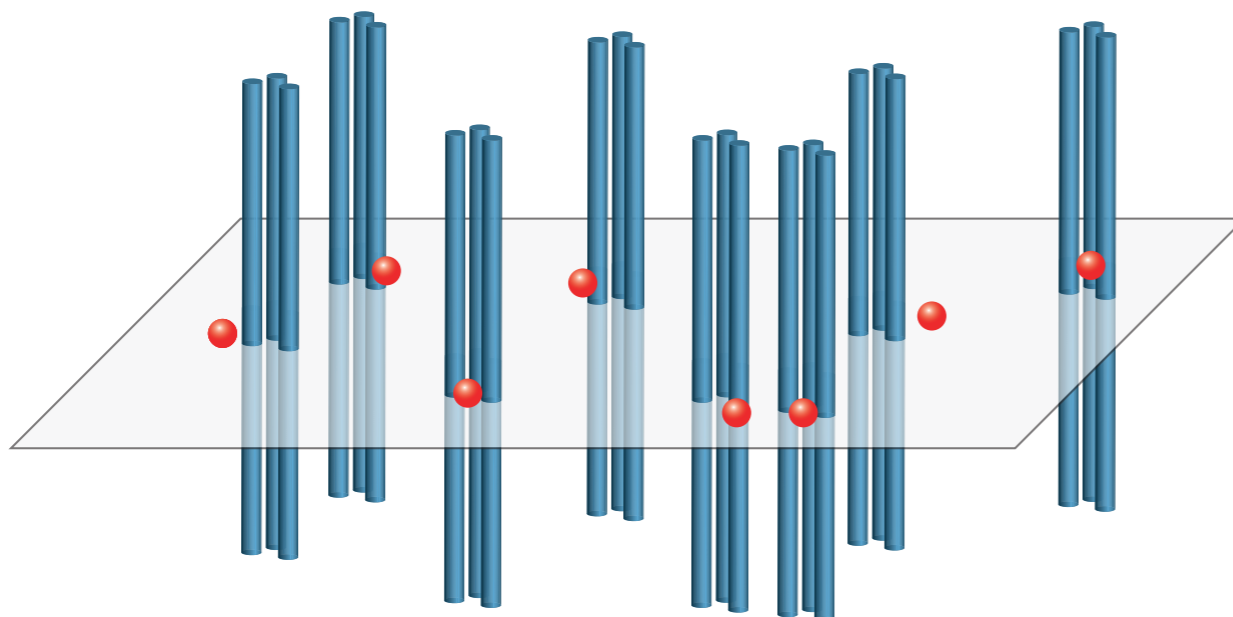
❖ Integer quantum Hall effect



$$\sigma_{xy} = \nu \frac{e^2}{h},$$

ν integer

❖ Fractional quantum Hall effect



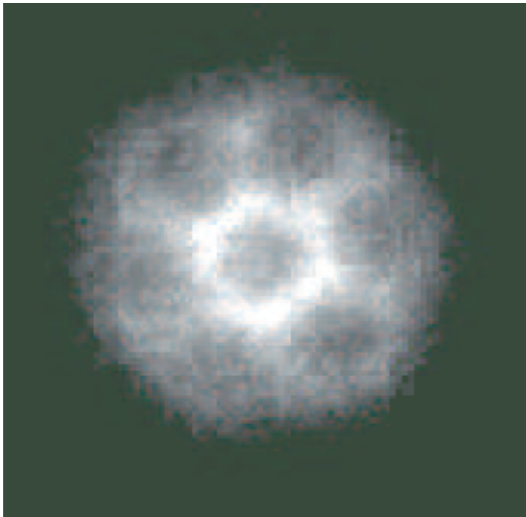
Laughlin state at $\nu = 1/3$

- flux quantum $\phi_0 = h/ec$
- electron



Artificial B fields with ultracold atoms

❖ Rotation

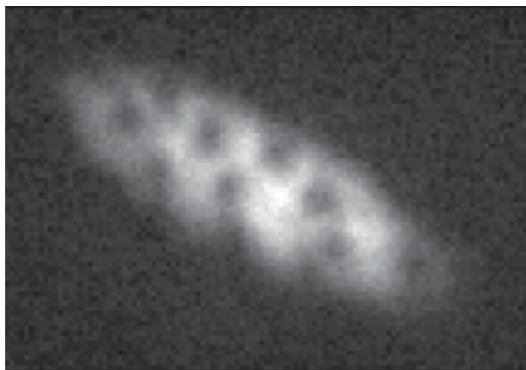


The Coriolis force $\mathbf{F}_C = 2m \mathbf{v} \times \boldsymbol{\Omega}_{\text{rot}}$ is analogous to the Lorentz force $\mathbf{F}_L = q \mathbf{v} \times \mathbf{B}$

Issue: typically $\gamma > 1000$

K. Madison *et al*, Phys. Rev. Lett. **84**, 806 (2000)
J. R. Abo-Shaeer *et al*, Science **292**, 476 (2001)

❖ Raman-induced gauge field



Spatially dependent optical couplings lead to a Berry phase analogous to the Aharonov-Bohm phase.

Issues: small B fields, heating from Raman lasers.

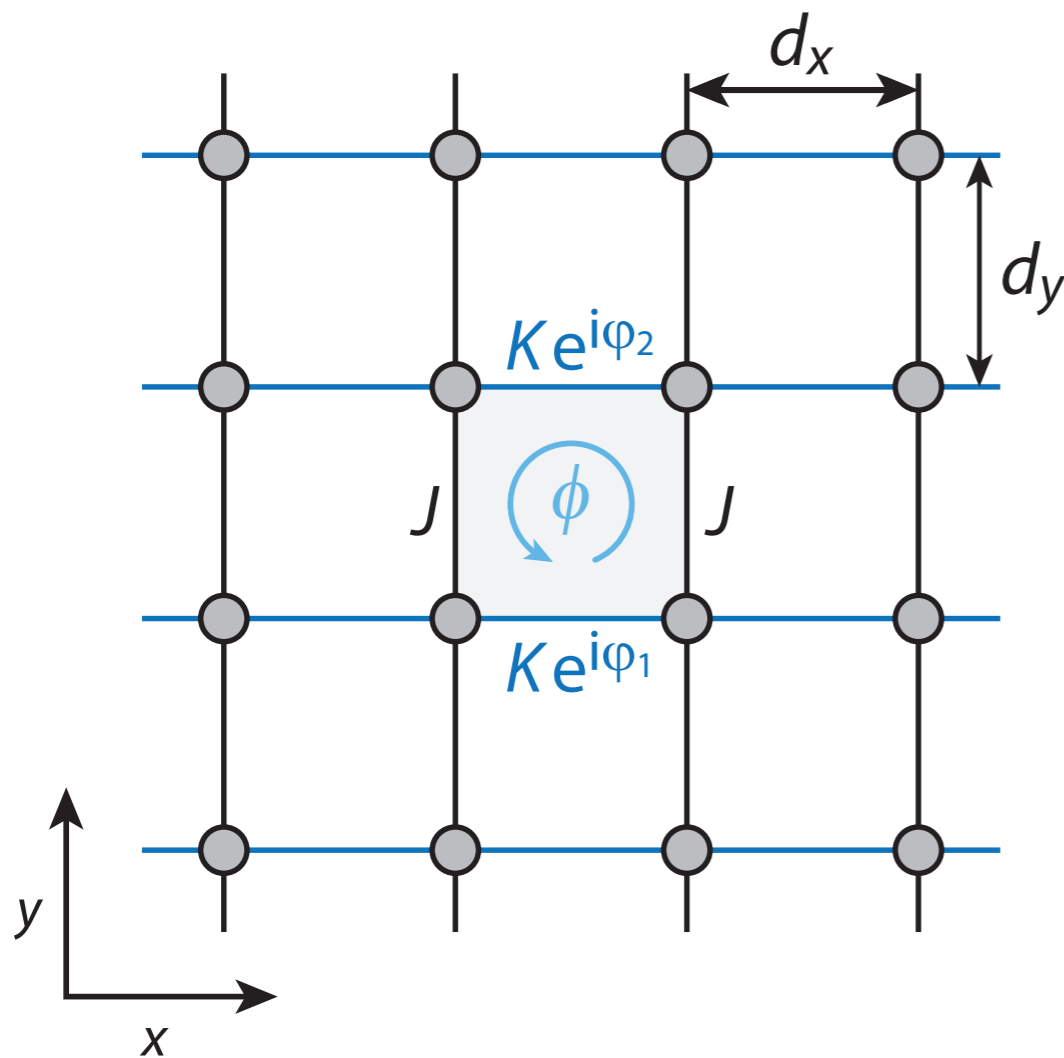
Y. Lin *et al*, Nature **462**, 628 (2009)



Artificial B fields with ultracold atoms in OLs

Controlling atom tunneling along x with Raman lasers leads to effective tunnel couplings with spatially-dependent Peierls phases $\phi(\mathbf{R})$

$$\hat{H} = - \sum_{\mathbf{R}} \left(K e^{i\varphi(\mathbf{R})} \hat{a}_{\mathbf{R}}^\dagger \hat{a}_{\mathbf{R}+\mathbf{d}_x} + J \hat{a}_{\mathbf{R}}^\dagger \hat{a}_{\mathbf{R}+\mathbf{d}_y} \right) + \text{h.c.}$$



Magnetic flux through a plaquette

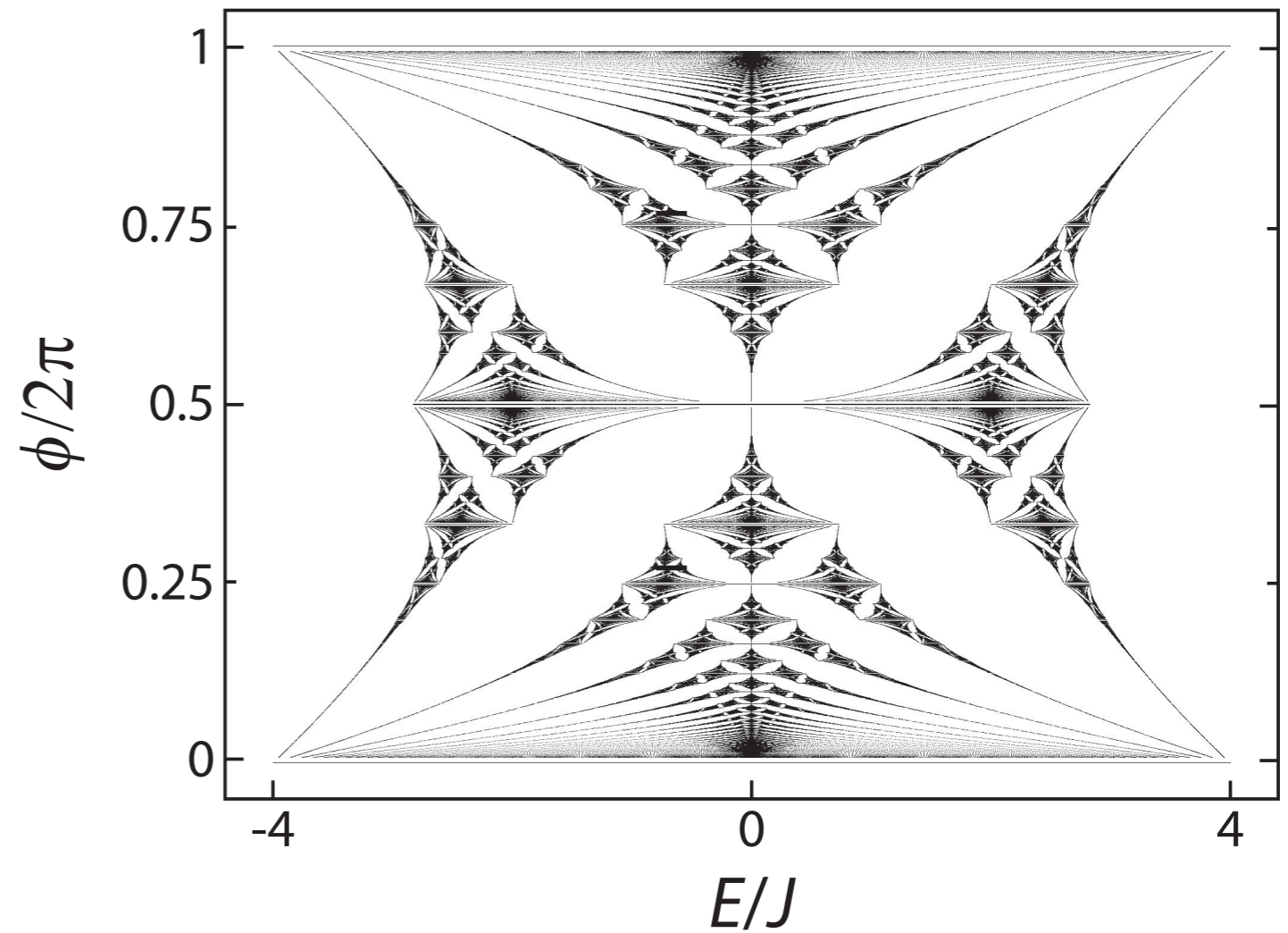
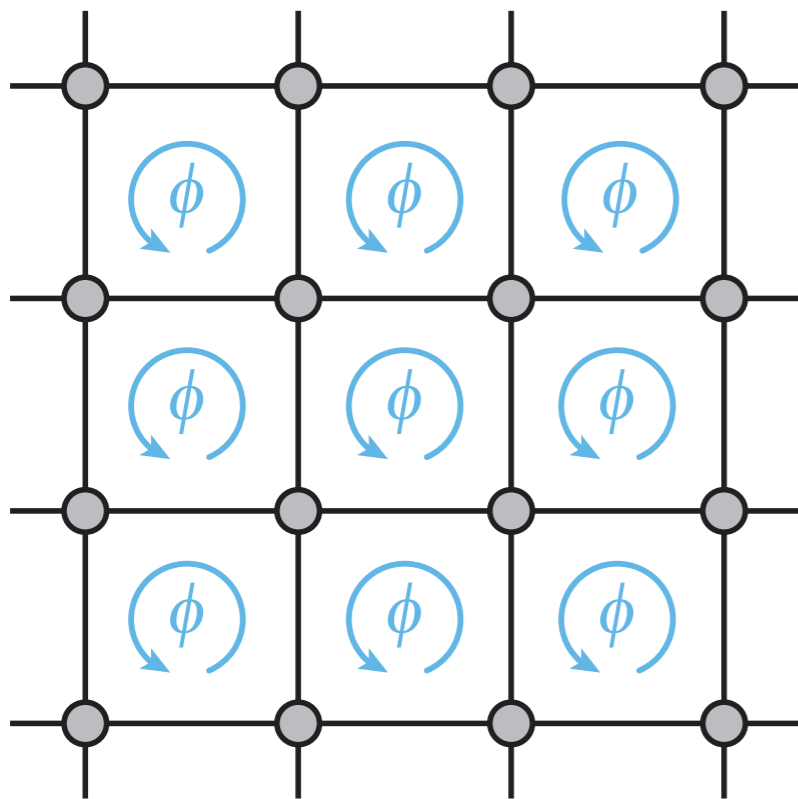
$$\phi = \int_{\square} B dS = \varphi_1 - \varphi_2$$

- D. Jaksch & P. Zoller, NJP **5**, 56 (2003)
- F. Gerbier & J. Dalibard, NJP **12**, 033007 (2010)
- E. Mueller, PPA **70**, 041603 (2004)
- A. Kolovsky, Europhys. Lett. **93**, 20003 (2011)



Harper Hamiltonian and Hofstadter butterfly

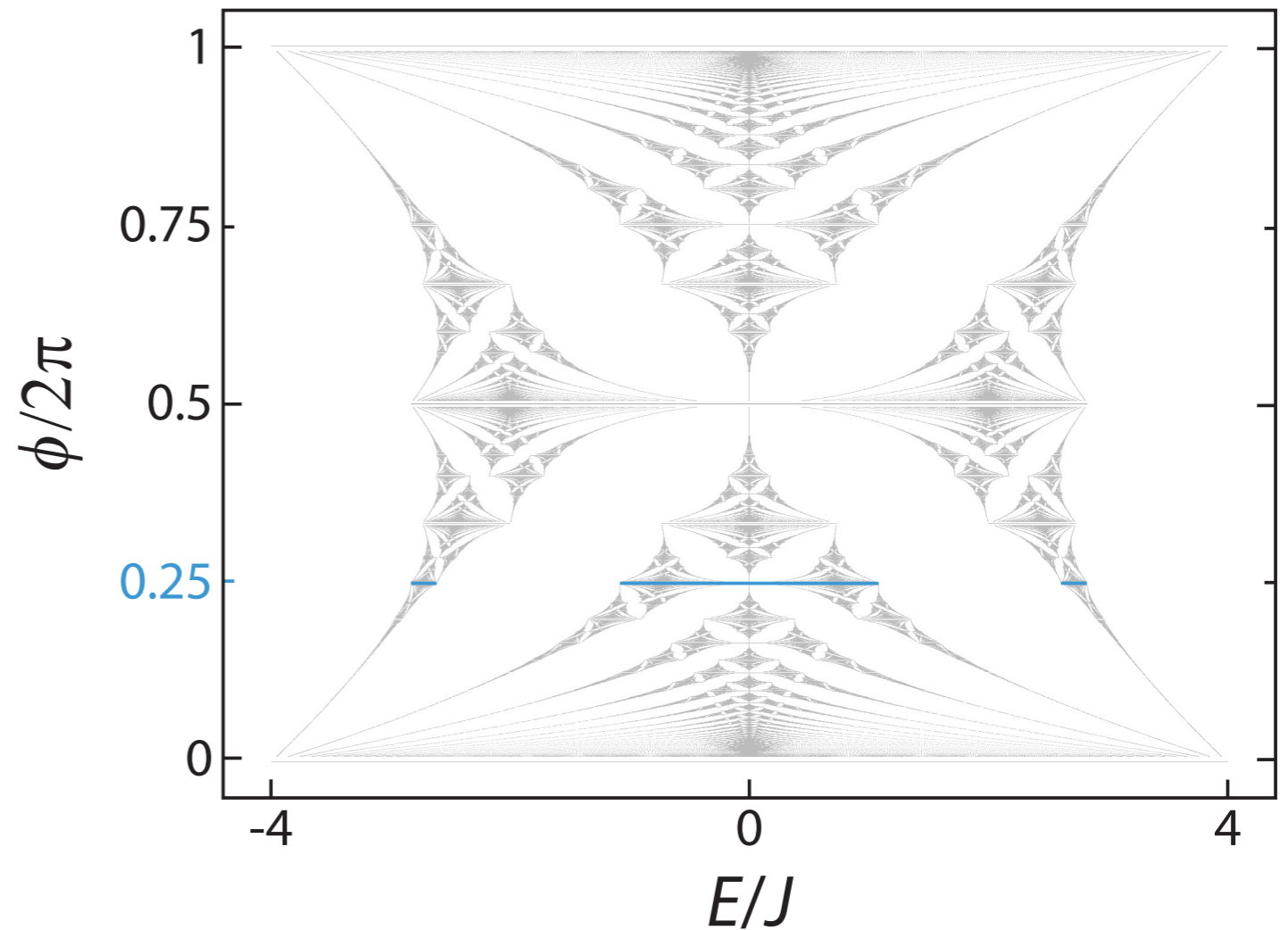
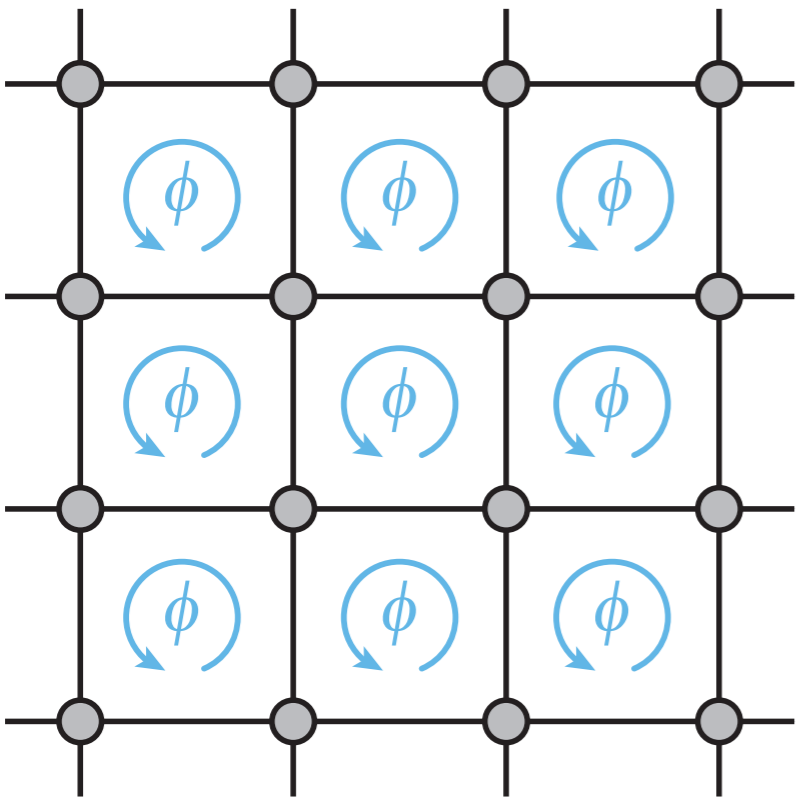
Harper Hamiltonian: $J=K$ and ϕ uniform.





Harper Hamiltonian and Hofstadter butterfly

Harper Hamiltonian: $J=K$ and ϕ uniform.

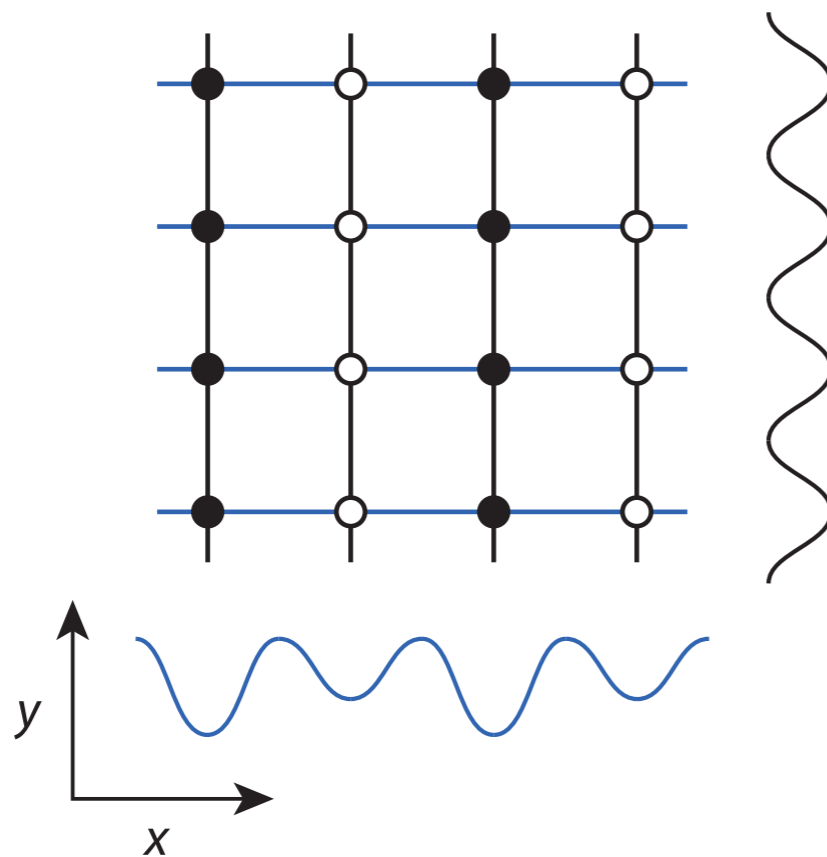


- Lowest band is topologically equivalent to lowest Landau level
- $\nu=1/2$ + repulsive interactions \longrightarrow Laughlin state for Bosons.



Staggered flux lattice with Rb atoms

Consider a 2D optical lattice, where tunneling is inhibited along the x direction by a superlattice potential



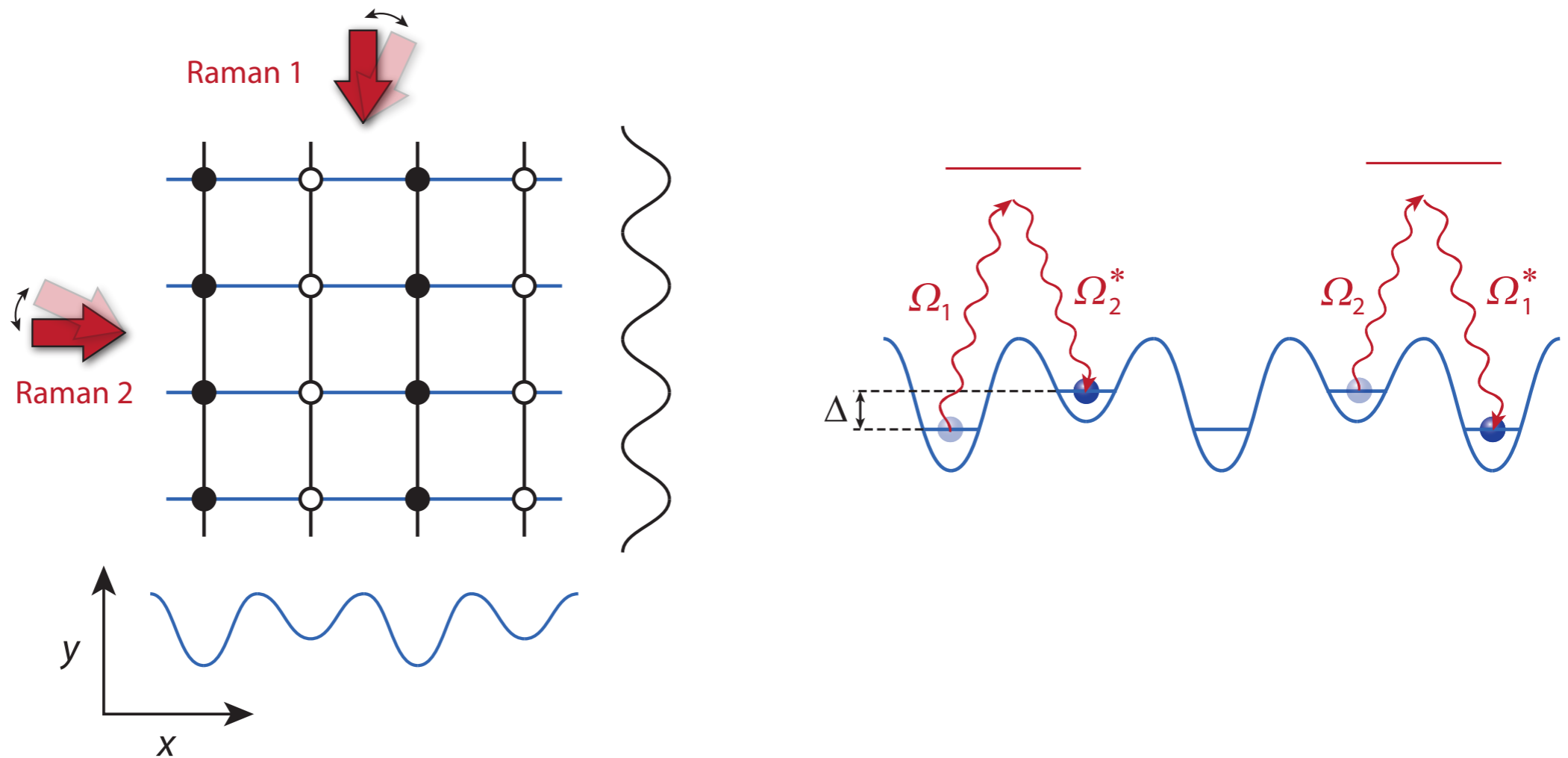
D. Jaksch & P. Zoller, NJP **5**, 56 (2003)
F. Gerbier & J. Dalibard, NJP **12**, 033007 (2010)
A. Kolovsky, Europhys. Lett. **93**, 20003 (2011)





Staggered flux lattice with Rb atoms

Tunneling along this direction can be restored using Raman beams.

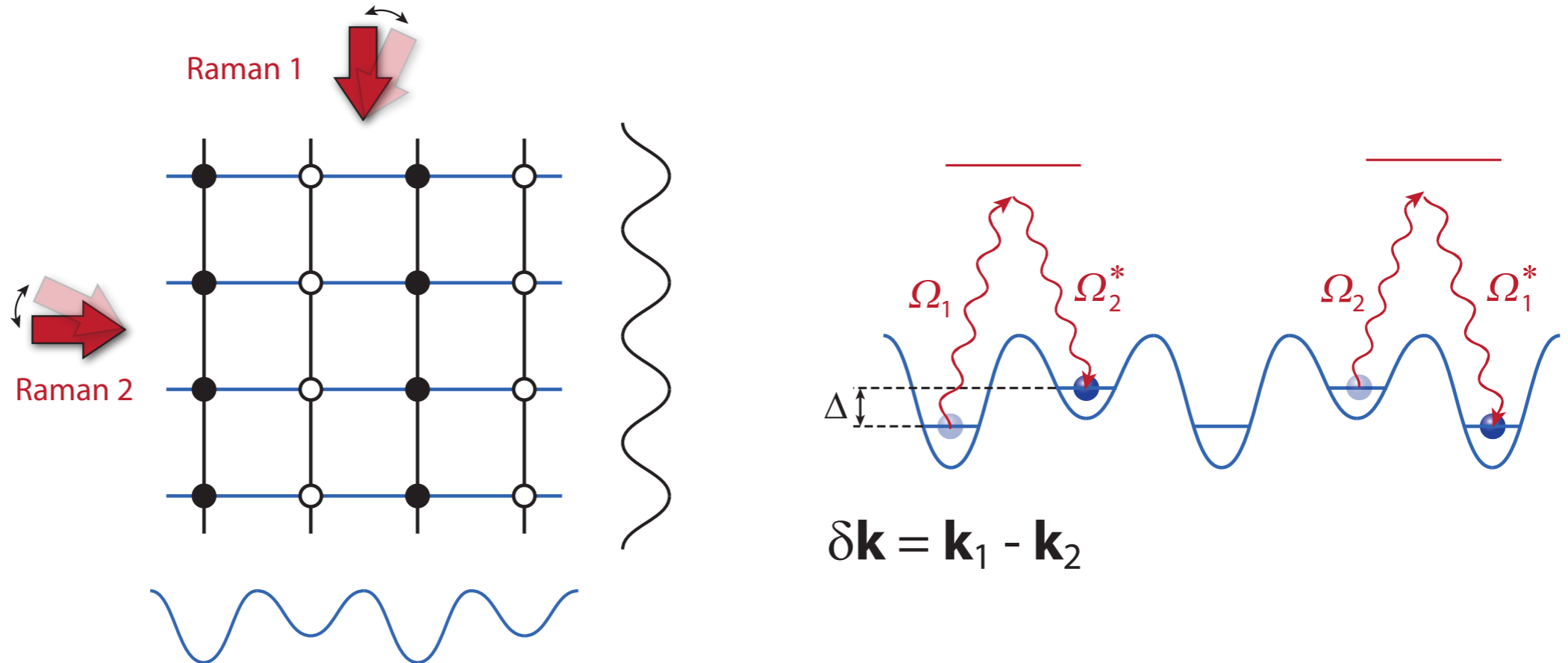


D. Jaksch & P. Zoller, NJP **5**, 56 (2003)
F. Gerbier & J. Dalibard, NJP **12**, 033007 (2010)
A. Kolovsky, Europhys. Lett. **93**, 20003 (2011)





Staggered flux lattice with Rb atoms



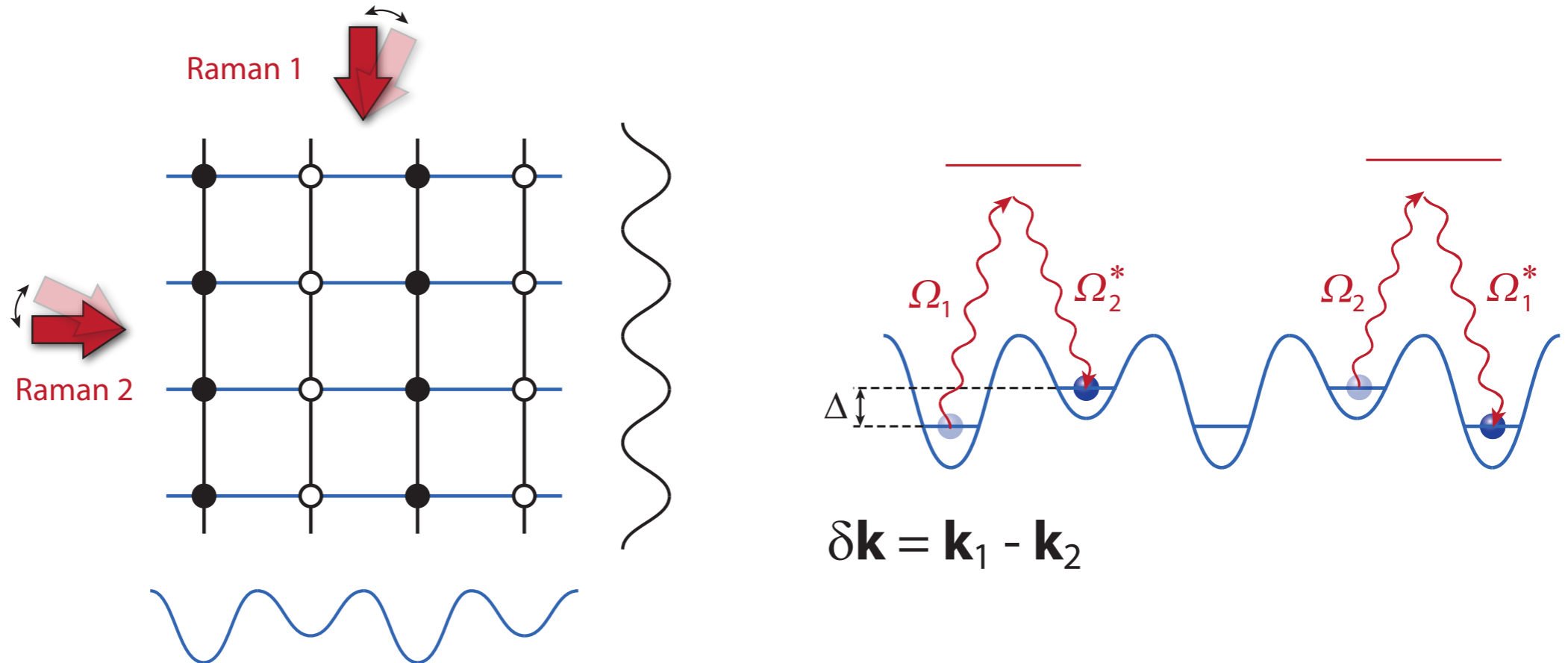
$$\begin{aligned}
 K_{|\bullet\rangle \rightarrow |\circ\rangle}(\mathbf{R}) &= \int d\mathbf{r} w_{\bullet}^*(\mathbf{r} - \mathbf{R}) w_{\circ}(\mathbf{r} - \mathbf{R} - \mathbf{d}_x) \Omega(\mathbf{r}) \\
 &= K e^{i\delta\mathbf{k} \cdot \mathbf{R}} \quad \text{for} \quad \Omega(\mathbf{r}) = V_K e^{i\delta\mathbf{k} \cdot \mathbf{r}}
 \end{aligned}$$

D. Jaksch & P. Zoller, *NJP* **5**, 56 (2003)
 F. Gerbier & J. Dalibard, *NJP* **12**, 033007 (2010)
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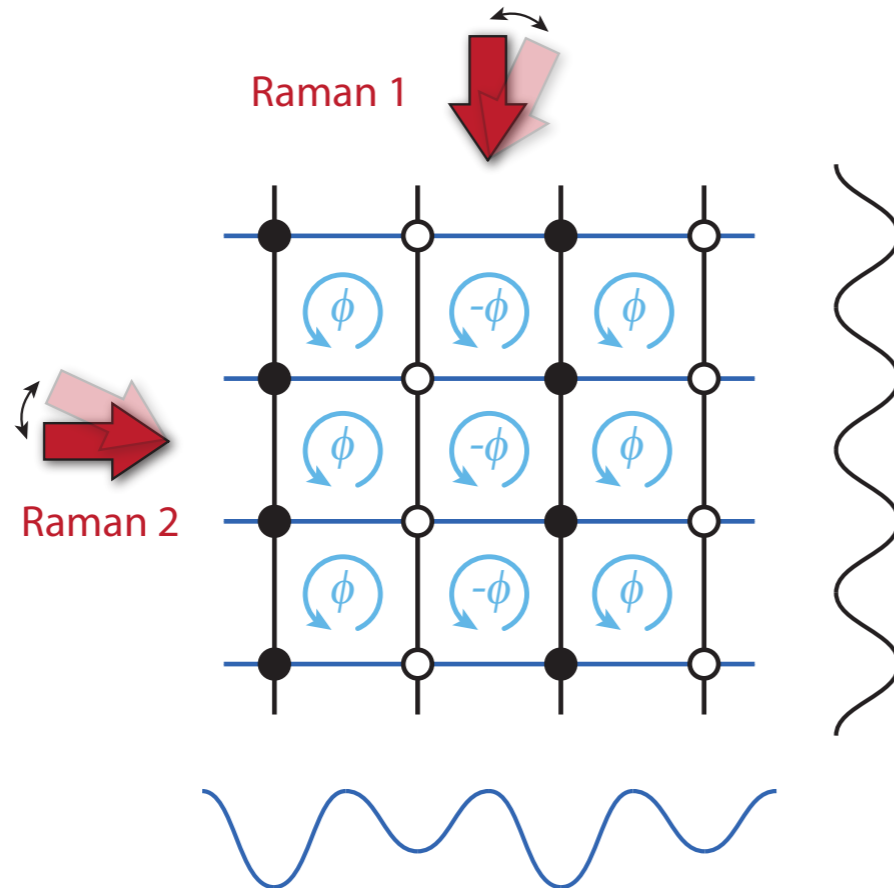
$$K_{|\circ\rangle \rightarrow |\bullet\rangle}(\mathbf{R}') = K e^{-i\delta\mathbf{k} \cdot \mathbf{R}'}$$

- D. Jaksch & P. Zoller, *NJP* **5**, 56 (2003)
 F. Gerbier & J. Dalibard, *NJP* **12**, 033007 (2010)
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Staggered flux lattice with Rb atoms



- Staggered flux ϕ with zero mean
- Tunable flux value, $\delta\mathbf{k} = \mathbf{k}_1 - \mathbf{k}_2$ (our setup: $\phi = \pi/2$)

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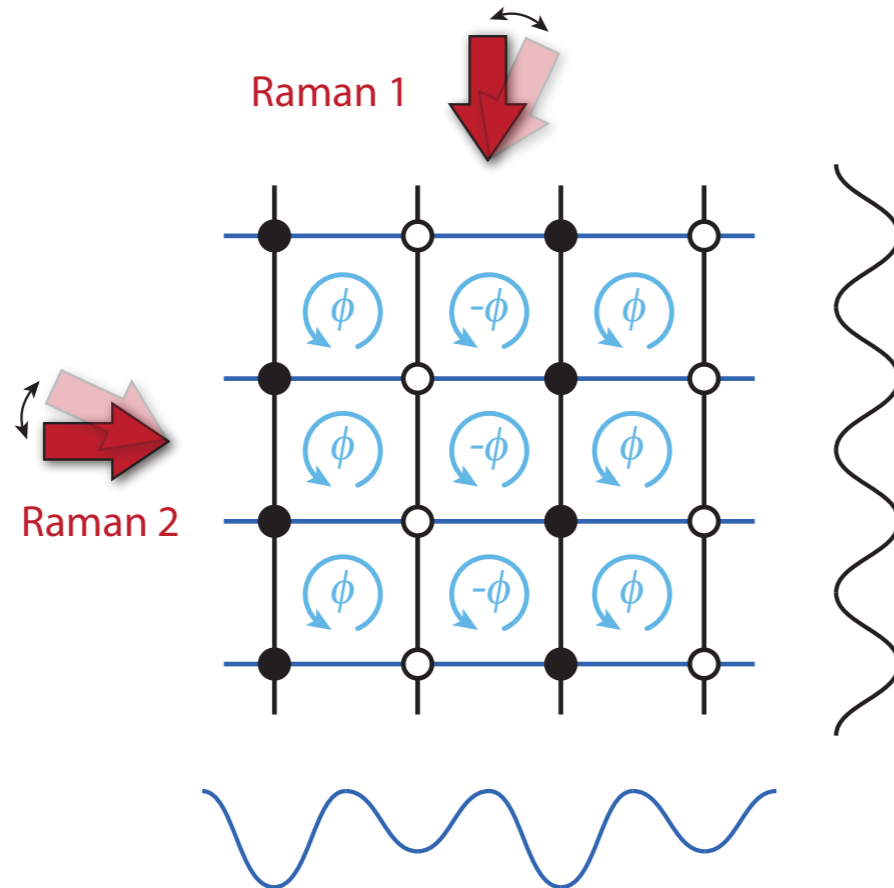
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Methods to rectify the flux:

- Linear potential gradient
- State-dependent lattices

D. Jaksch & P. Zoller, NJP **5**, 56 (2003)
 F. Gerbier & J. Dalibard, NJP **12**, 033007 (2010)
 A. Kolovsky, Europhys. Lett. **93**, 20003 (2011)

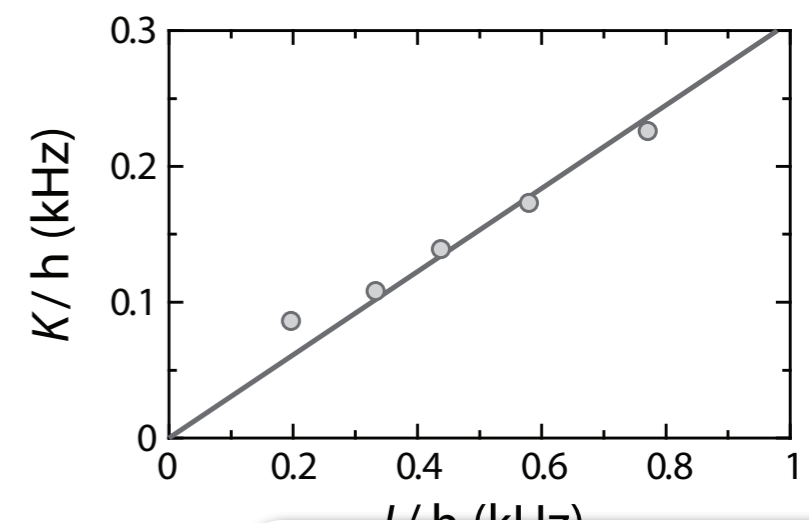
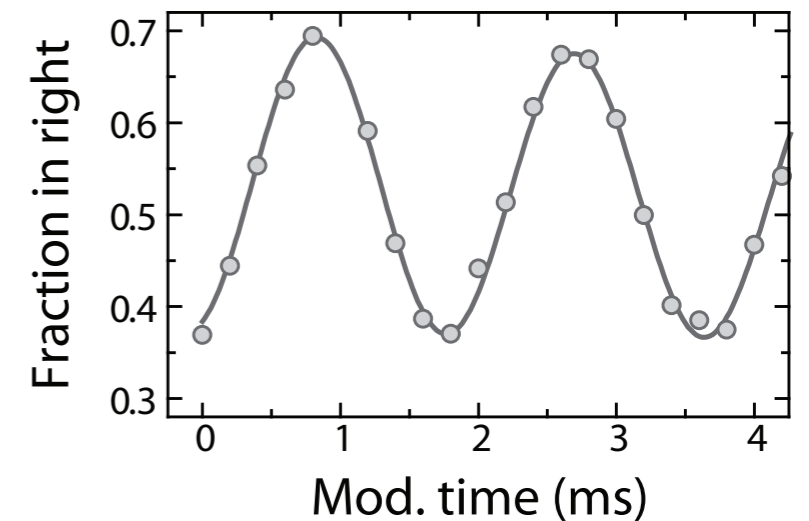
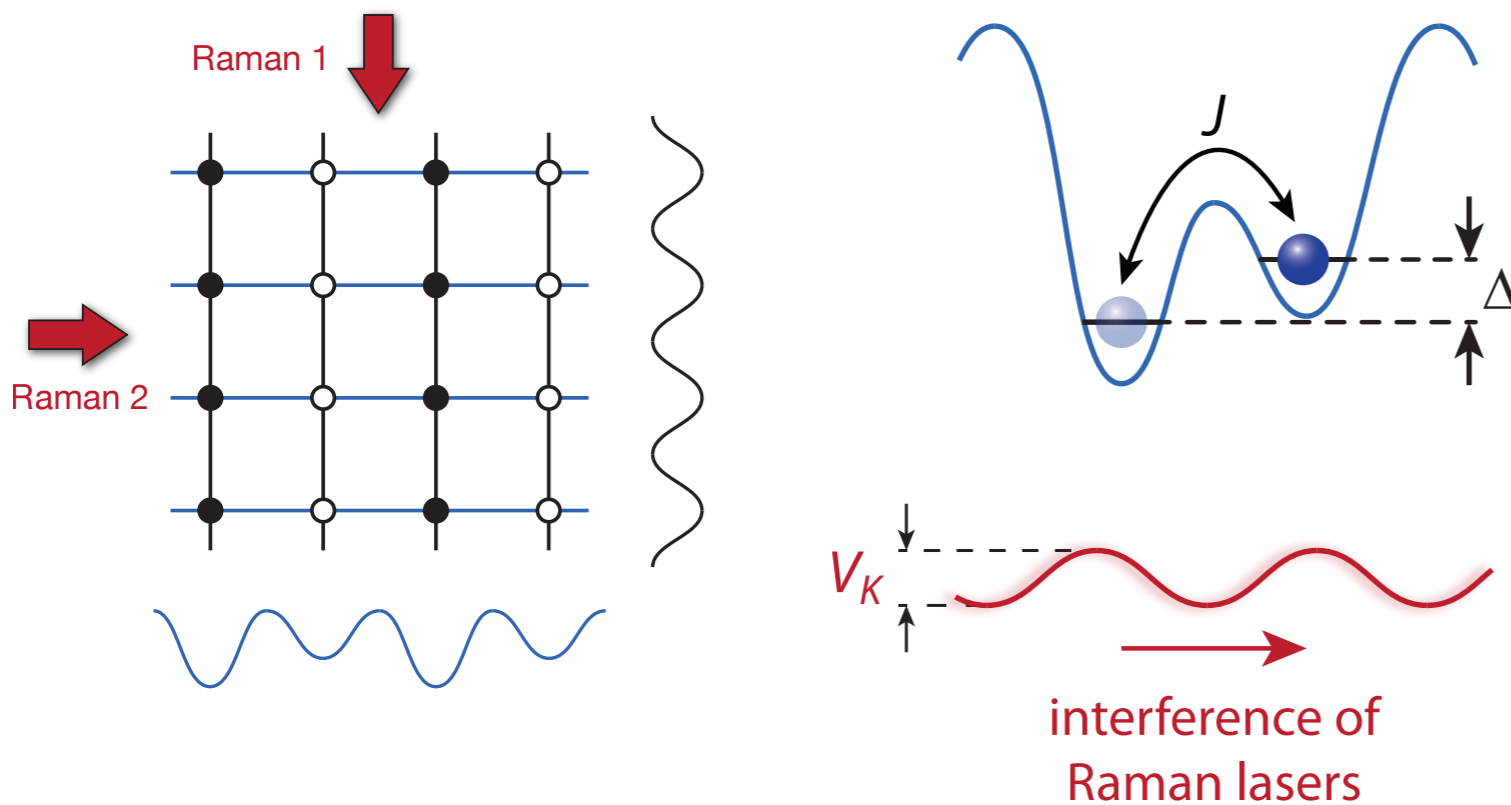




Raman-assisted tunneling

In the limit of $V_K \ll \Delta$ the amplitude of the Raman-assisted tunneling is given by:

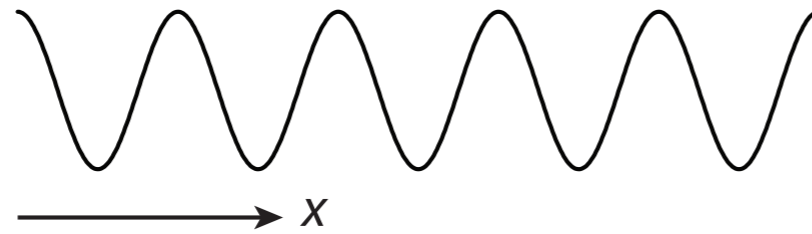
$$K \simeq \frac{1}{2\sqrt{2}} \frac{V_K J}{\Delta}$$





Experimental sequence

- ❖ We load a ^{87}Rb condensate into a 2D-optical lattice.

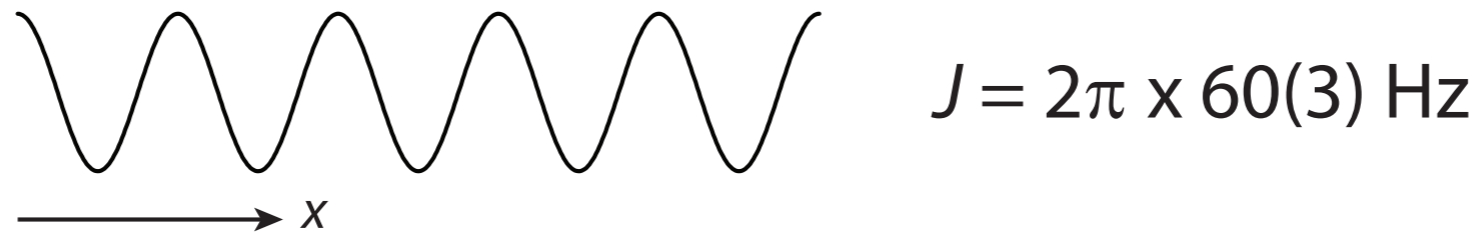


$$J = 2\pi \times 60(3) \text{ Hz}$$

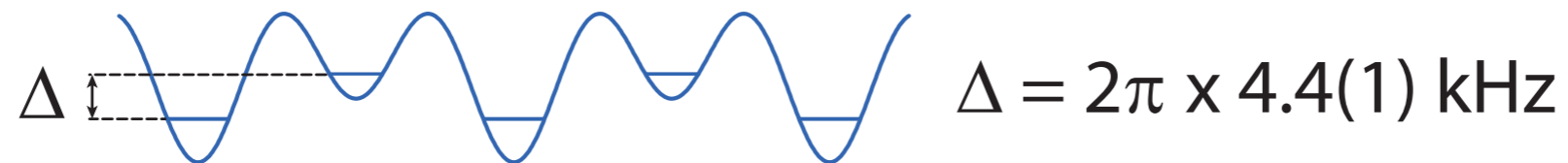


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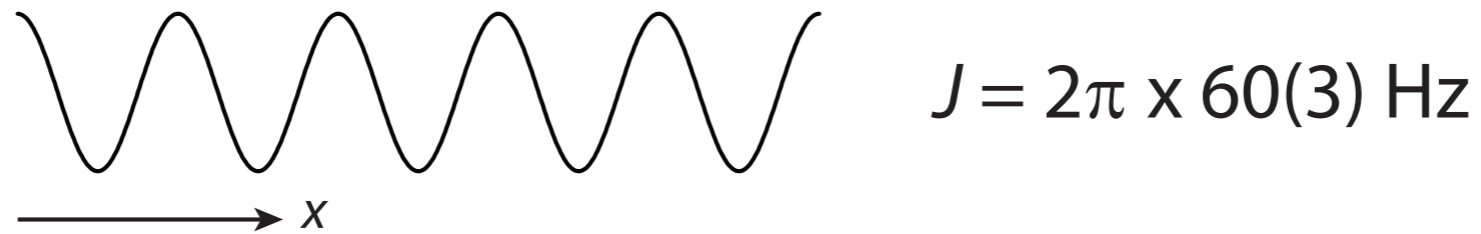
- ❖ We inhibit tunneling along x with a superlattice.



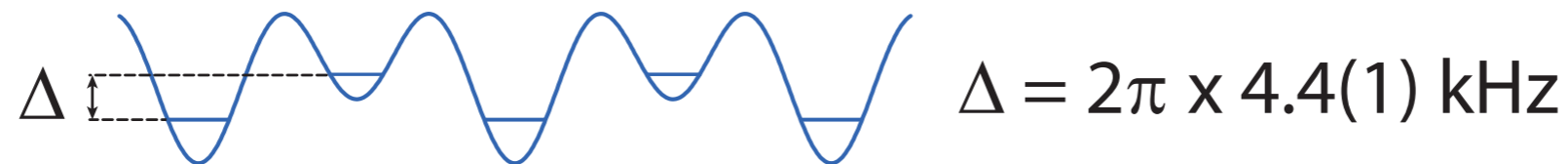


Experimental sequence

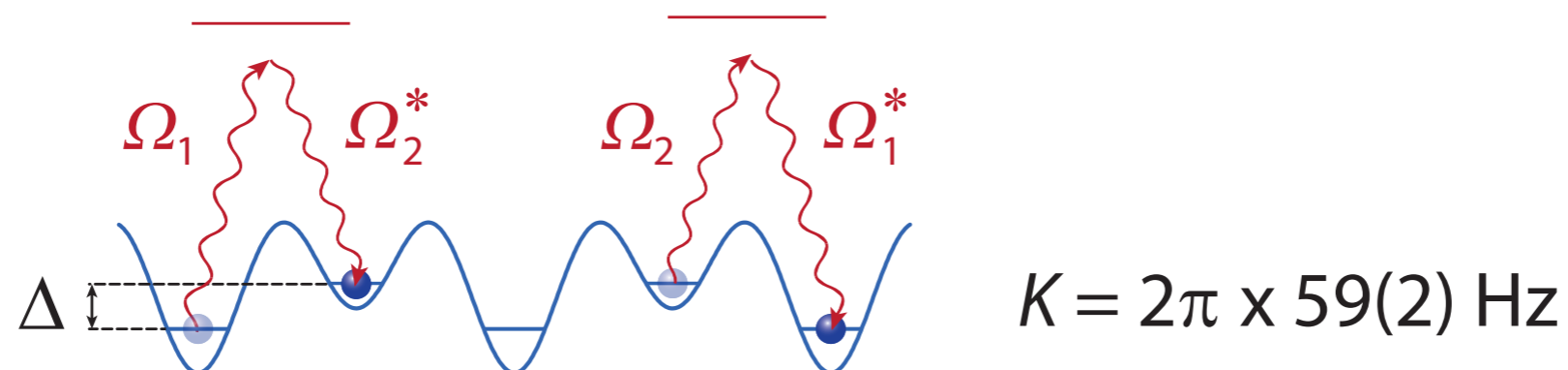
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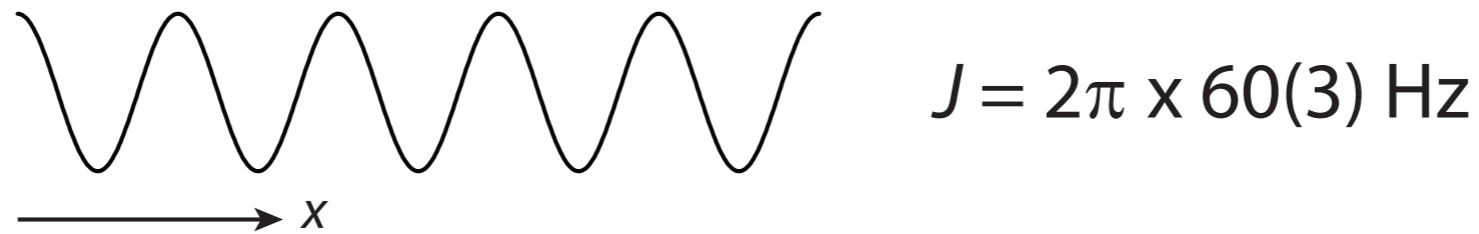
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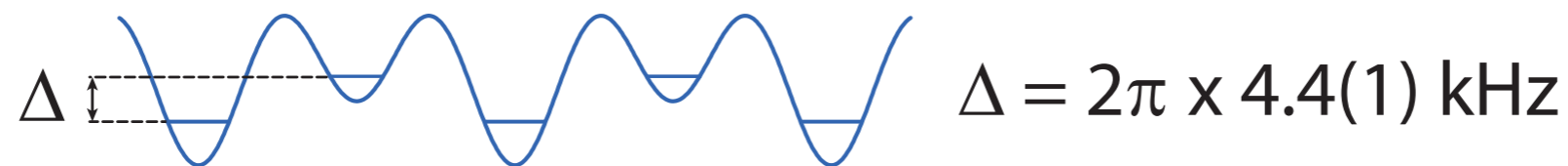


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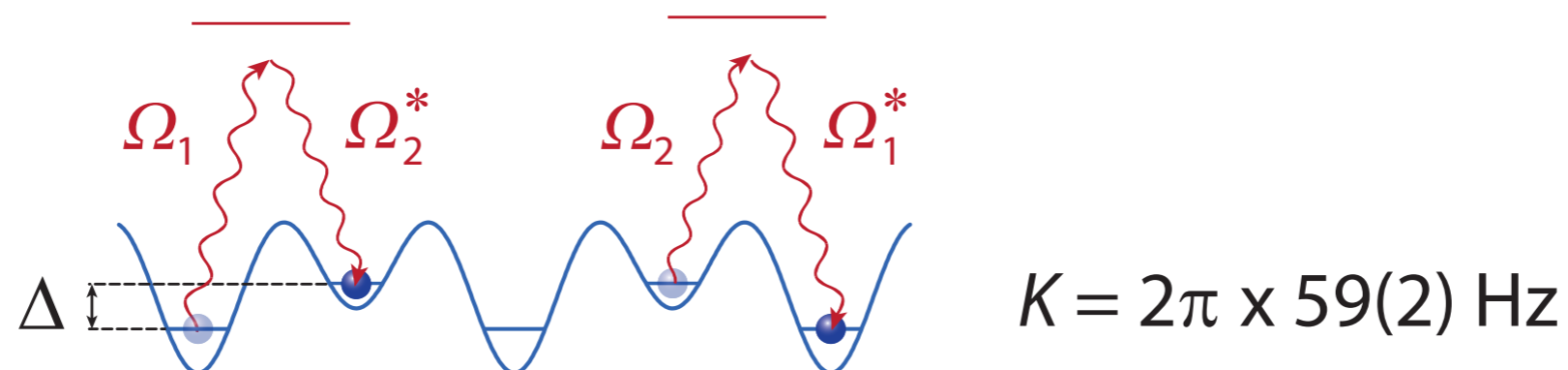
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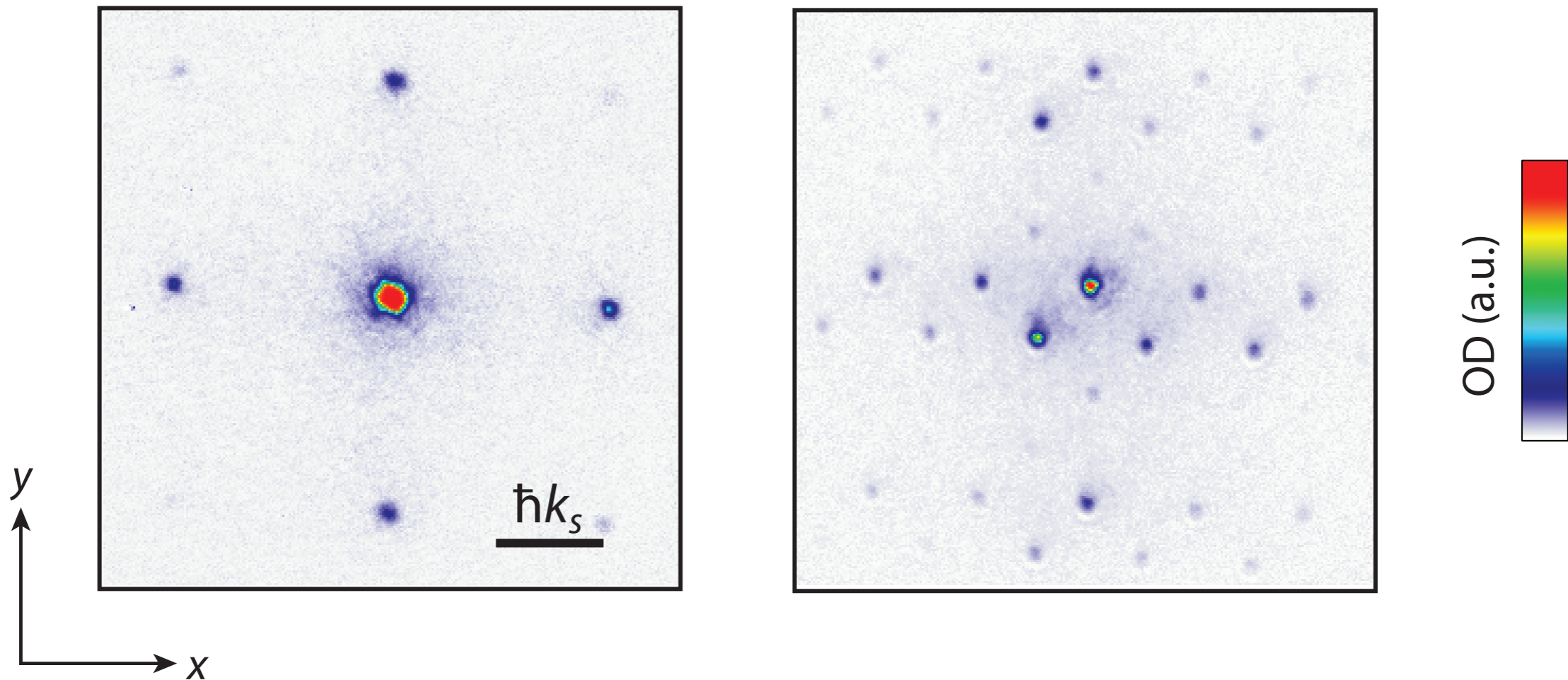
- ❖ After 10 ms hold time, TOF images.



Momentum distribution ($J/K=1$): observations

Reference: cubic lattice
(no Raman drive)

$J/K=1.0(1)$

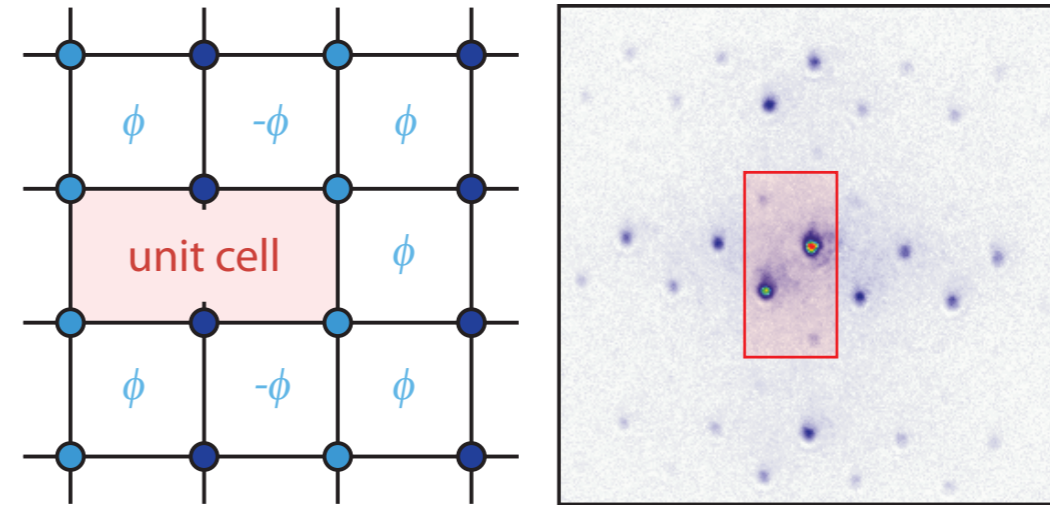
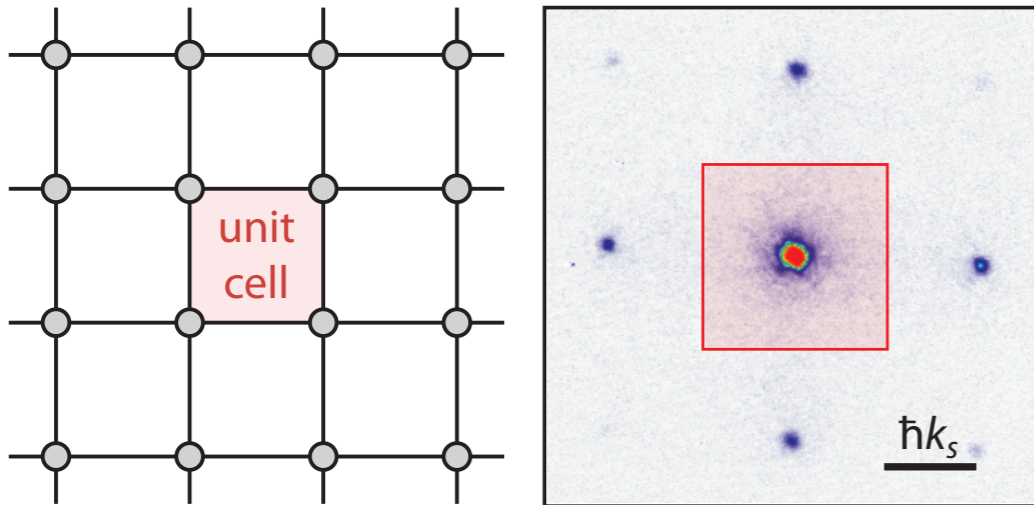


Due to the frustration introduced by the phase factors in $K(\mathbf{R})$, the condensation occurs for non-zero momenta.



Band structure

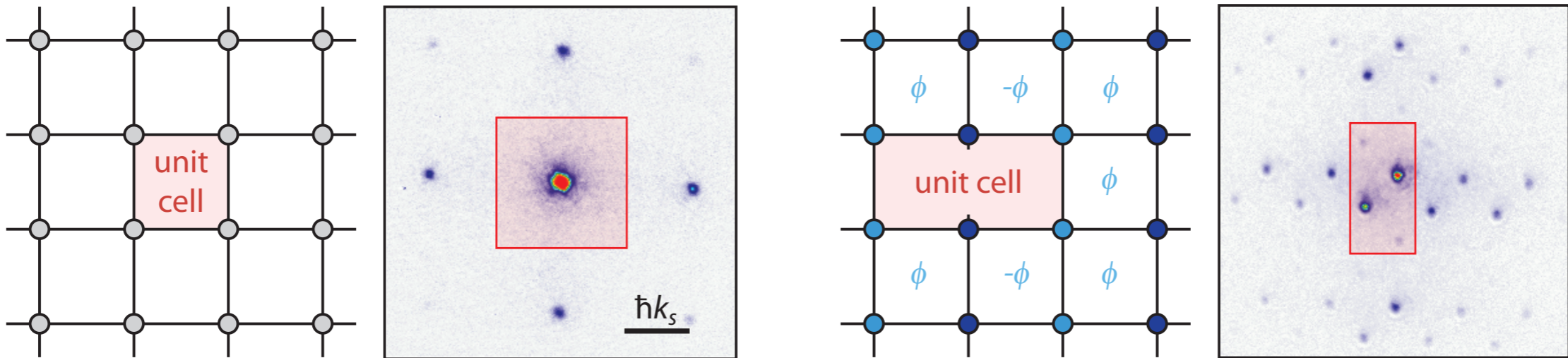
- 'Magnetic' Brillouin zone





Band structure

- 'Magnetic' Brillouin zone



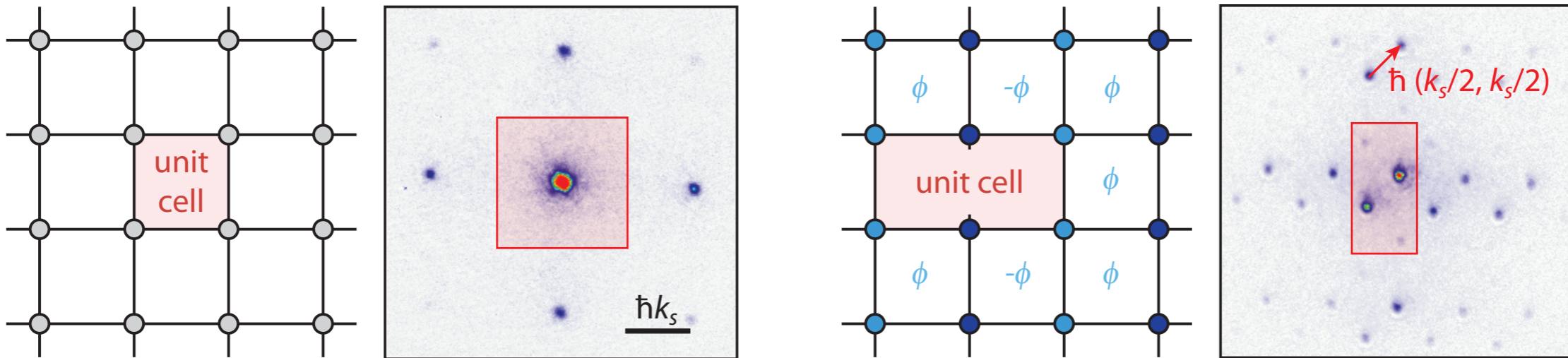
- Single-particle spectrum in the tight-binding approximation
From the magnetic translation symmetries:

$$\psi_{|k_x, k_y\rangle}(\mathbf{R} = m \mathbf{d}_x + n \mathbf{d}_y) = e^{i(m \cdot k_x d_x + n \cdot k_y d_y)} \times \begin{cases} \psi_e & m \text{ even} \\ \psi_o e^{i \frac{\pi}{2}(m+n)} & m \text{ odd} \end{cases} ,$$



Band structure

- 'Magnetic' Brillouin zone



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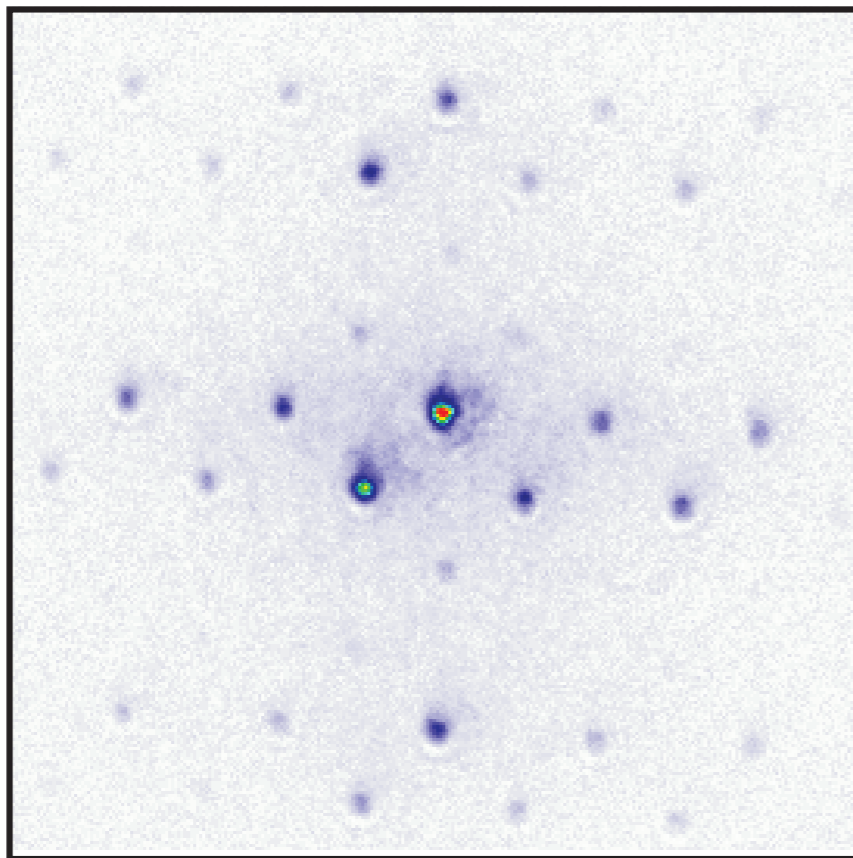
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An eigenstate $|k_x, k_y\rangle$ has two momentum components at (k_x, k_y) and $(k_x, k_y) + (k_s/2, k_s/2)$

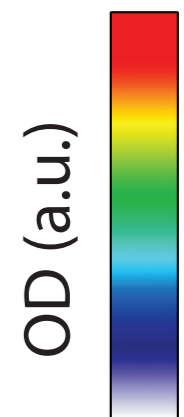
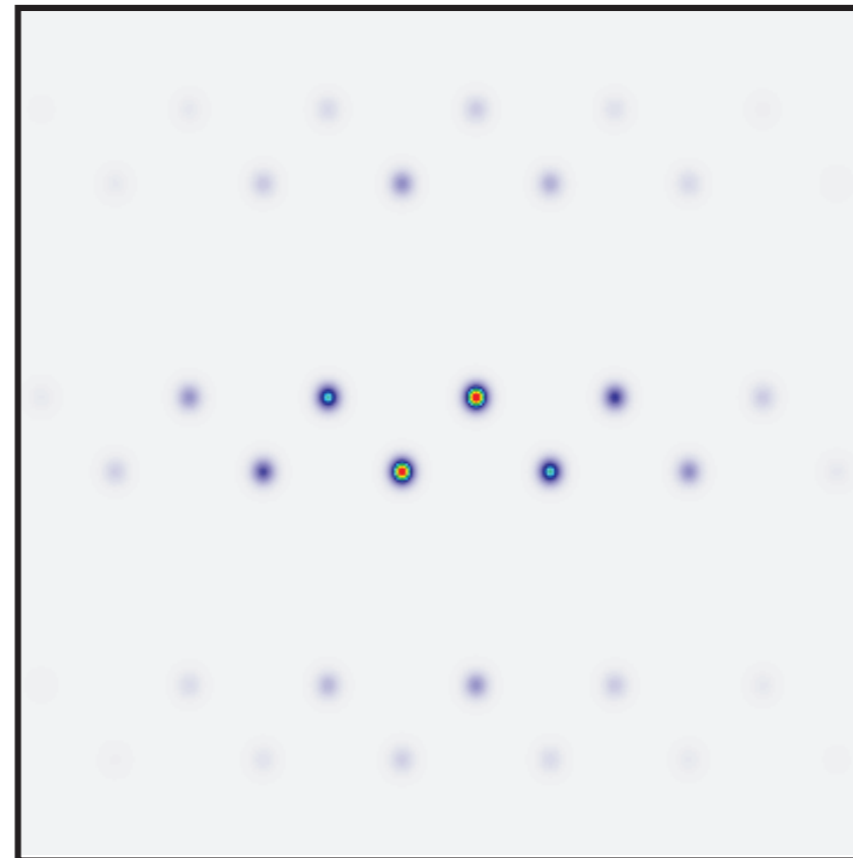


Momentum distribution ($J/K=1$): comparison with theory

experiment



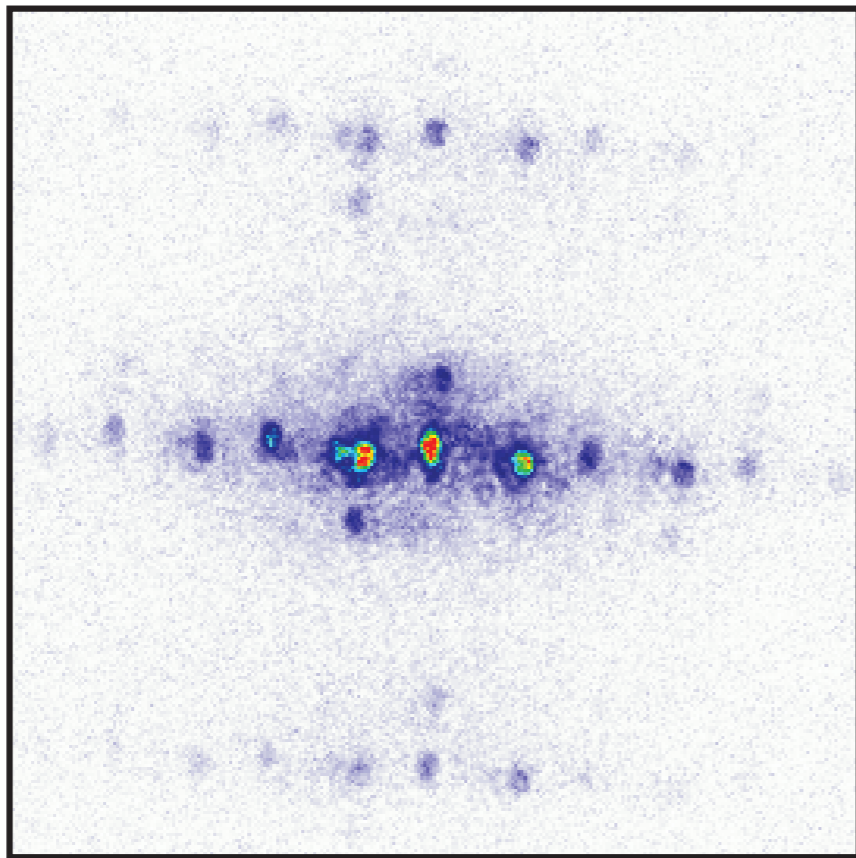
theory





Momentum distribution ($J/K=2.5$)

experiment

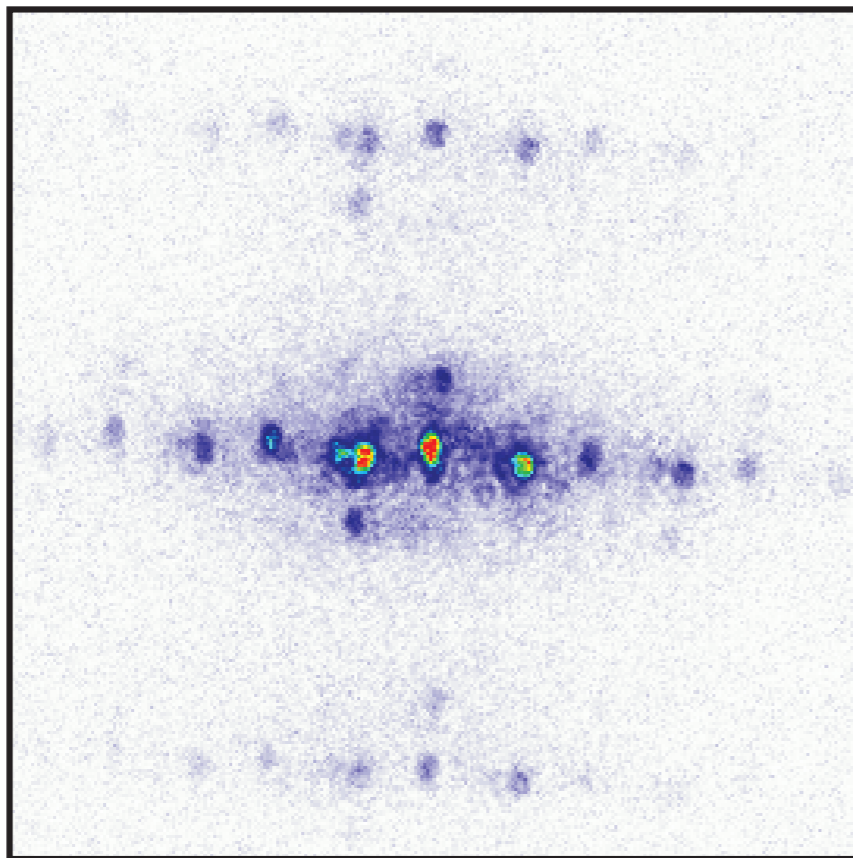


The diffraction peaks are splitted \longrightarrow two-fold ground state degeneracy

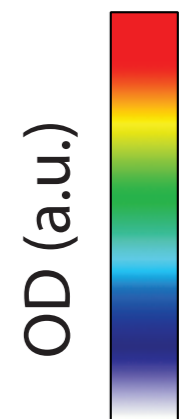
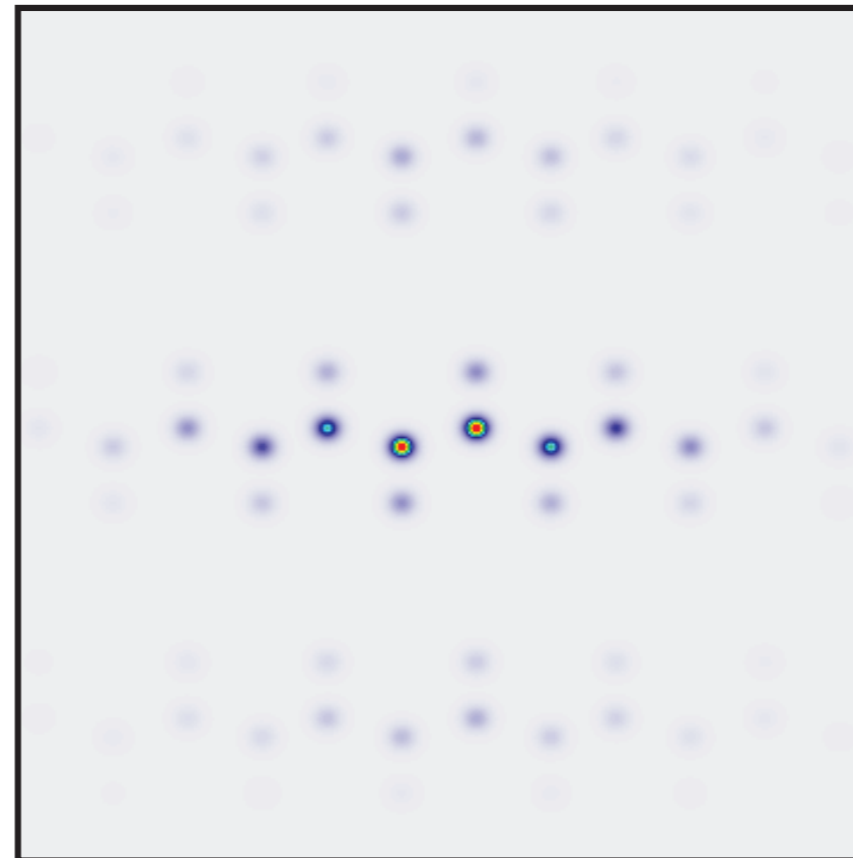


Momentum distribution ($J/K=2.5$)

experiment



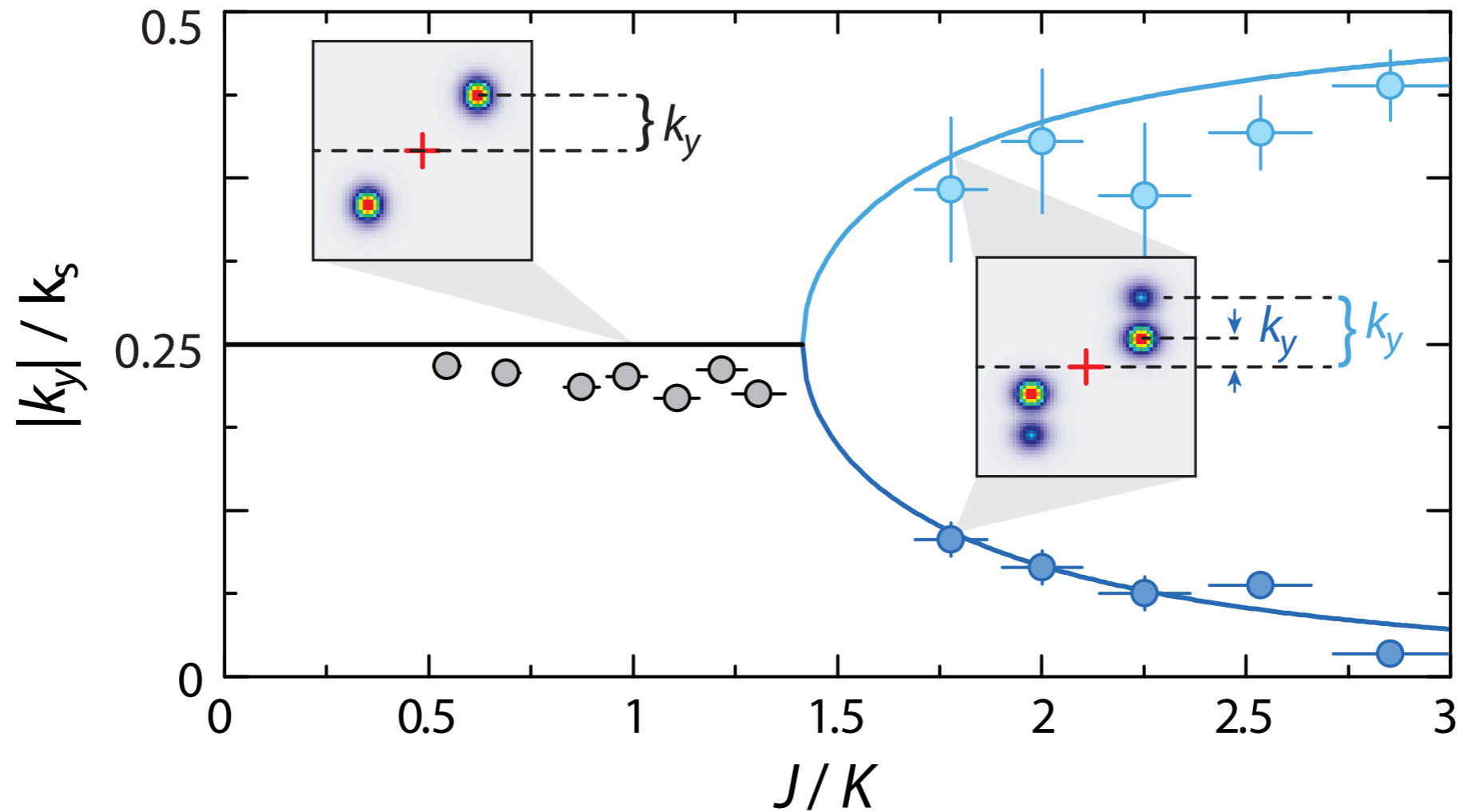
theory



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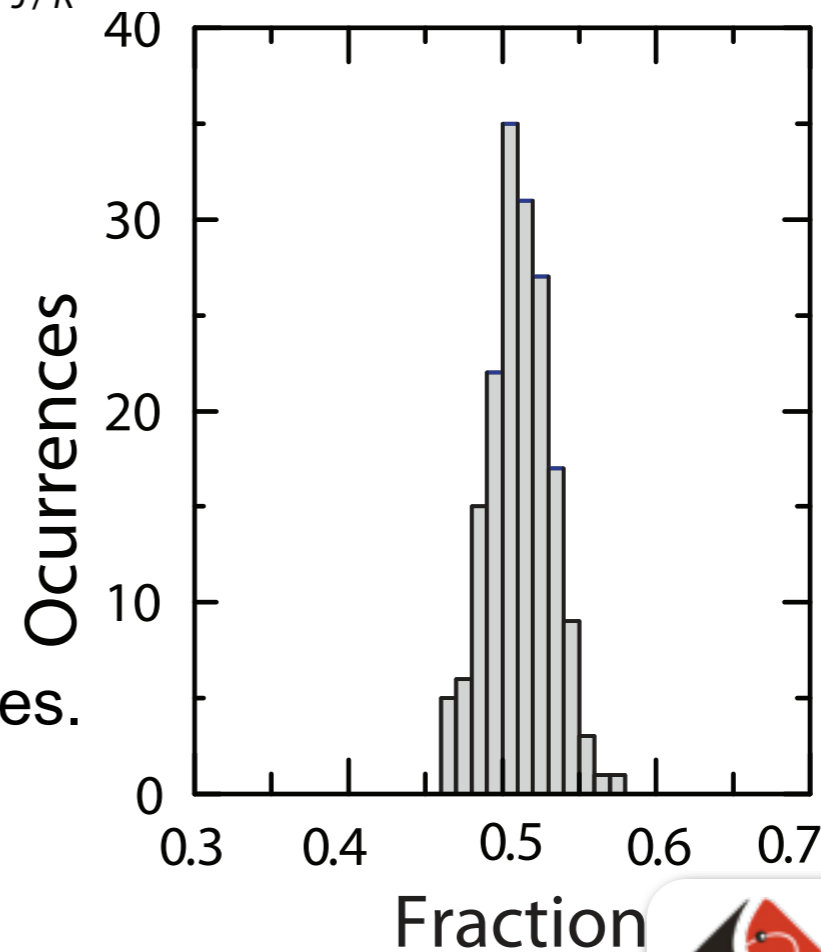
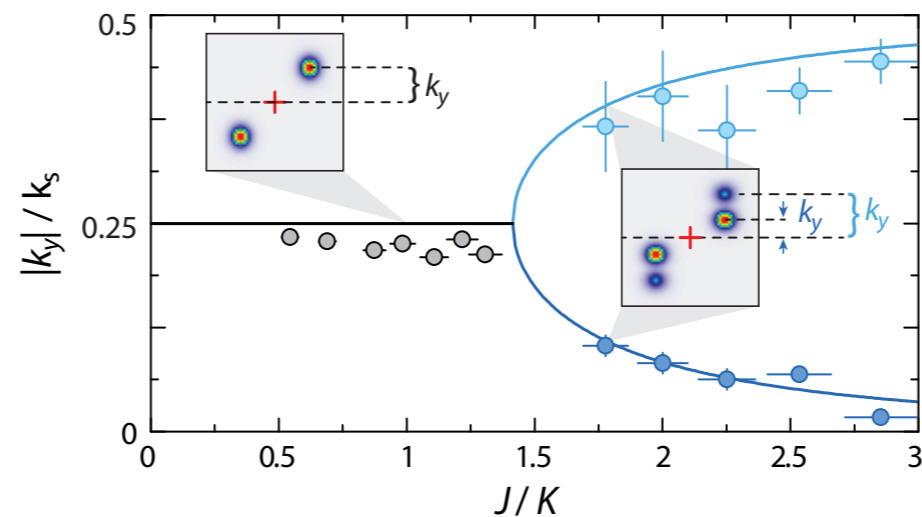
Momenta of the two degenerate ground states



G. Moeller, N. Cooper, PRA **82**, 1 (2010)
J. Struck *et al.*, Science **333**, 996 (2011)



Momenta of the two degenerate ground states



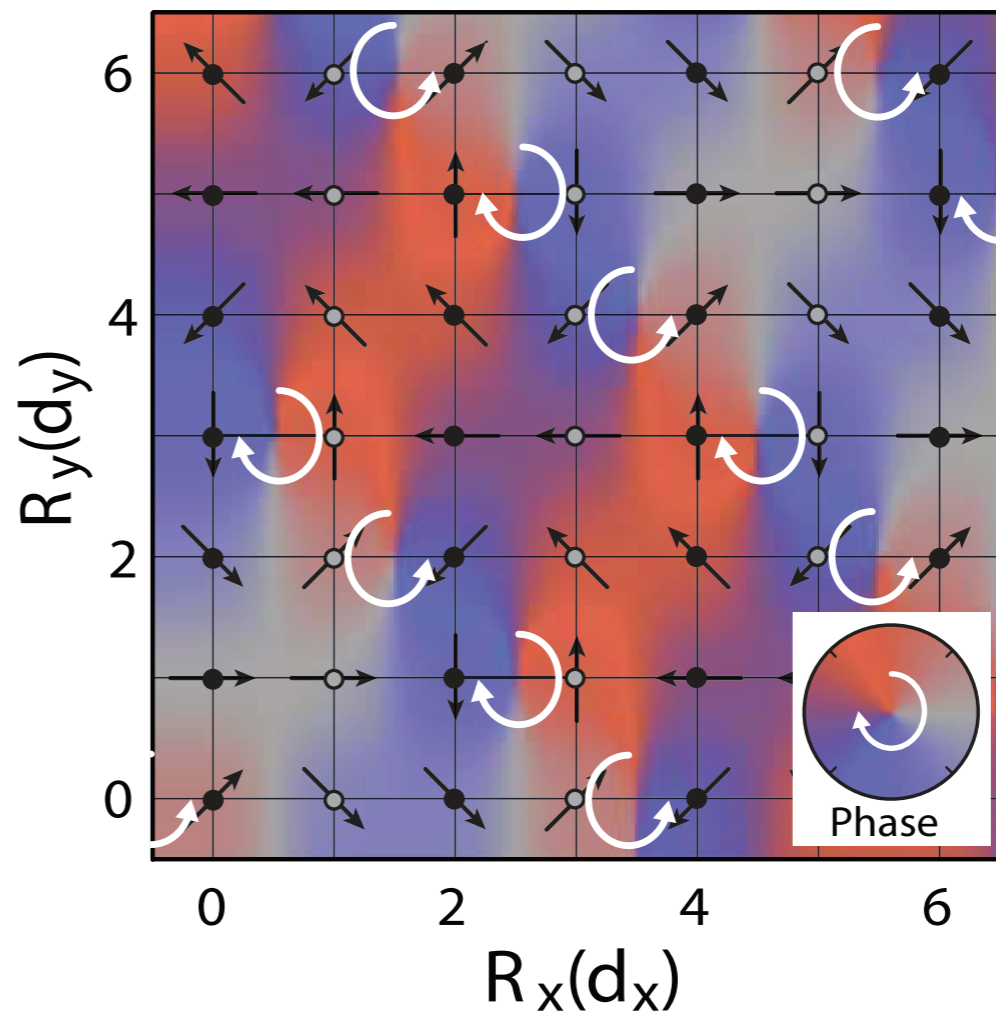
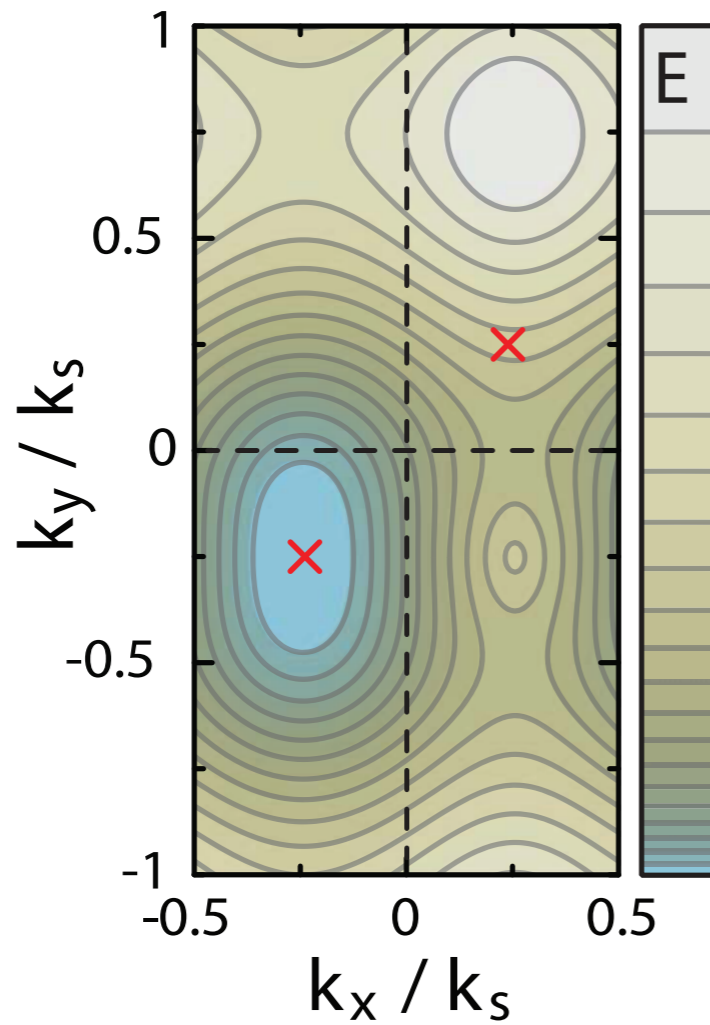
In the case of ground state degeneracy, we observe an equal population of both states.

G. Moeller, N. Cooper, PRA **82**, 1 (2010)
J. Struck *et al.*, Science **333**, 996 (2011)



Dispersion relation and spatial phase distribution

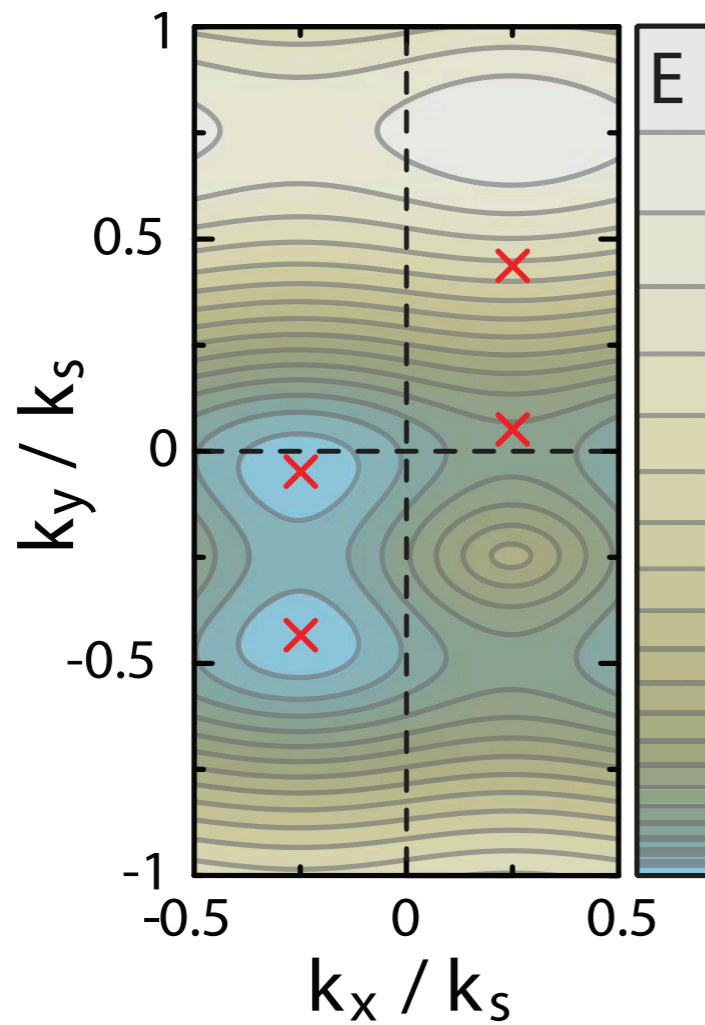
$$J = K$$



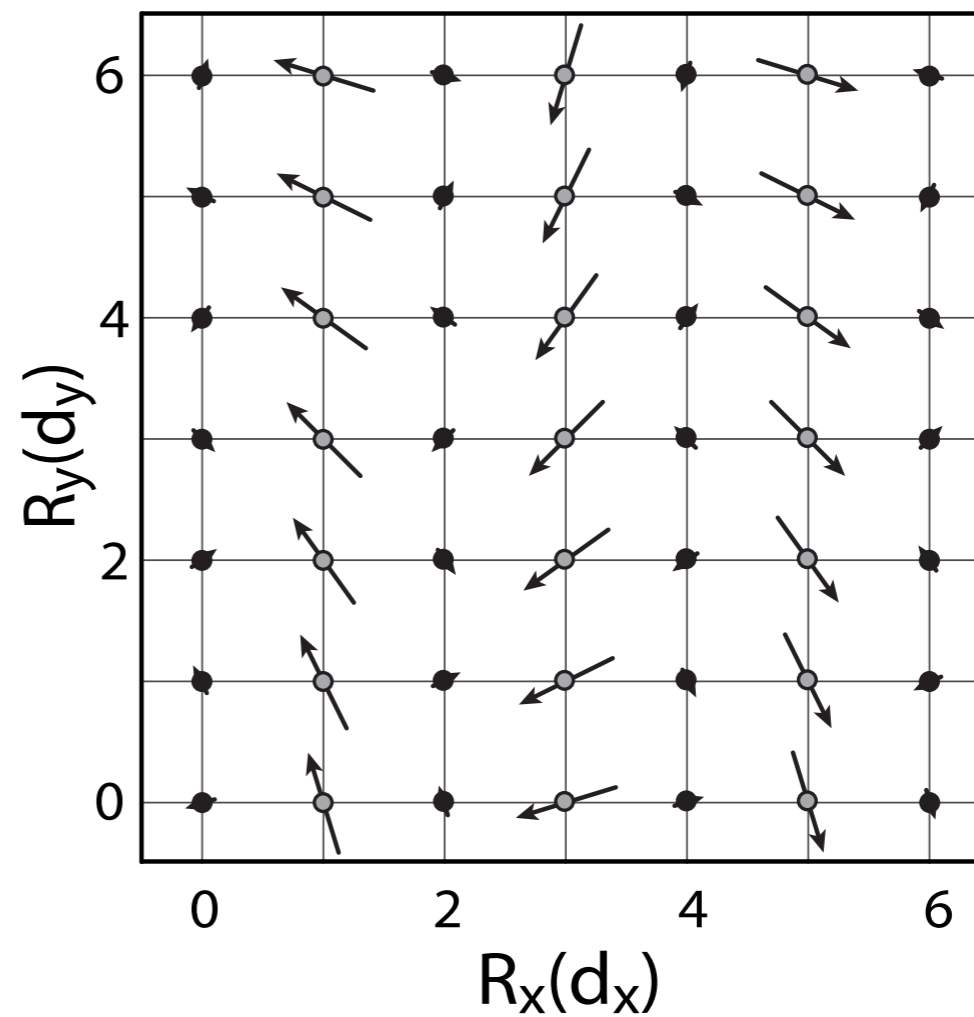


Dispersion relation and spatial phase distribution

$J=2.5 K$

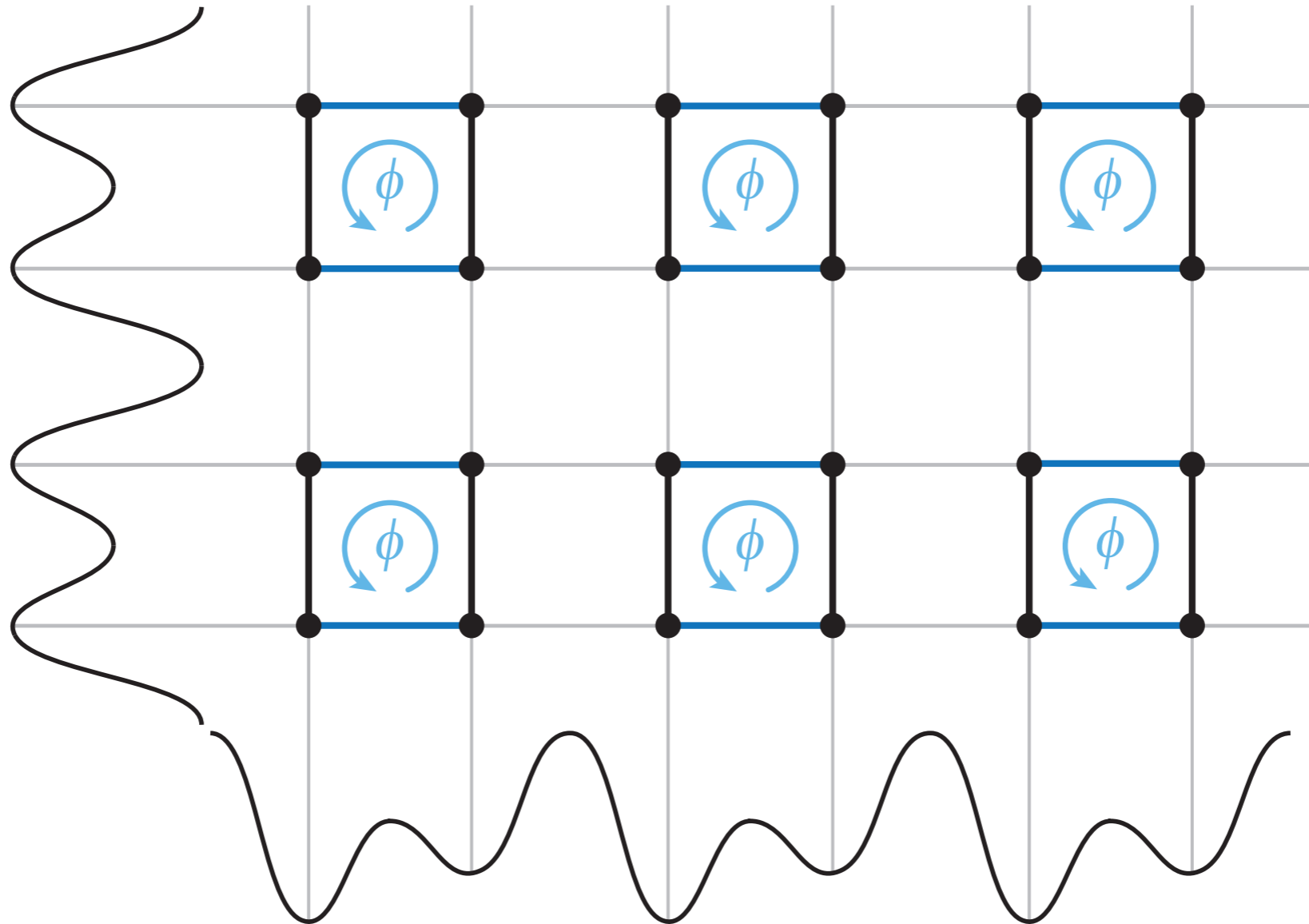


Striped density pattern





A lattice of plaquettes

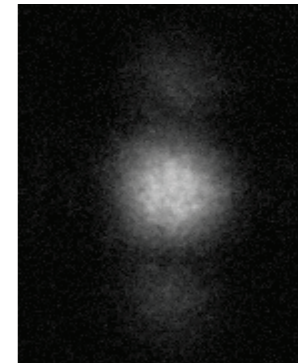
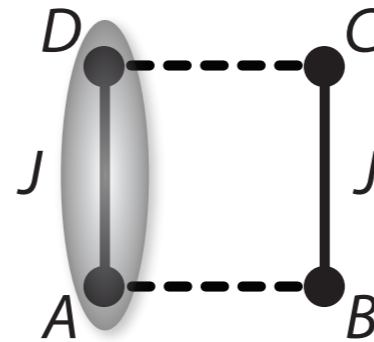




State preparation for phase imprinting

- Load single atoms into ground state of tilted plaquettes:

$$|\psi_0\rangle = \frac{|A\rangle + |D\rangle}{\sqrt{2}}$$

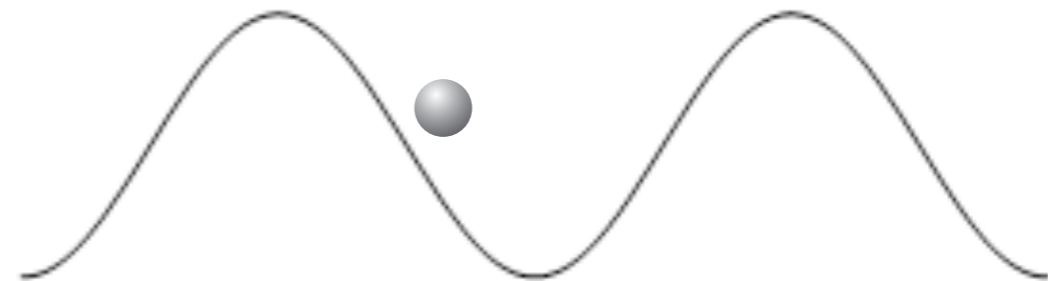
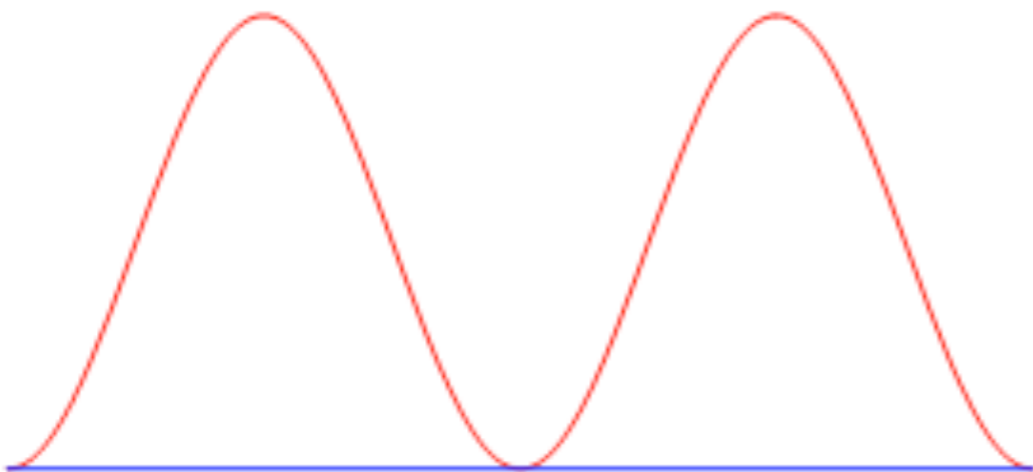
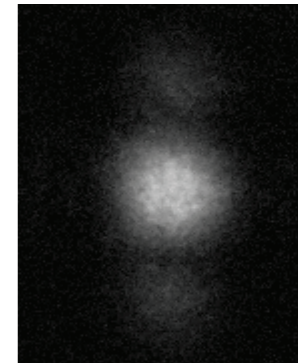
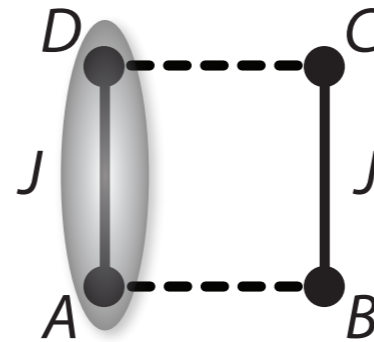




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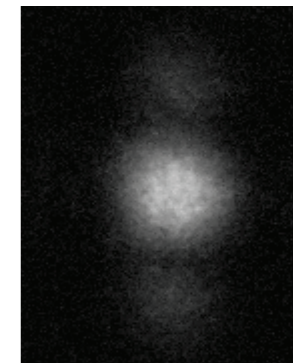
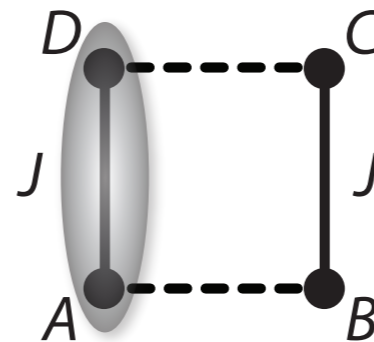




State preparation for phase imprinting

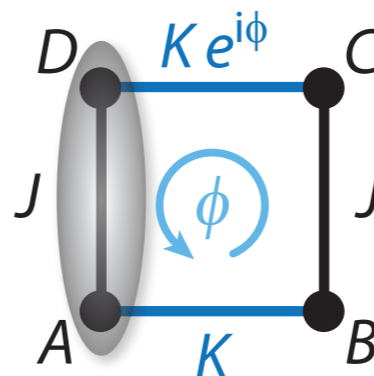
- Load single atoms into ground state of tilted plaquettes:

$$|\psi_0\rangle = \frac{|A\rangle + |D\rangle}{\sqrt{2}}$$



- Switch on Raman coupling to induce atom transfer to the B, C sites
In the limit $J \ll K$ the state is coupled to

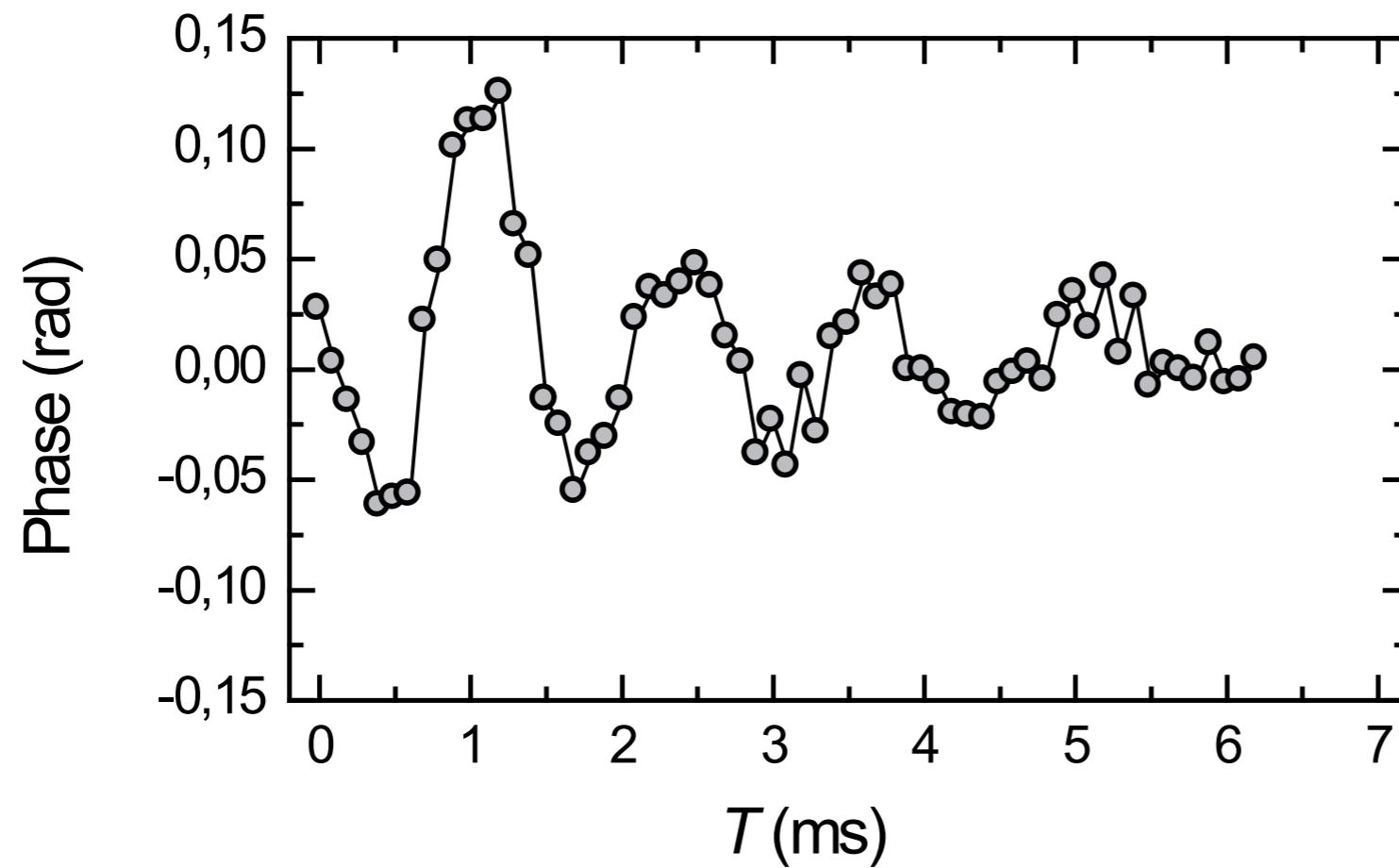
$$|\psi_1\rangle = \frac{|B\rangle + i|C\rangle}{\sqrt{2}}$$



Non-trivial
phase evolution

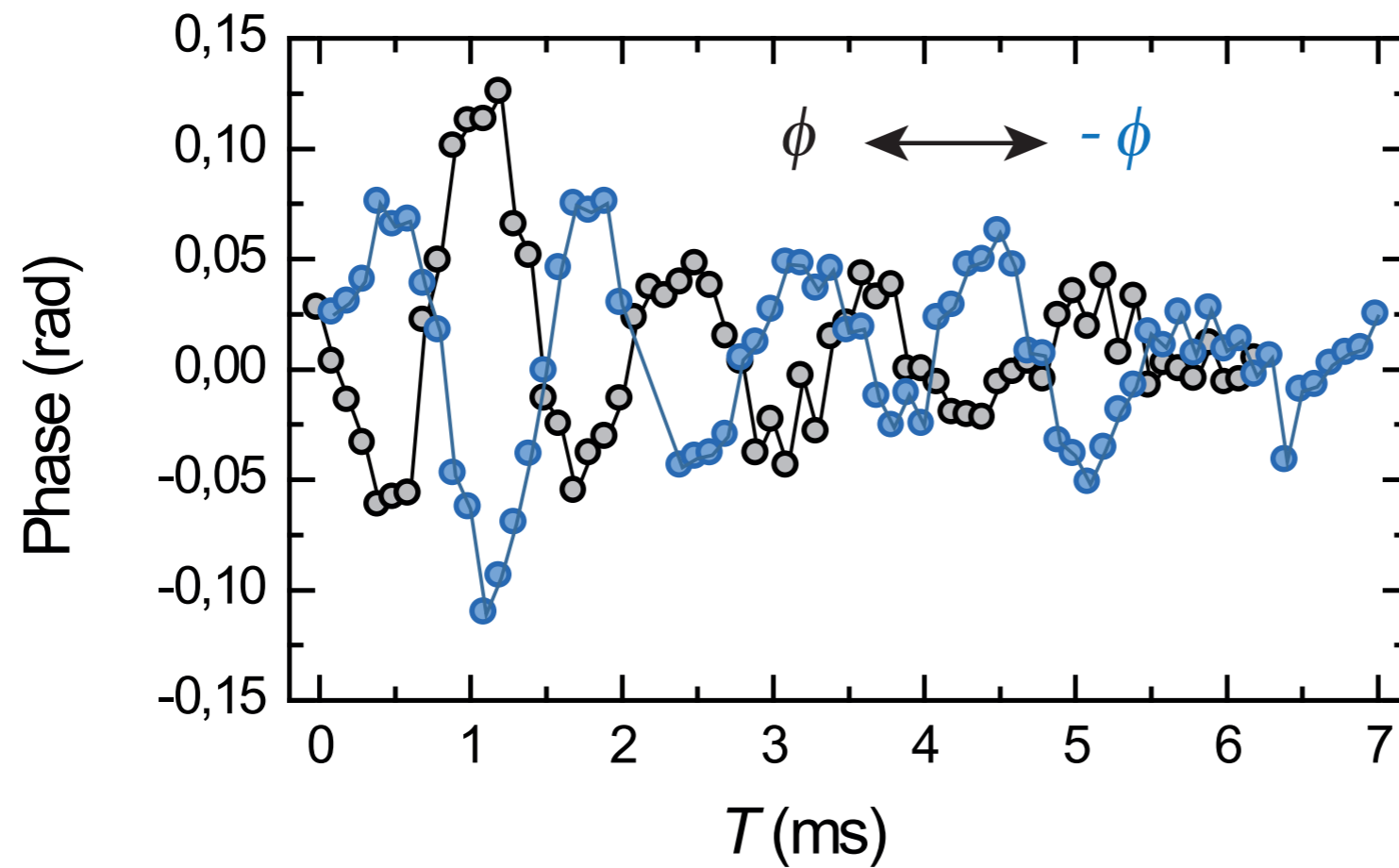


Observation of phase-imprinting



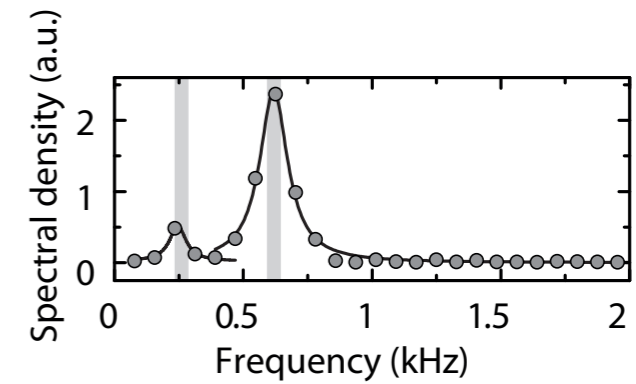
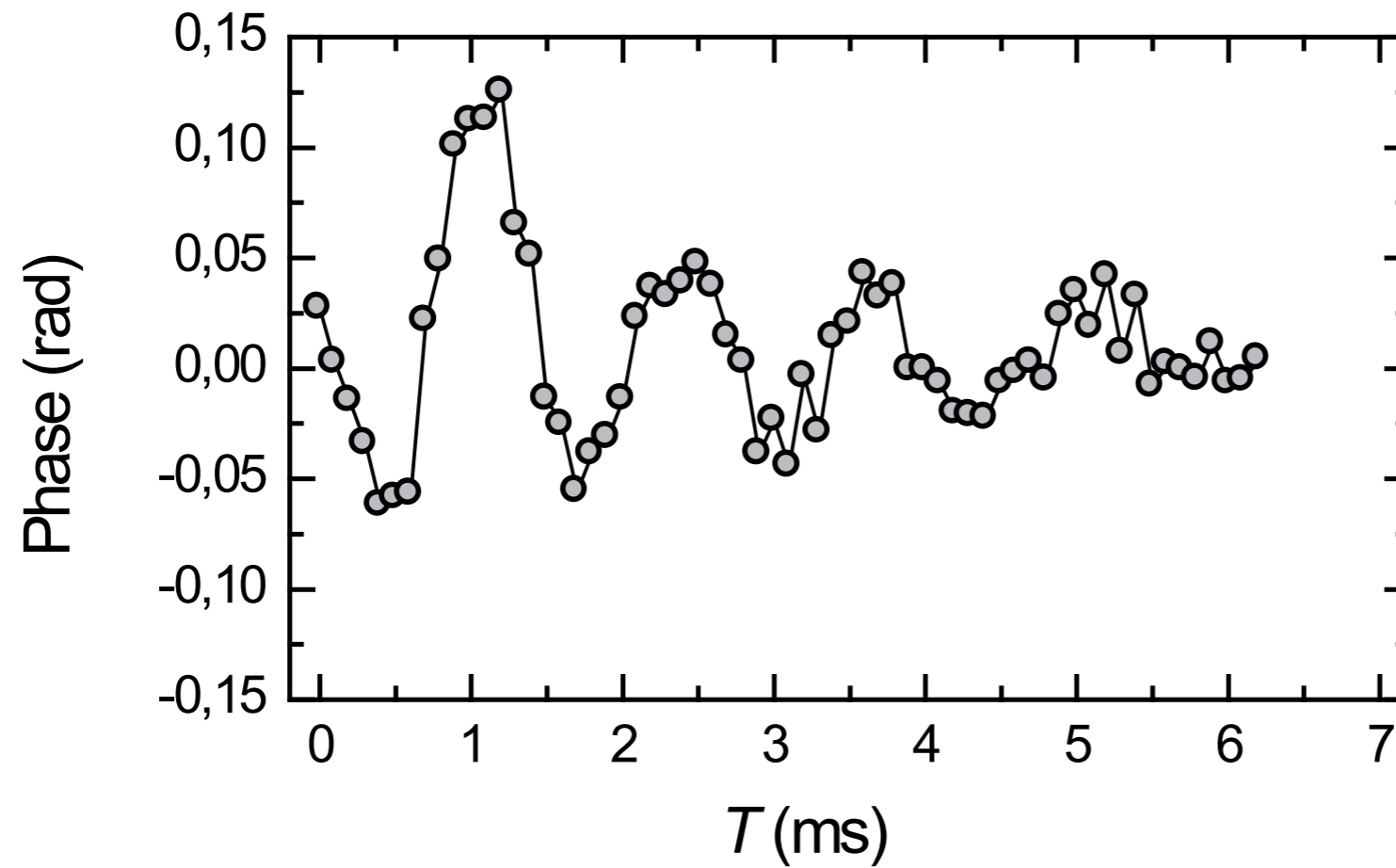


Observation of phase-imprinting





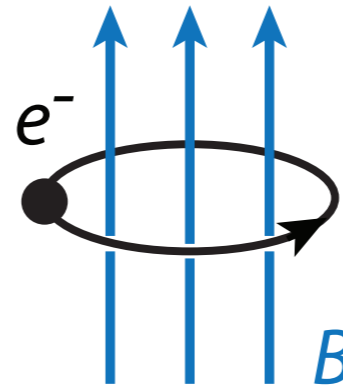
Observation of phase-imprinting





Quantum 'Cyclotron' Orbit

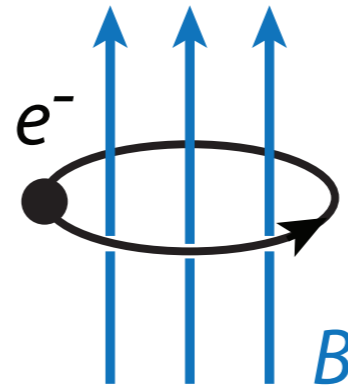
- Classical:
Charged particle in
a uniform magnetic field



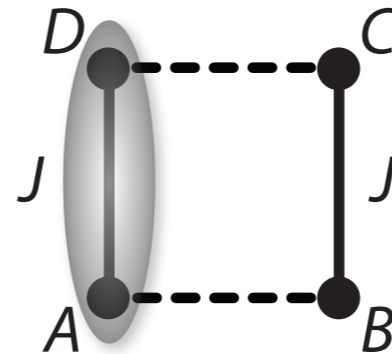


Quantum 'Cyclotron' Orbit

- Classical:
Charged particle in
a uniform magnetic field



- Measure quantum analogue:
Initial state:
Single atom in ground
state of tilted plaquette

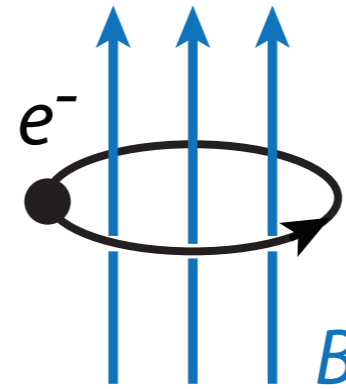


$$|\psi_0\rangle = \frac{|A\rangle + |D\rangle}{\sqrt{2}}$$

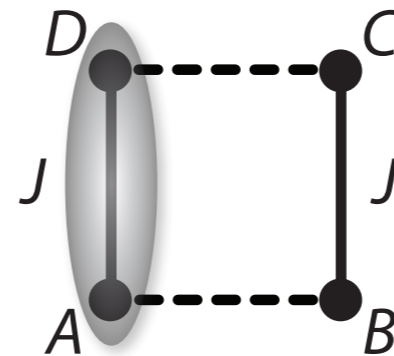


Quantum 'Cyclotron' Orbit

- Classical:
Charged particle in a uniform magnetic field

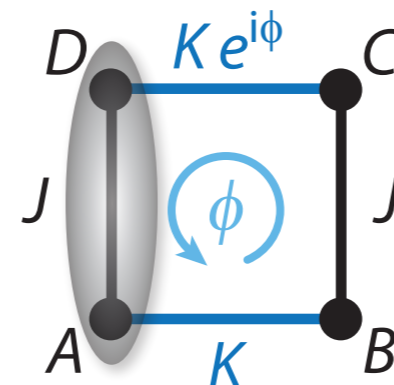


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Initial state:
Single atom in ground state of tilted plaquette



$$|\psi_0\rangle = \frac{|A\rangle + |D\rangle}{\sqrt{2}}$$

Switch on Raman coupling to induce atom transfer

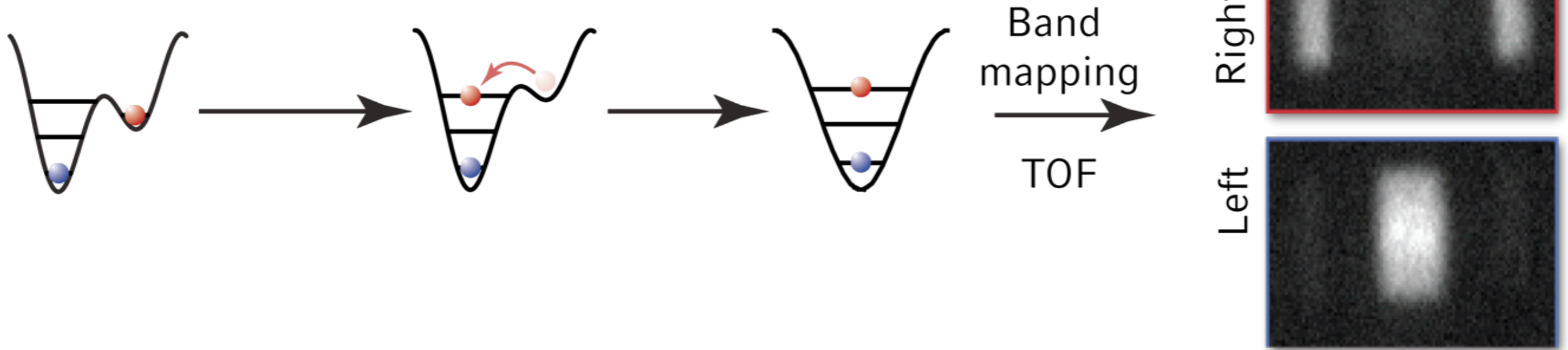


Site-resolved detection



Site-resolved detection

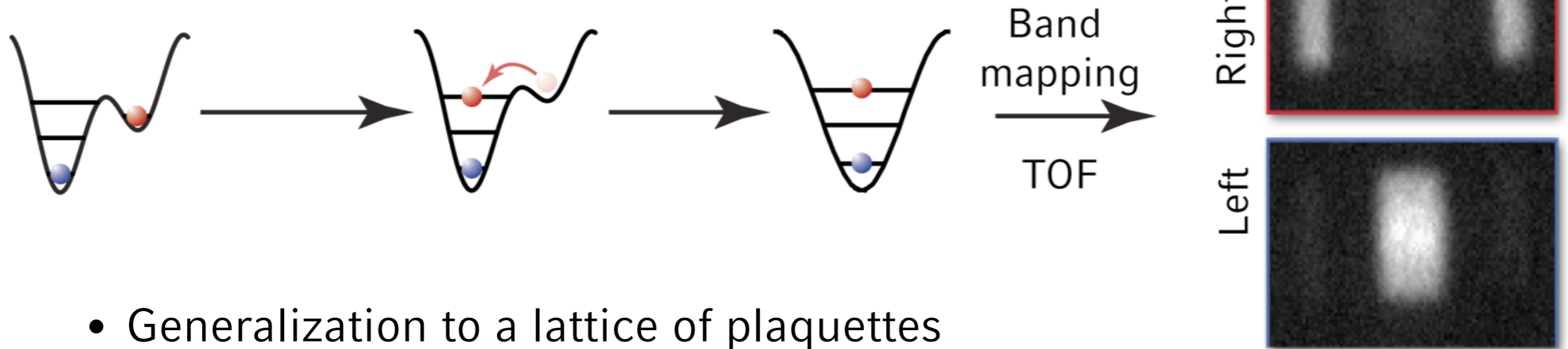
- Site-resolved detection in double-wells



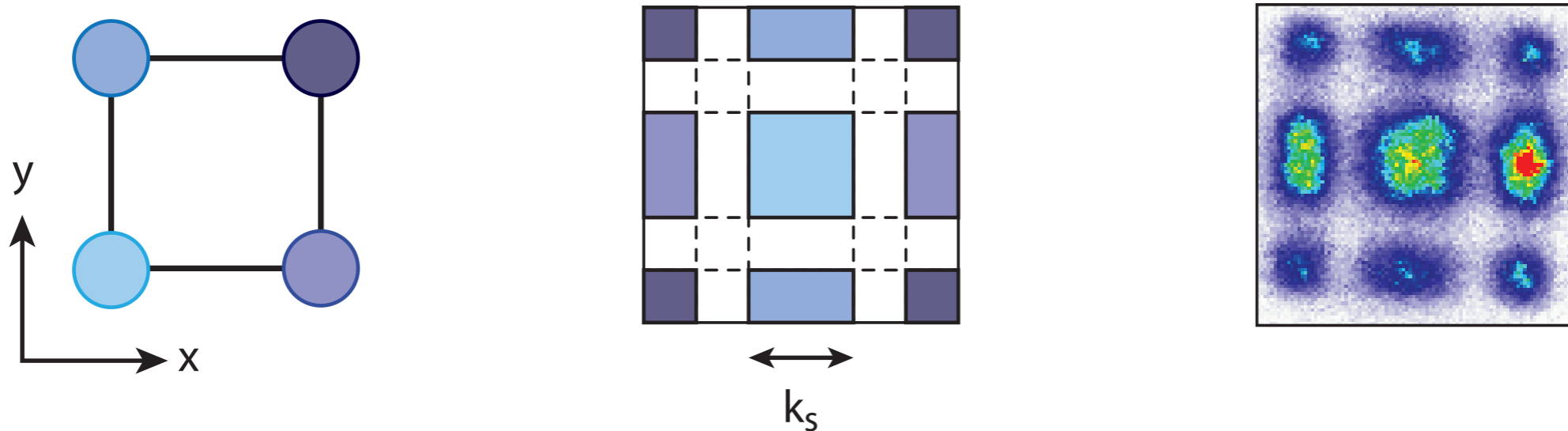


Site-resolved detection

- Site-resolved detection in double-wells



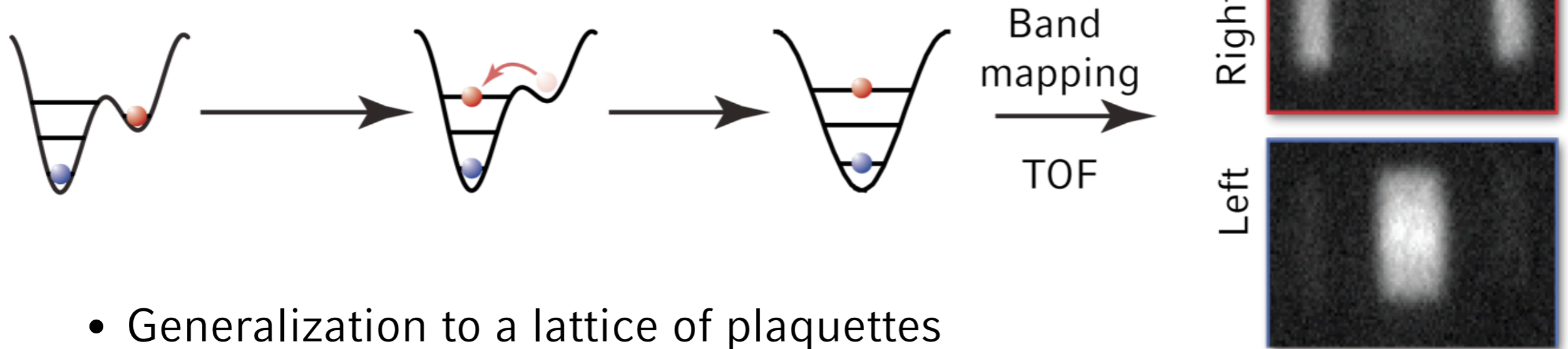
- Generalization to a lattice of plaquettes



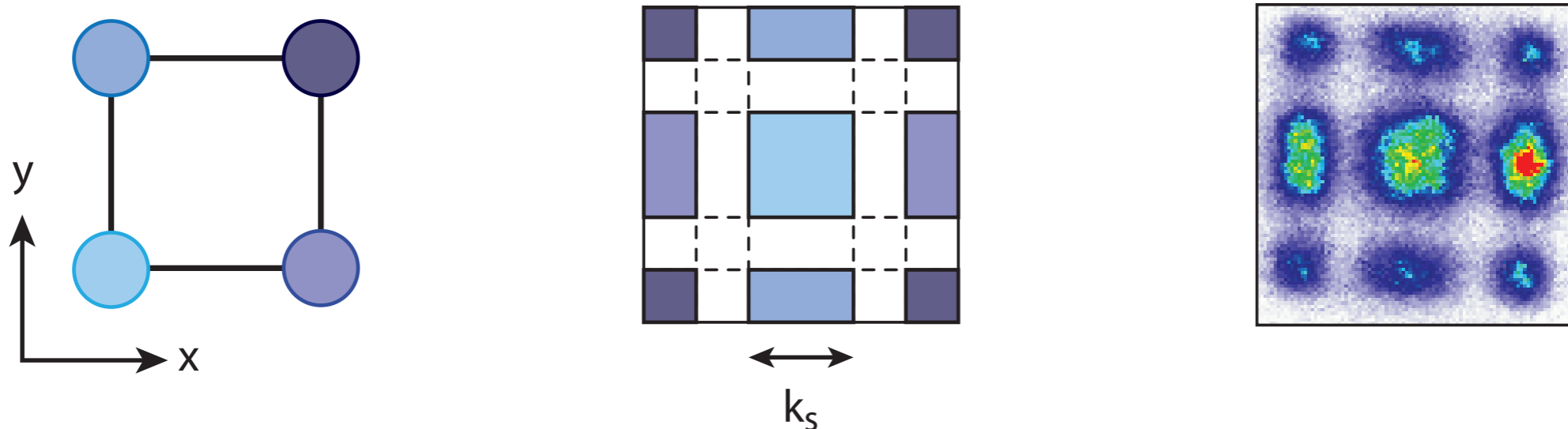


Site-resolved detection

- Site-resolved detection in double-wells



- Generalization to a lattice of plaquettes

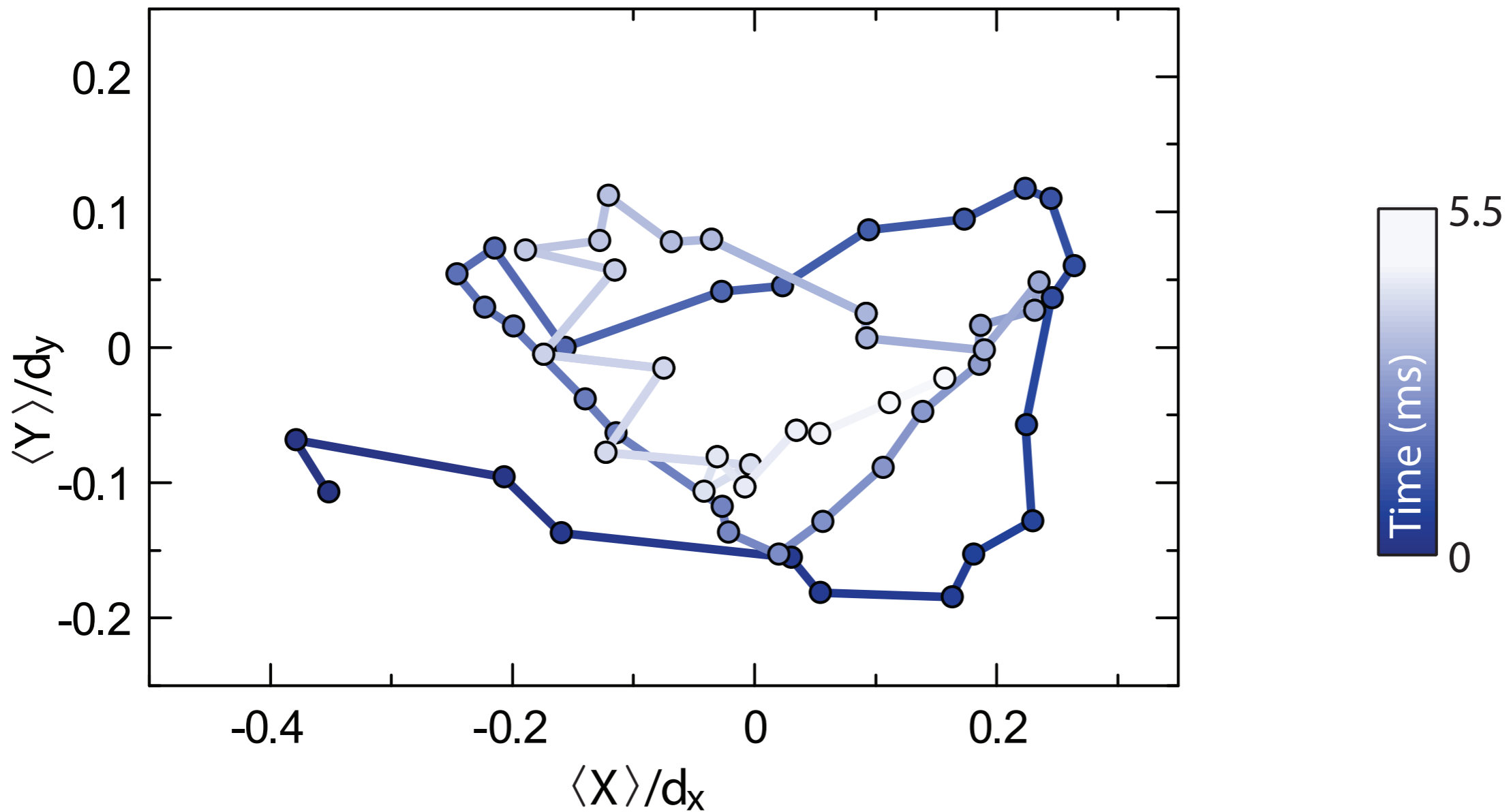


- Evaluate mean atom positions $\langle X \rangle$ and $\langle Y \rangle$



'Cyclotron' Orbit

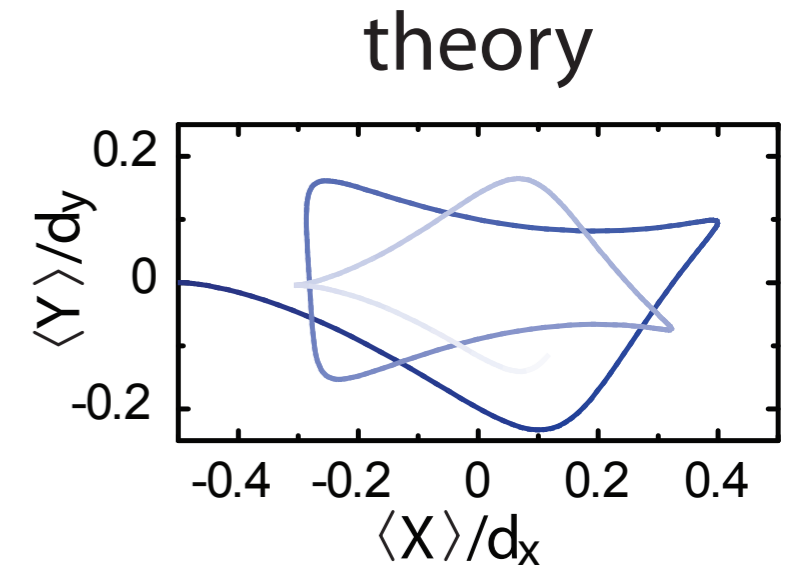
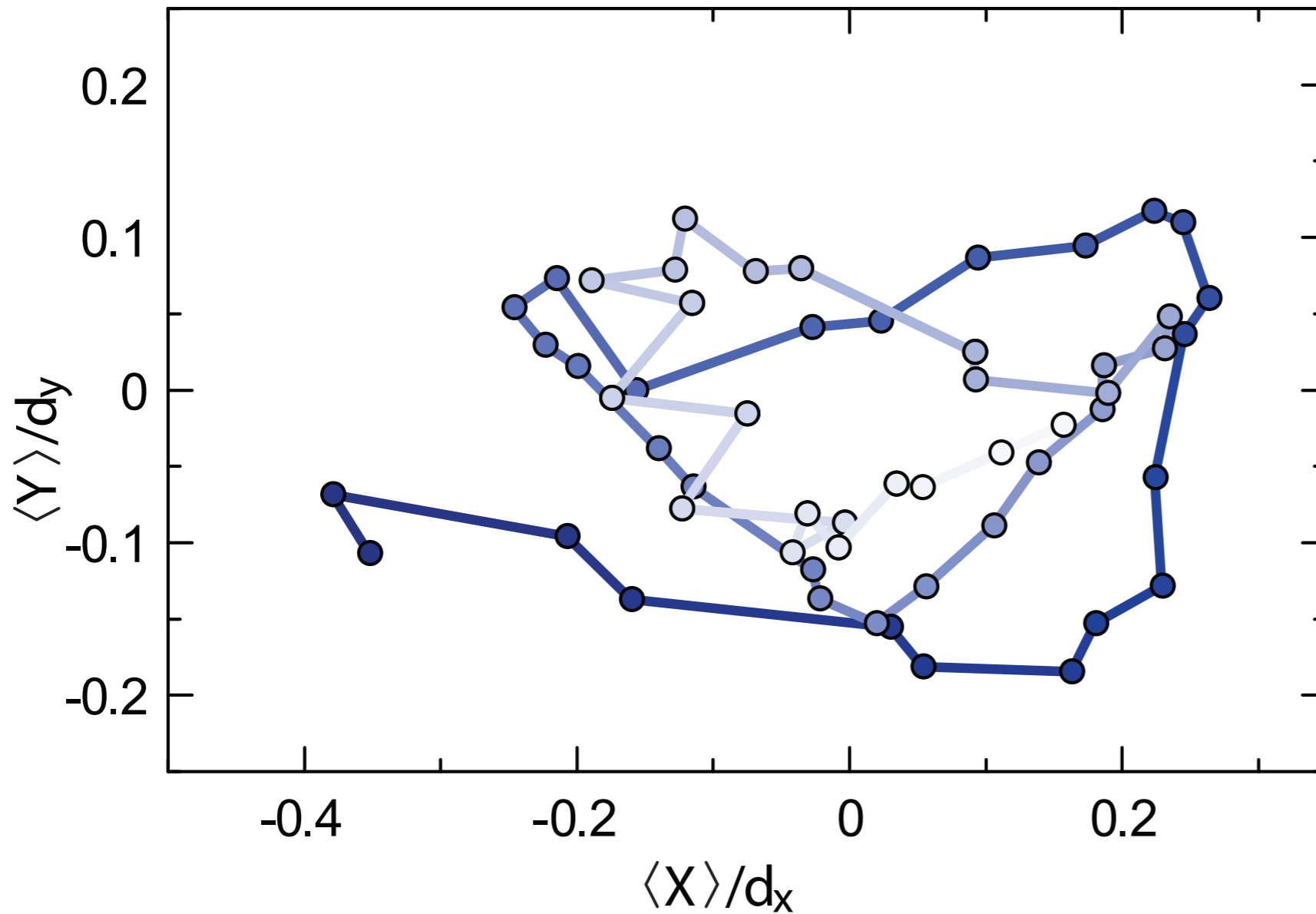
The mean atom position during the evolution.





'Cyclotron' Orbit

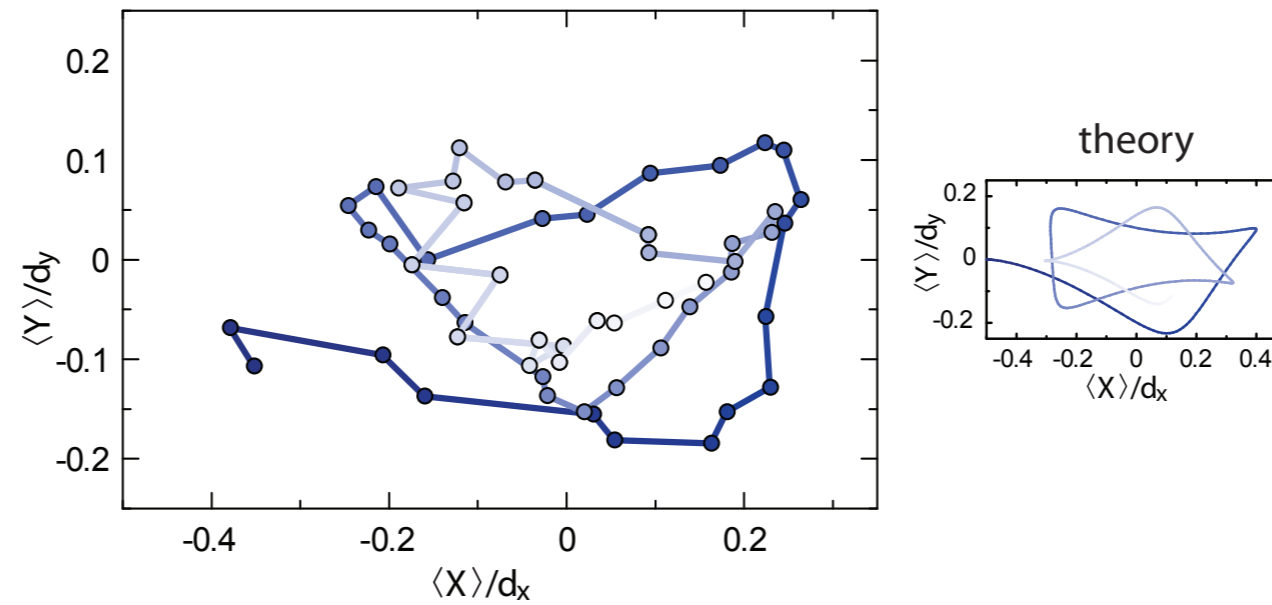
The mean atom position during the evolution.





'Cyclotron' Orbit

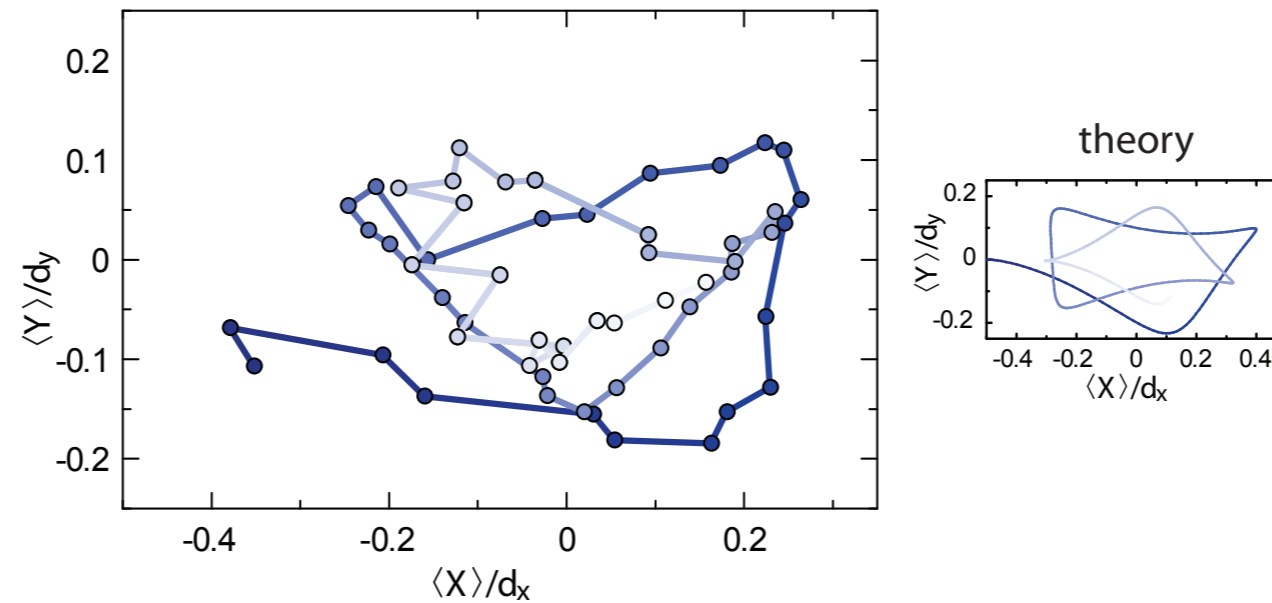
The mean atom position during the evolution.





'Cyclotron' Orbit

The mean atom position during the evolution.



From this evolution we fit the value of the phase

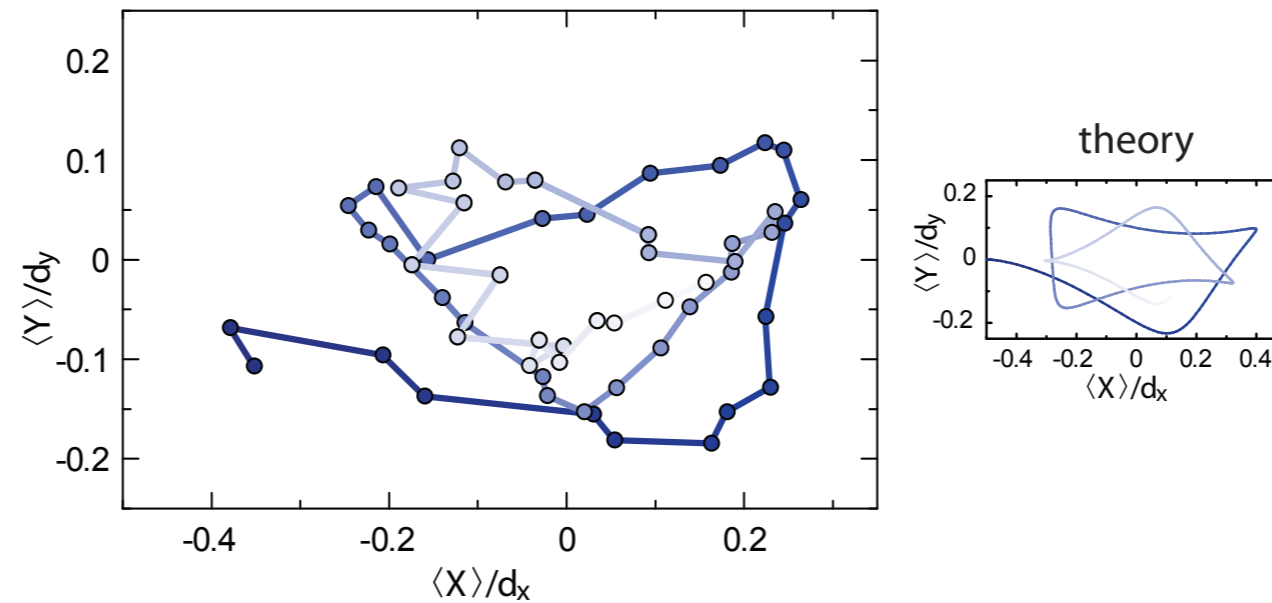
$$\phi = 0.73(5) \pi/2$$

Deviation from $\phi = \pi/2$



'Cyclotron' Orbit

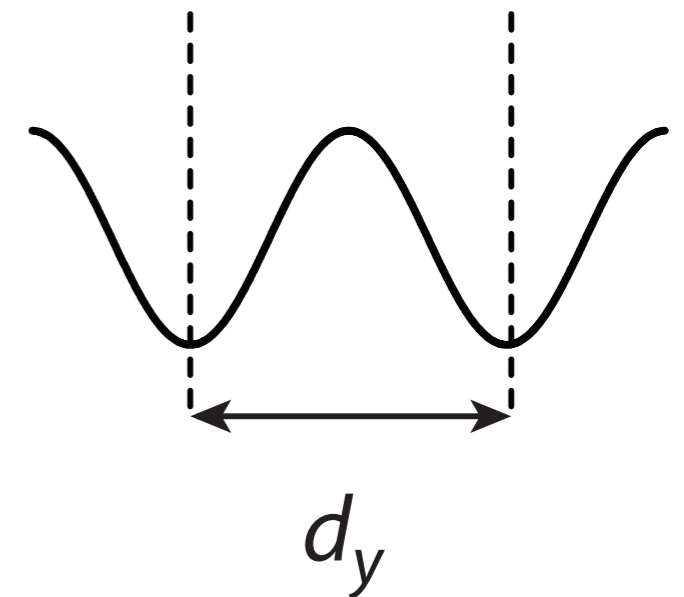
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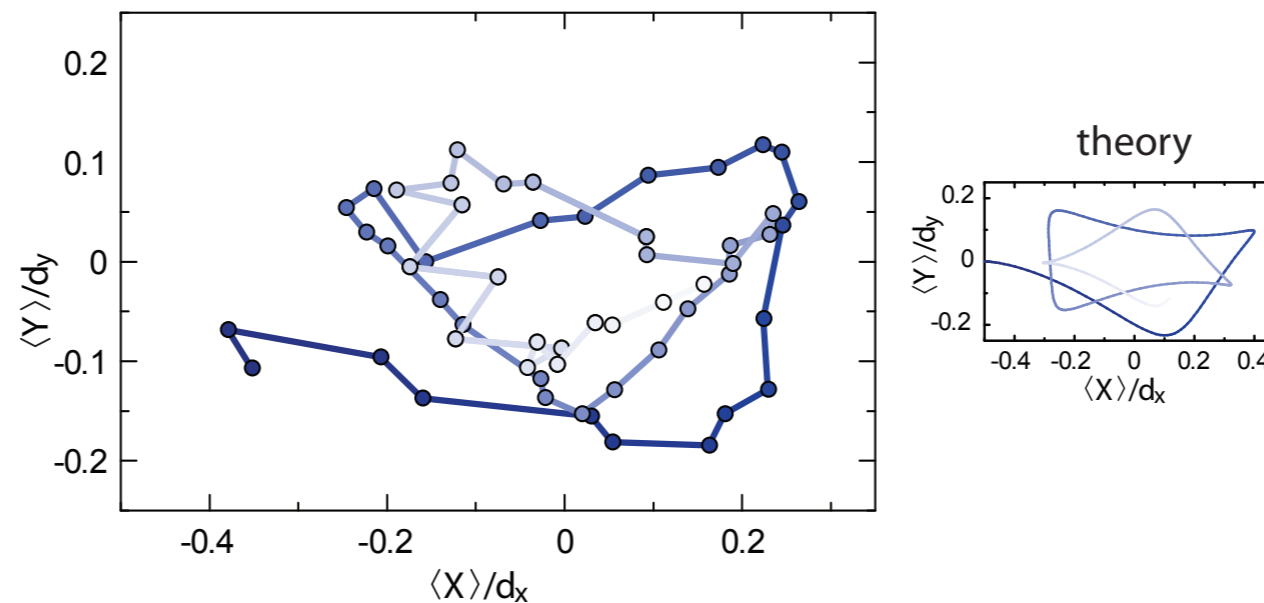
Deviation from $\phi = \pi/2$





'Cyclotron' Orbit

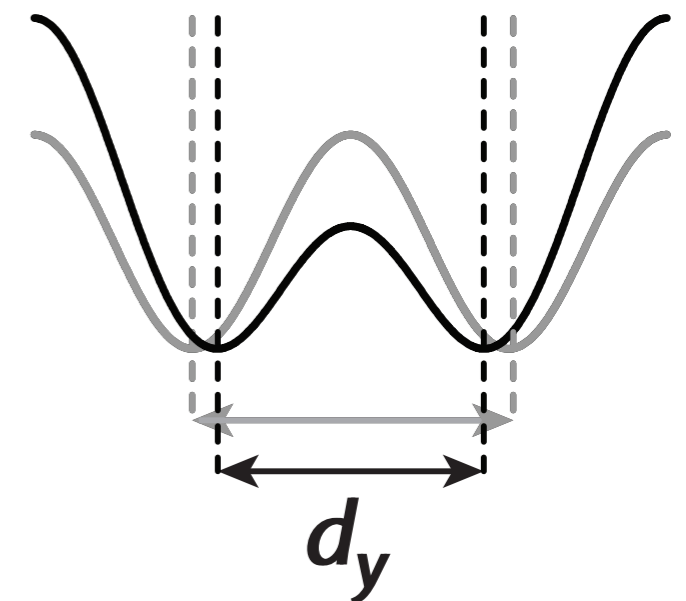
The mean atom position during the evolution.



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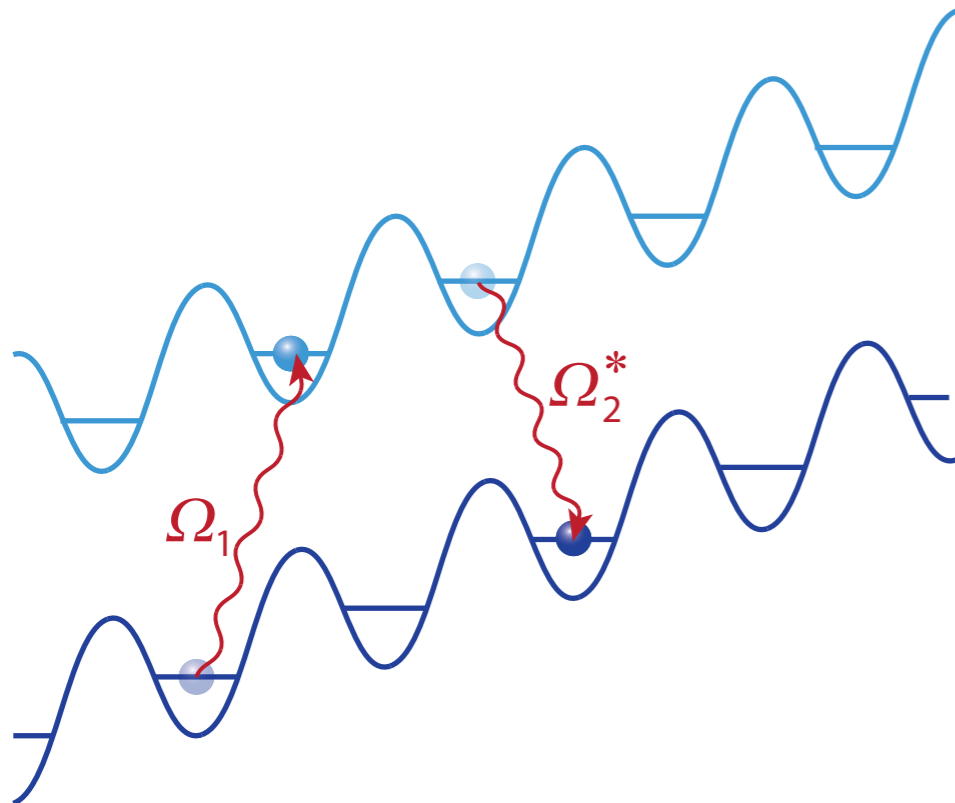
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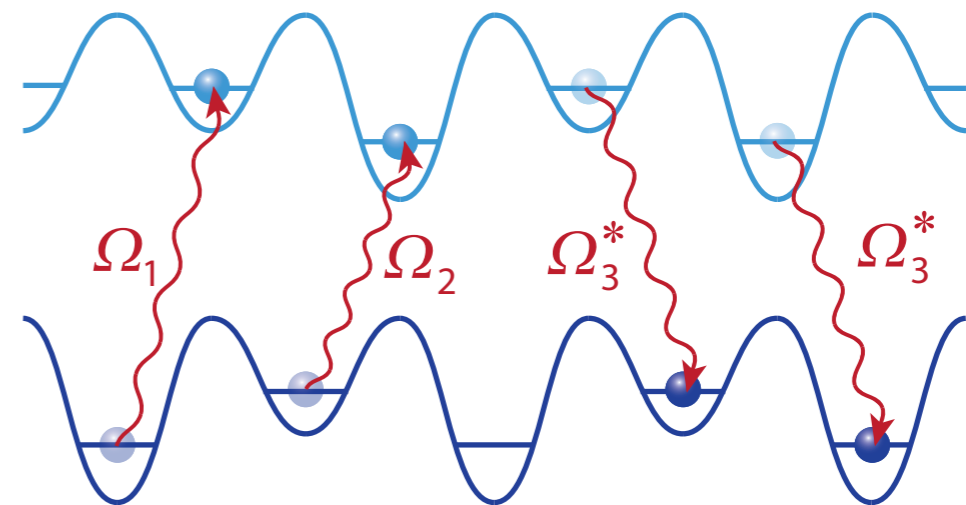


Outlook: Rectify the flux



Using a linear potential

D. Jaksch & P. Zoller, NJP 5, 56 (2003)

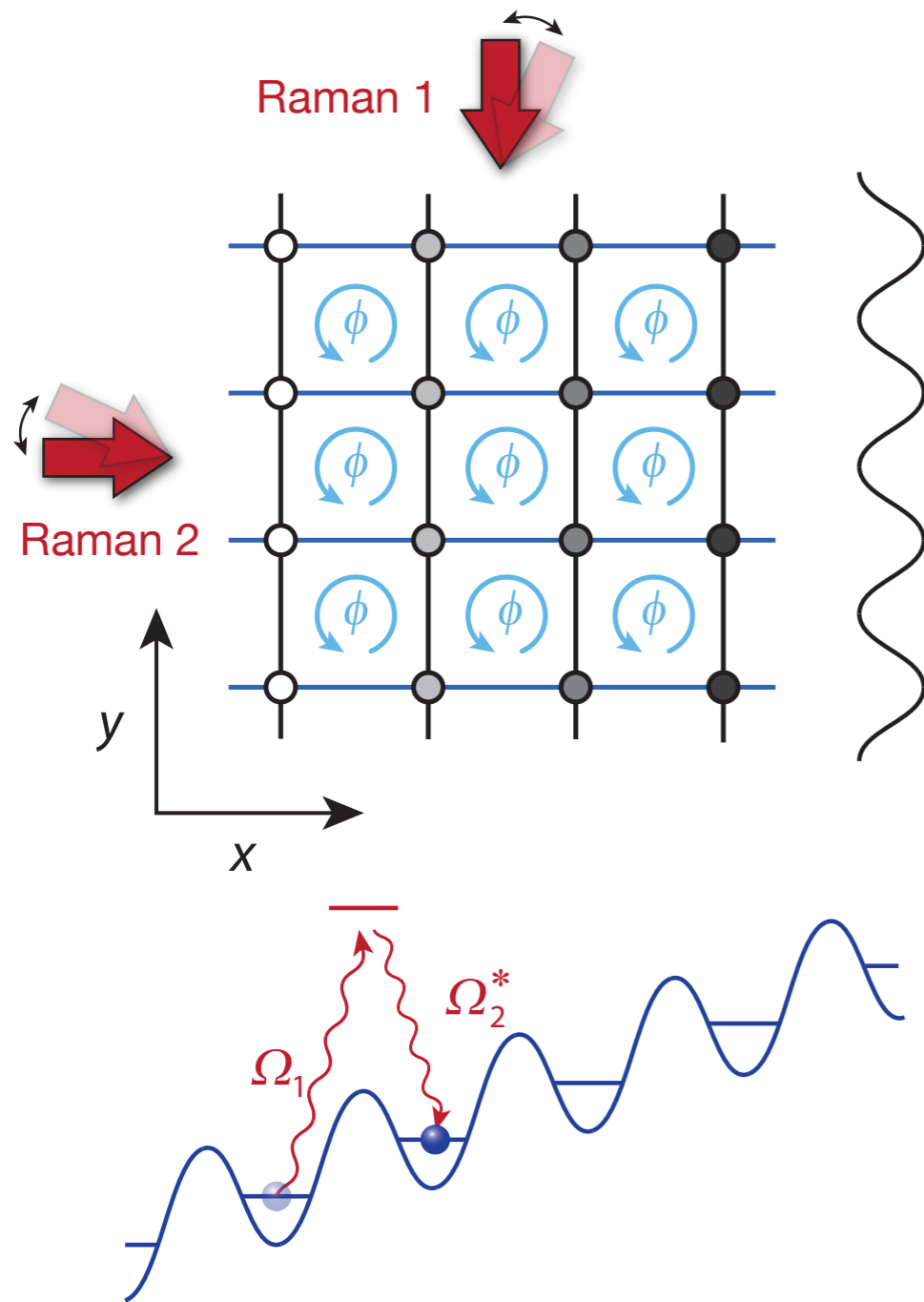


Using a superlattice

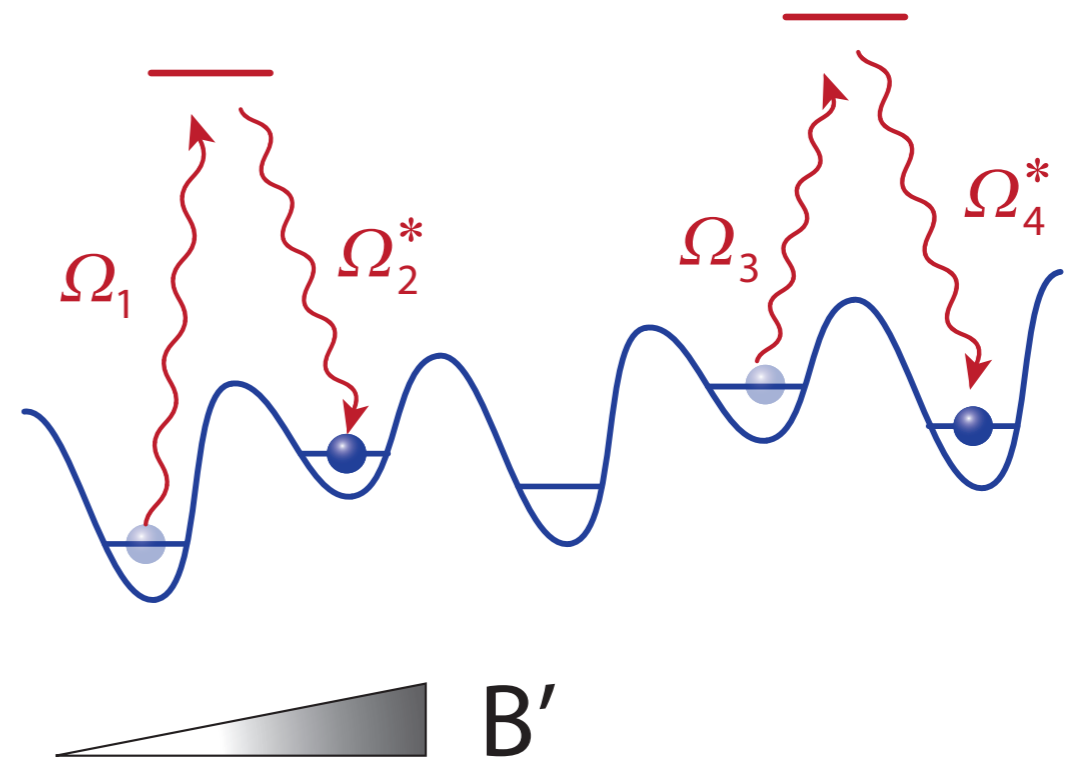
F. Gerbier & J. Dalibard, NJP 12, 033007 (2010)



Outlook: Rectify the flux



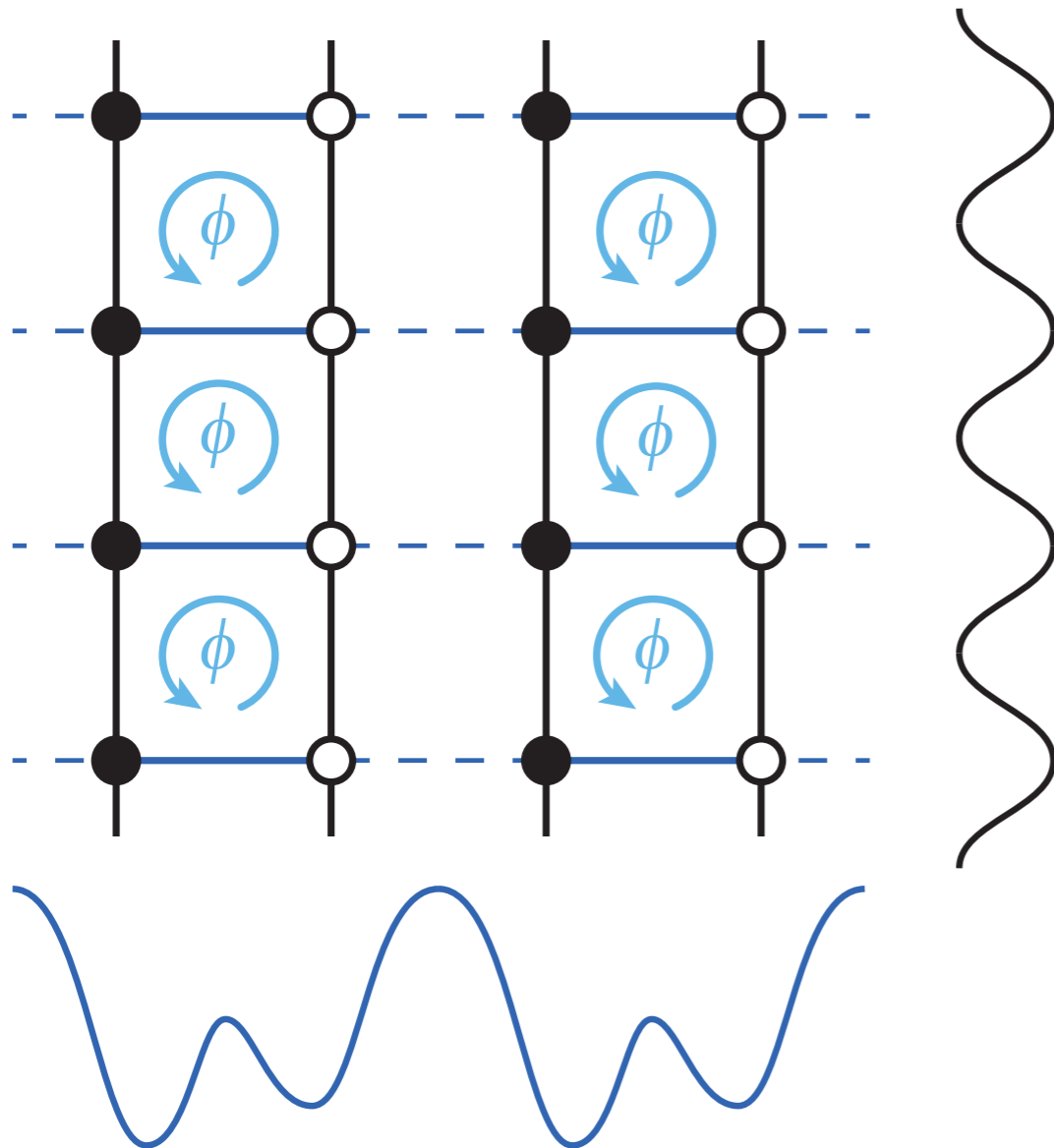
Using a linear potential
D. Jaksch & P. Zoller, NJP **5**, 56 (2003)



Using a superlattice + Gradient
F. Gerbier & J. Dalibard, NJP **12**, 033007 (2010)



Outlook: Rectify the flux

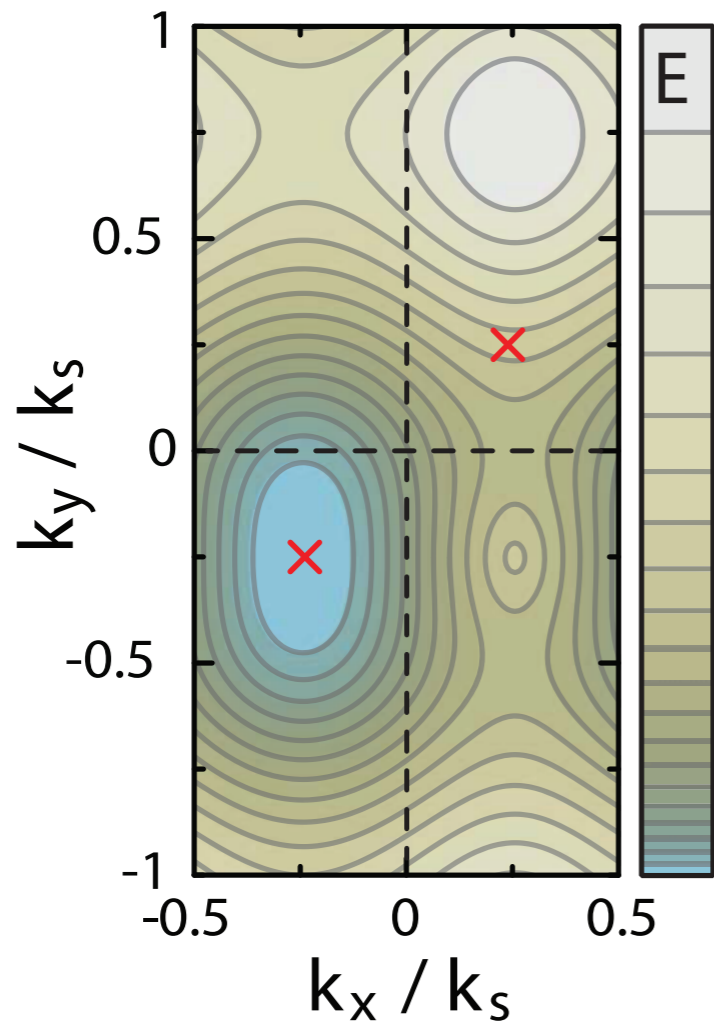


Ladders in a magnetic field:

- ? Observables
- ? Edge current
- ? Bifurcation point
- ? Strong interaction
- ? Dirac point



Outlook: Rectify the flux

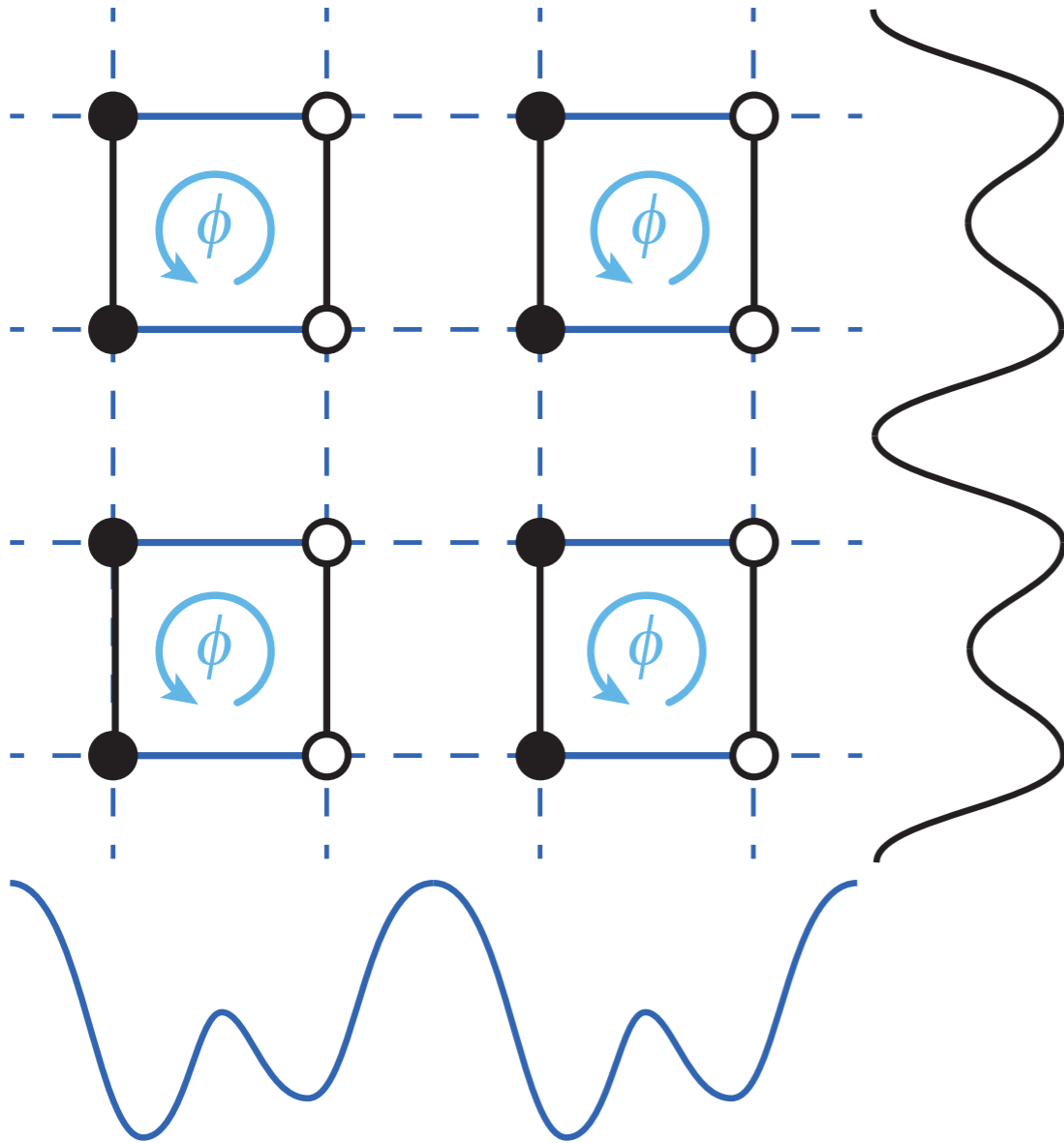


Ladders in a magnetic field:

- ? Observables
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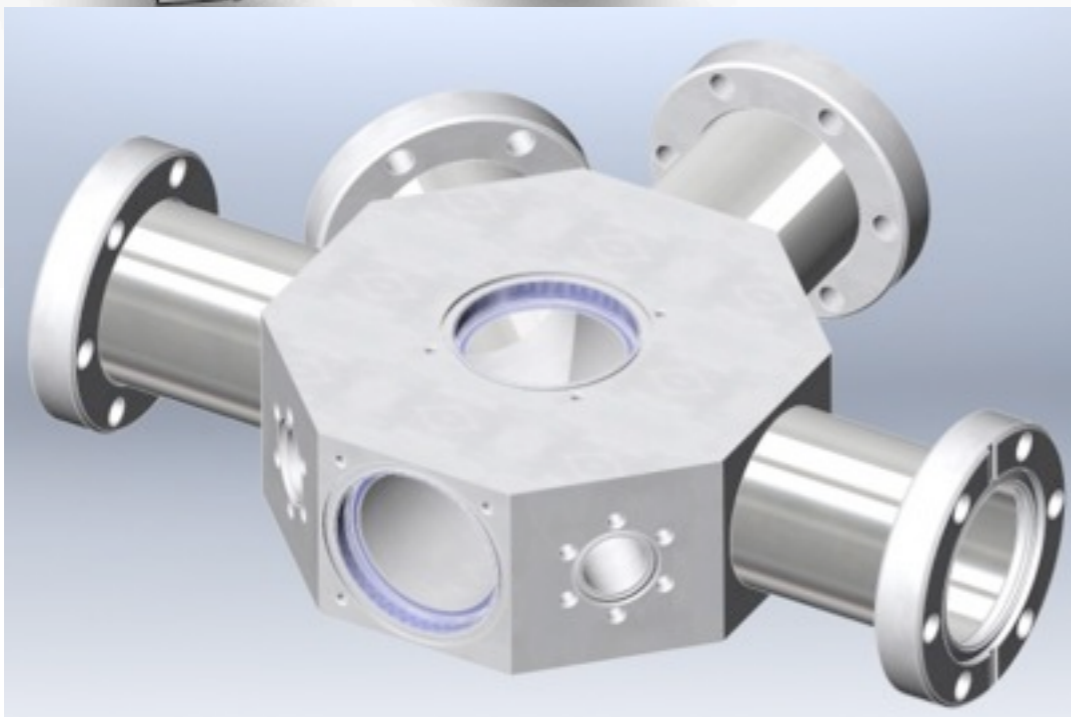
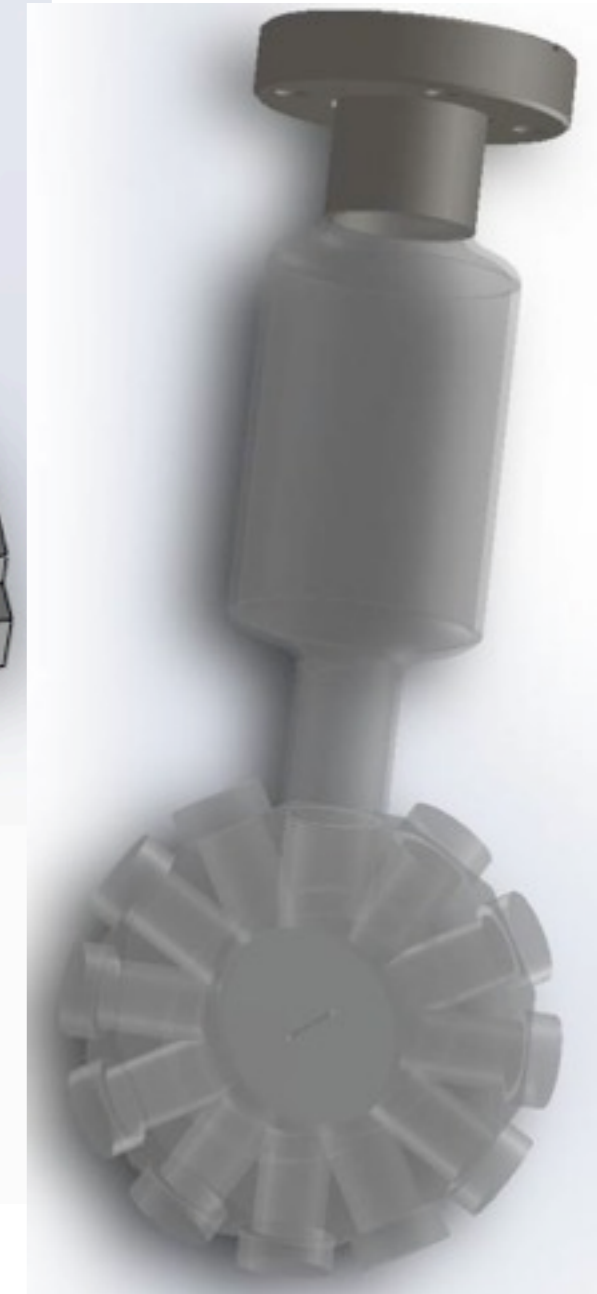
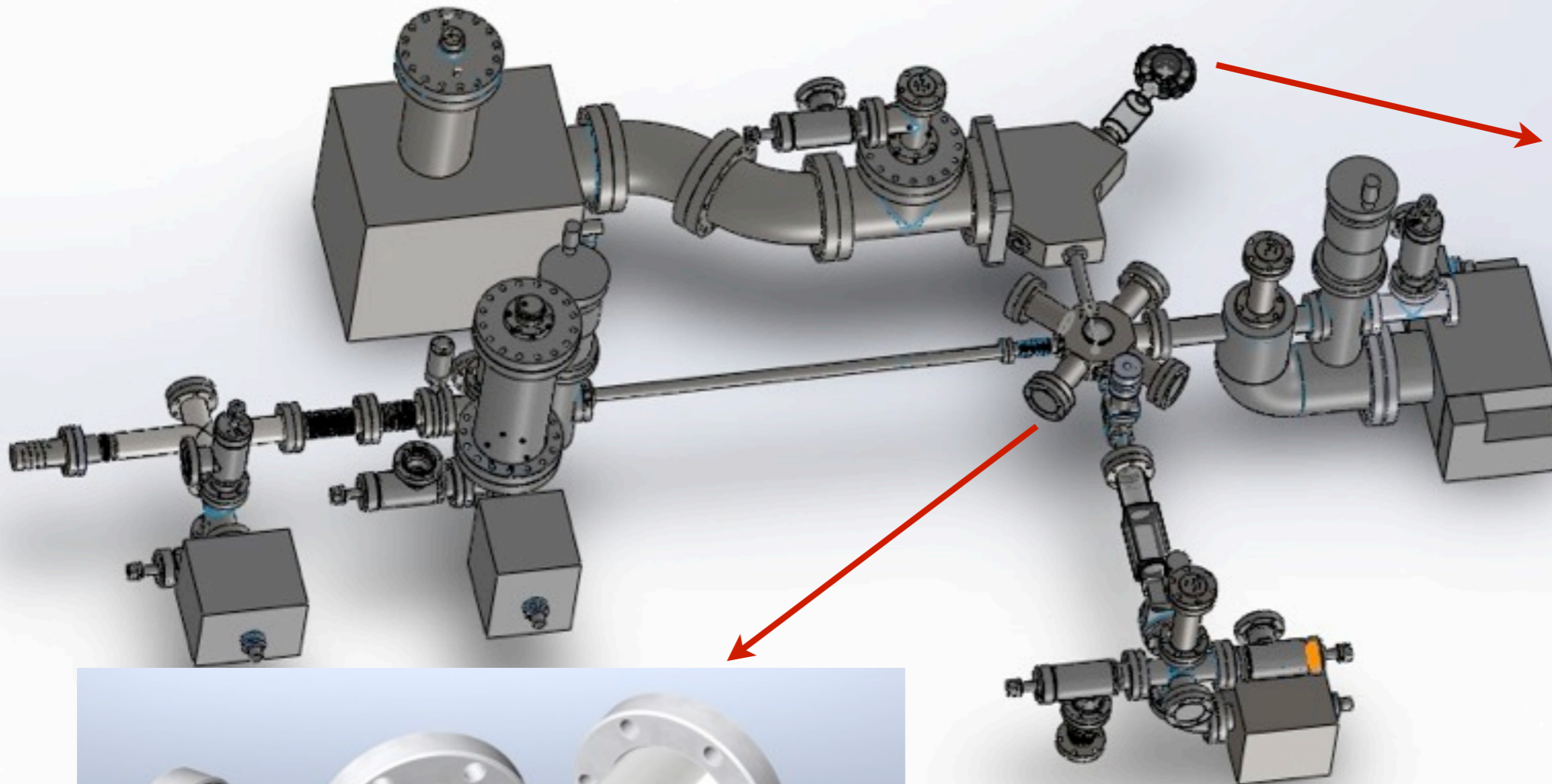
Outlook: Rectify the flux



Detection of a vortex prepared in isolated 4-site plaquettes



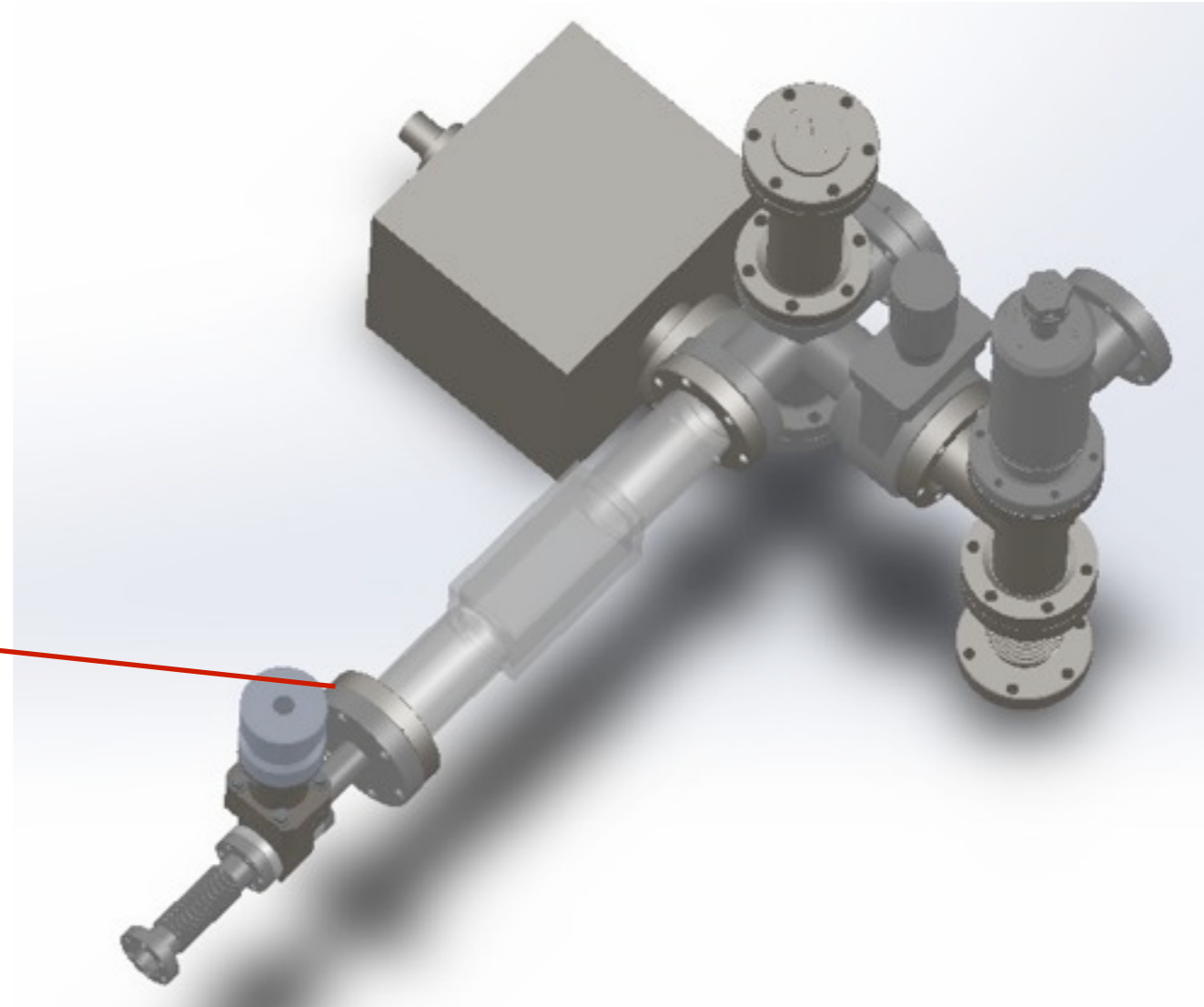
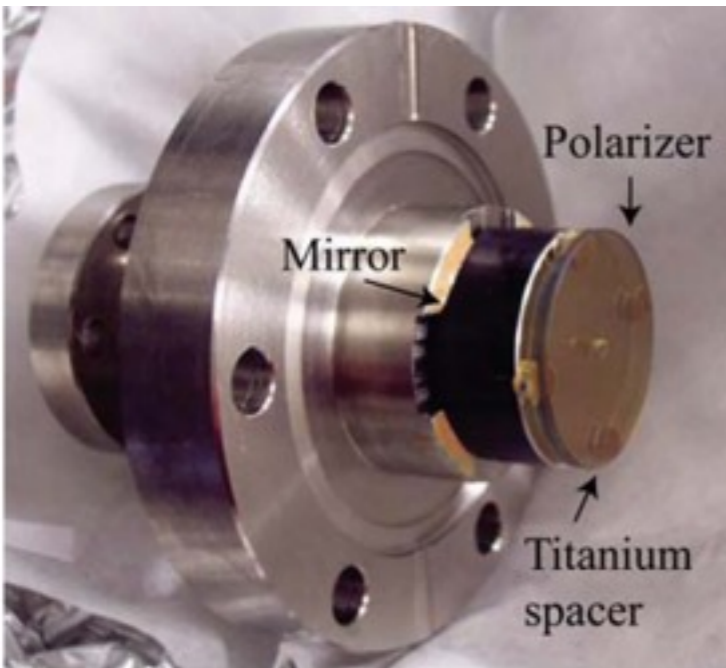
Ongoing Project: Fermi-Fermi Mixture in OLs





Ongoing Project: Fermi-Fermi Mixture in OLs

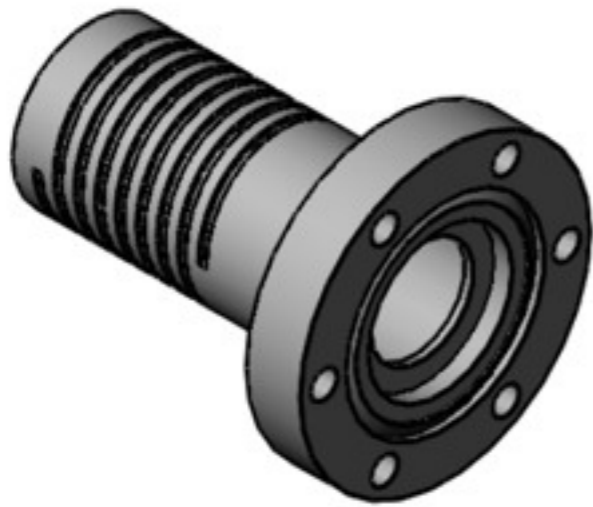
Vacuum design——2D+MOT for K40



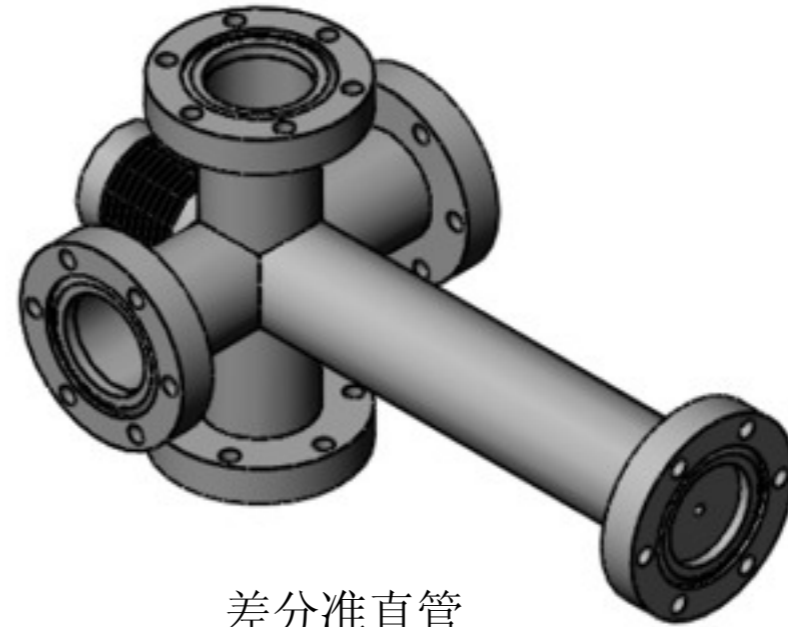


Ongoing Project: Fermi-Fermi Mixture in OLs

Vacuum design—Collimated Li6 atom beam



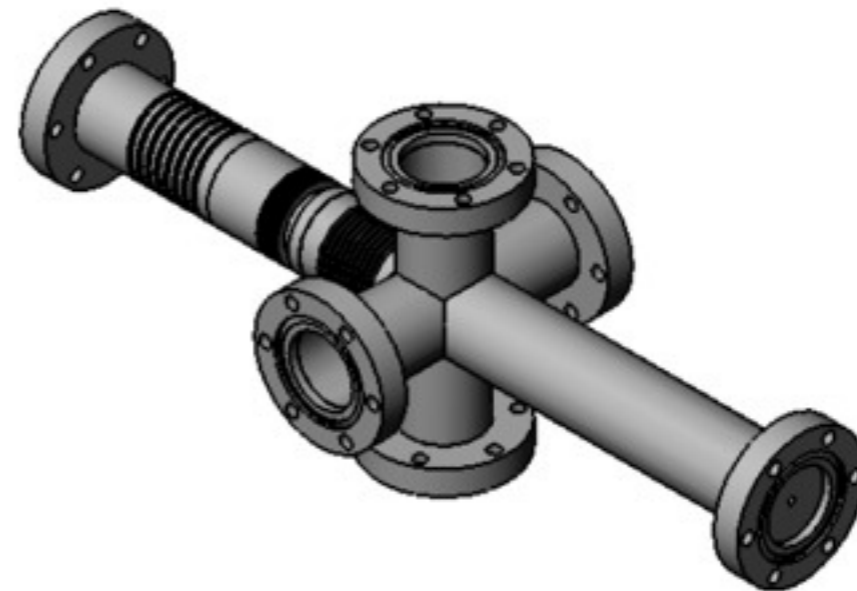
Li烤炉



差分准直管



回流炉



输出准直后的高通量Li原子束

