



Transport properties of recent discovered Weyl semimetals

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Ziquan Lin, Junfeng Wang



Introduction of Our Lab



Sample synthesis (single crystal growth)



Physical properties measurement

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OUTLINE

- Introduction of Weyl Semimetals
 - Band structure: surface Fermi arcs and bulk Weyl nodes
 - Chiral anomaly and other unique properties
- Theoretical Proposals of Weyl Semimetals
- Observation of the Weyl Semimetals in non-central-symmetric compounds
 - Ag_2Se
 - TaAs and Isostructural Compounds
 - Band structure calculation
 - ARPES
 - Transport properties
- Summary and Plan

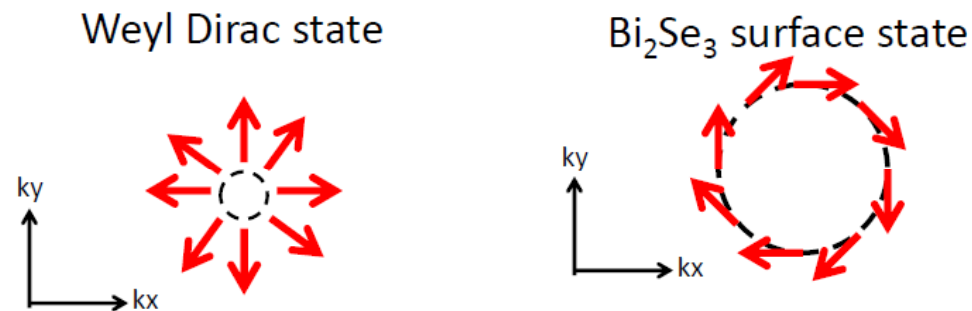


Weyl fermion: massless “Half” Dirac fermion

The Chirality of a massless fermion can be defined as:



A Dirac fermion is chiral symmetric: it can be seen as the sum of a pair of Weyl fermion

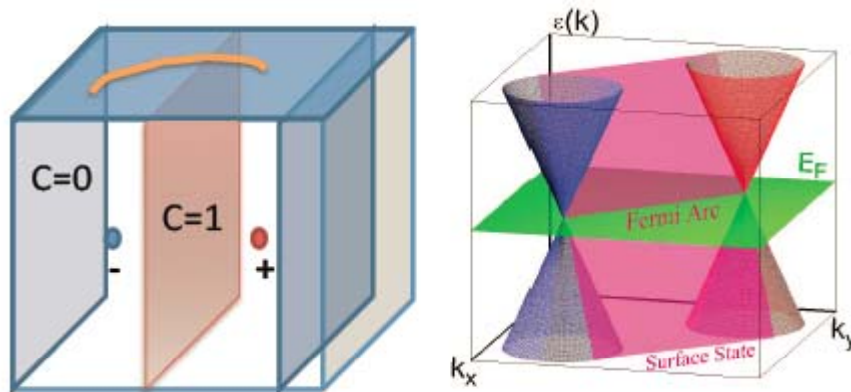


No Weyl fermion has been confirmed in high energy physics.
However, we can study it in condensed matter physics



Weyl semimetal: zero-gap semimetal with certain pairs of Weyl nodes

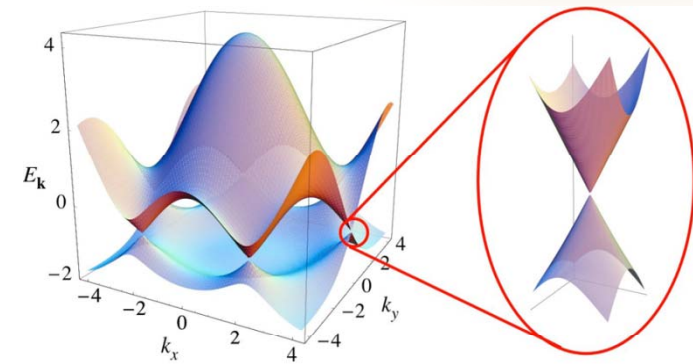
Weyl semimetal



$$H = v\vec{\sigma} \cdot \vec{k}$$

Weyl semimetal has a pair of Weyl nodes with opposite chirality

Graphene



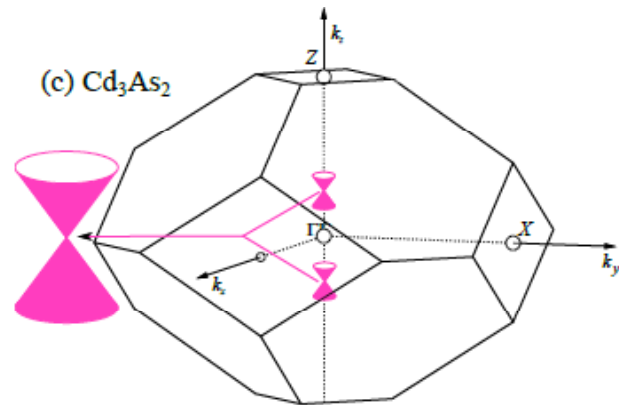
Band structure of graphene has six Dirac points

- Pairs of bands crossing at certain points (like graphene)
- The Weyl nodes are spin single degenerate. (spin textures are different)
- Fermi surface surrounding the Weyl nodes are topologically non-trivial
- The Weyl nodes with opposite chirality are connected by Fermi arcs on surfaces

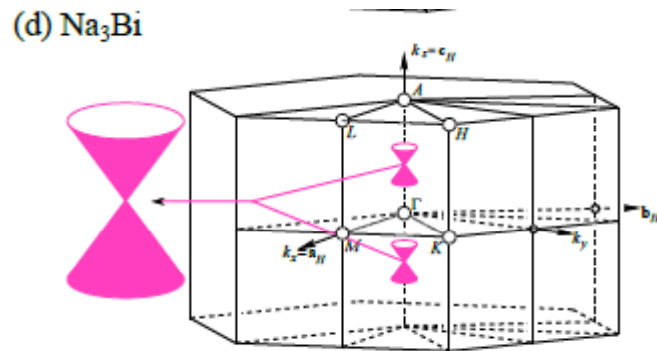
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Comparison: 3D Dirac Semimetals



- Dirac nodes along the high-symmetry lines
- Protected by rotation symmetry
- 3D analog of graphene



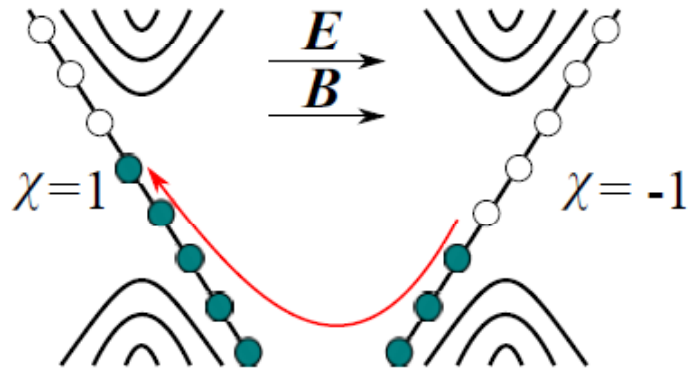
Difference with Weyl points

- Degeneracy
- Spin texture near the nodes
- Position in k space

Realization of a Weyl semimetal needs breaking a time-reversal symmetry or inversion symmetry



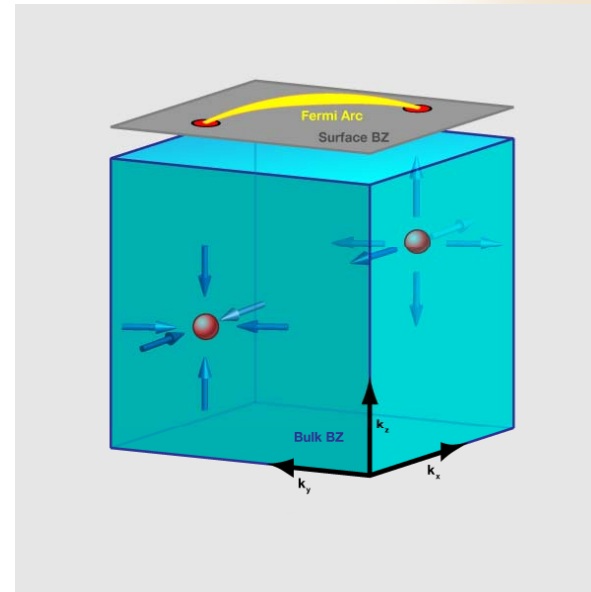
Chiral Anomaly: Negative MR



In a large field:
Beyond QL: 1D chiral anomaly
Charge pumping effect:

$$DOS \propto \vec{E} \cdot \vec{B} \propto \sigma$$

Linear positive longitudinal magneto-conductance



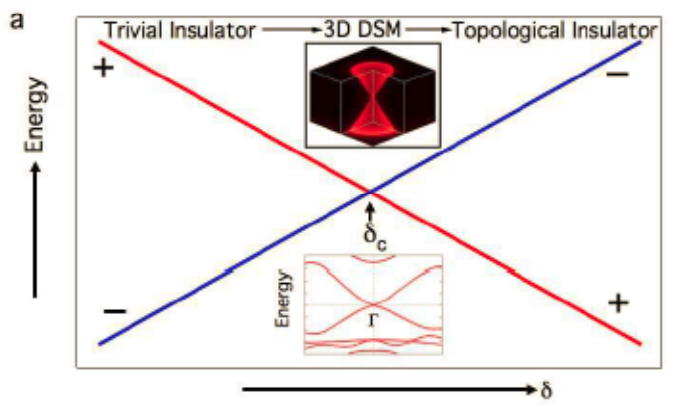
More general case:
Negative MR

D. Son & B. Spivak PRB 88 104412 (2013)

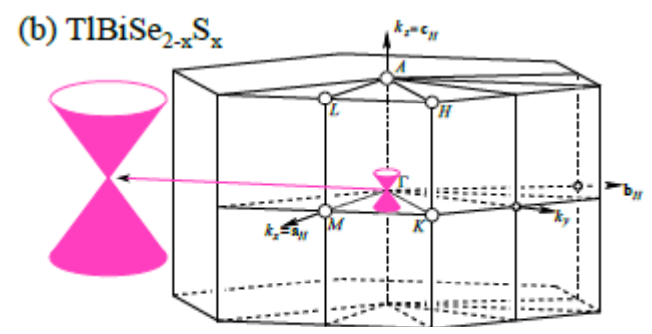
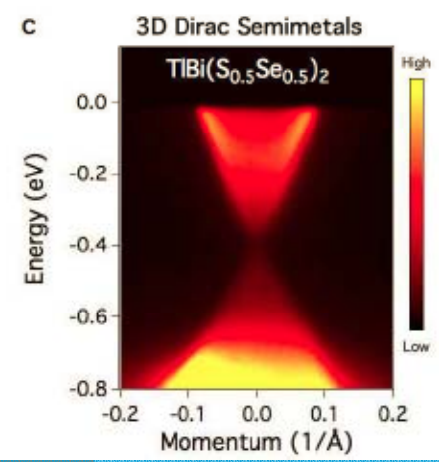
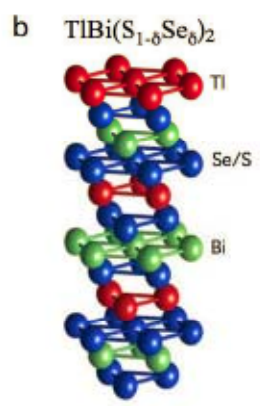
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How to Realize a Weyl Semimetal (Proposal I)

◆ A QCP realized by tuning the SOC interaction.



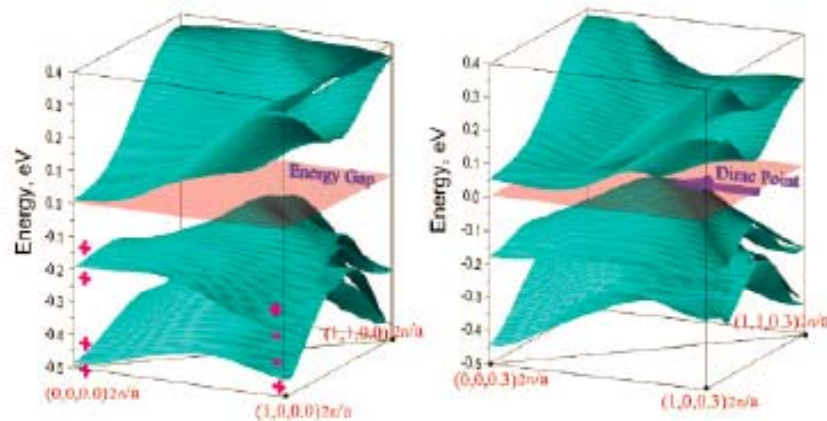
Example: $\text{TlBi}(\text{S}_{1-x}\text{Se}_x)_2$, $\text{Bi}_{1-x}\text{Sb}_x$, $\text{Bi}_{2-x}\text{In}_x\text{Se}_3$



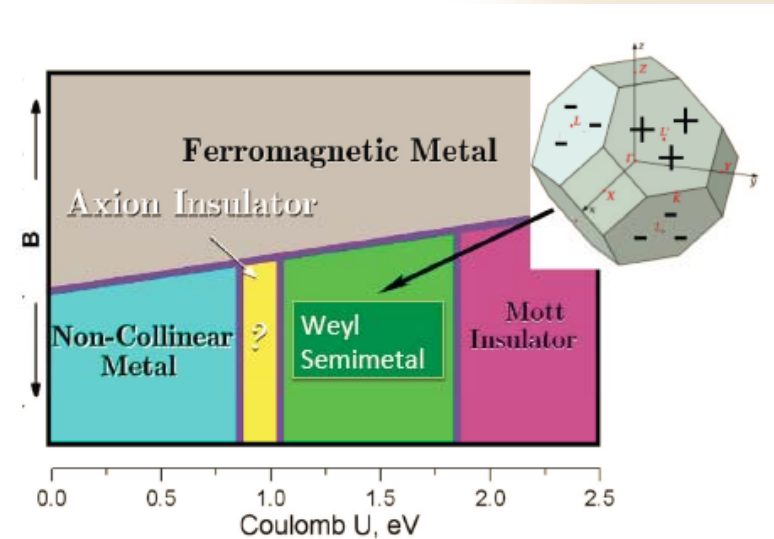
Disadvantage: doping induced disorder, need magnetic ground state

How to Realize a Weyl Semimetal (Proposal II)

- ◆ A magnetic ordered ground state with strong spin-orbital coupling.



Band structure of $Y_2Ir_2O_7$
 X. Wan *et al.* PRB 83 205101(2011)



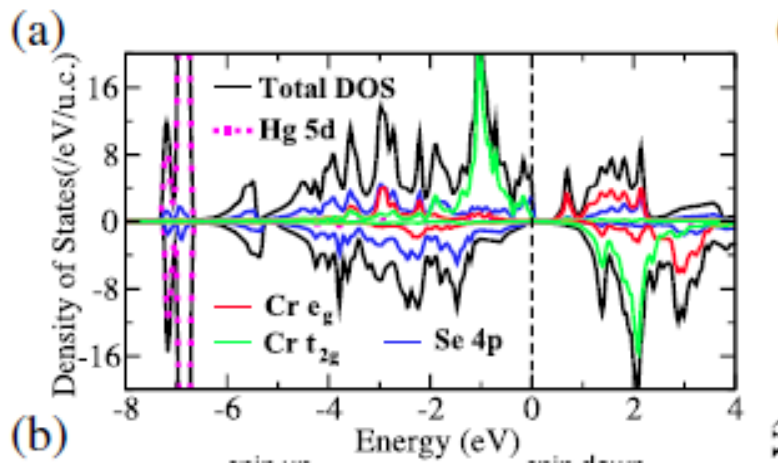
Calculated phase diagram of $Y_2Ir_2O_7$ with interaction strength
 X. Wan *et al.* PRB 83 205101(2011)

Example: $Y_2Ir_2O_7$ (geometric frustrated)
 Disadvantage: difficulty of sample growth, unknown magnetic structure

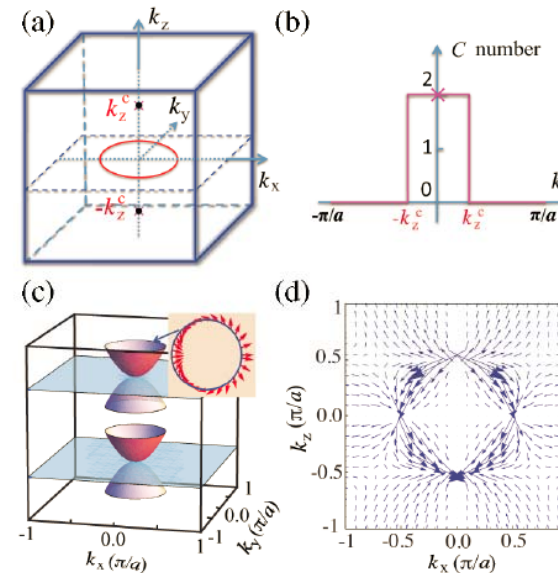


How to Realize a Weyl Semimetal (Proposal II)

- ◆ A magnetic ordered ground state with strong spin-orbital coupling.



Band structure of HgCr_2Se_4
G. Xu *etal.* PRL 107 186806(2011)



Weyl nodes in HgCr_2Se_4
G. Xu *etal.* PRL 107 186806(2011)

Example: HgCr_2Se_4 (FM spinel)

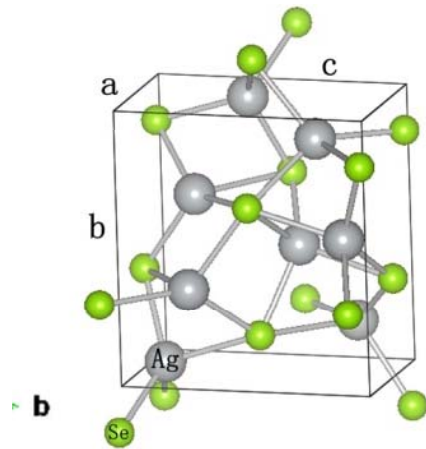
Disadvantage: difficulty of sample growth, magnetic domain

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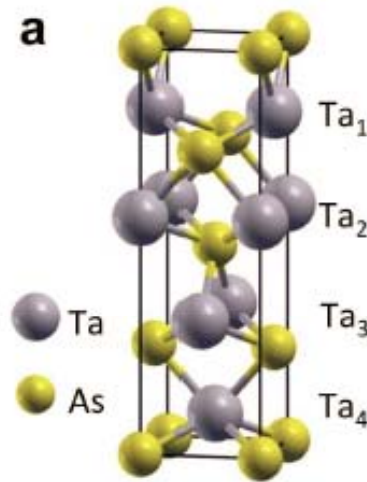
How to Realize a Weyl Semimetal (Proposal III)



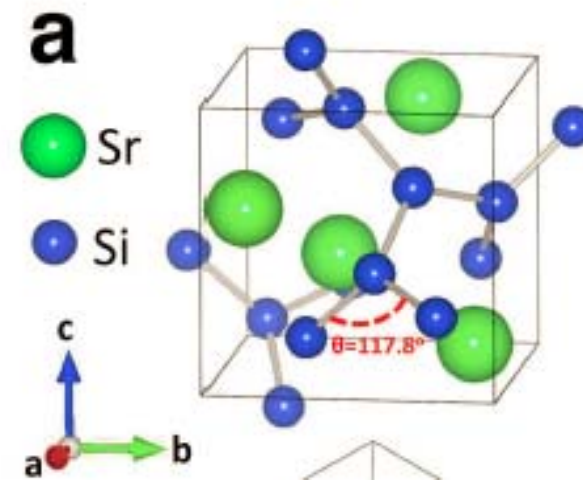
◆ non-central-symmetric semimetal with proper SOC



Ag_2Se : $P2_12_12_1$



TaAs : $I4_1md$



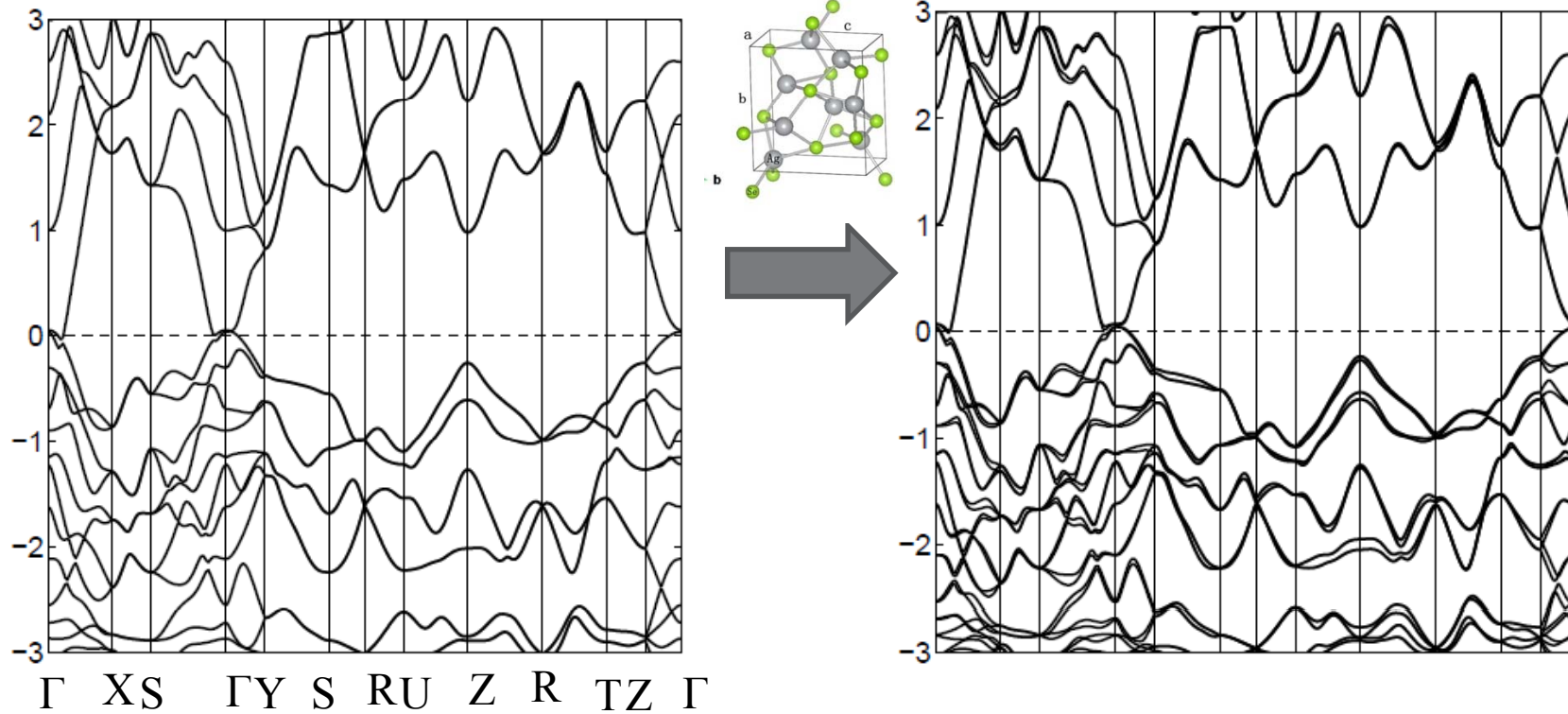
SrSi_2 : $P4_332$

Difficulty:

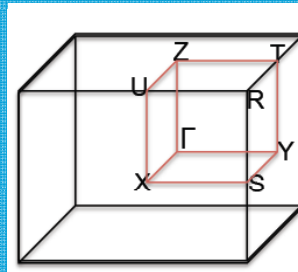
- choose right systems (need more calculations, structural information)
- single crystal growth (thermally unstable low-T phases)

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Band Structure of Ag_2Se



With no SOC: gapless in high-symmetry points



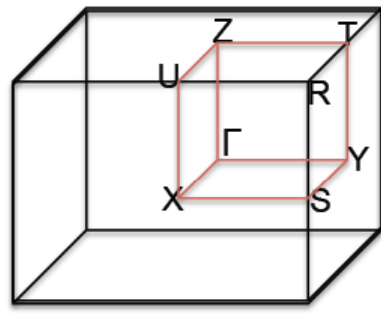
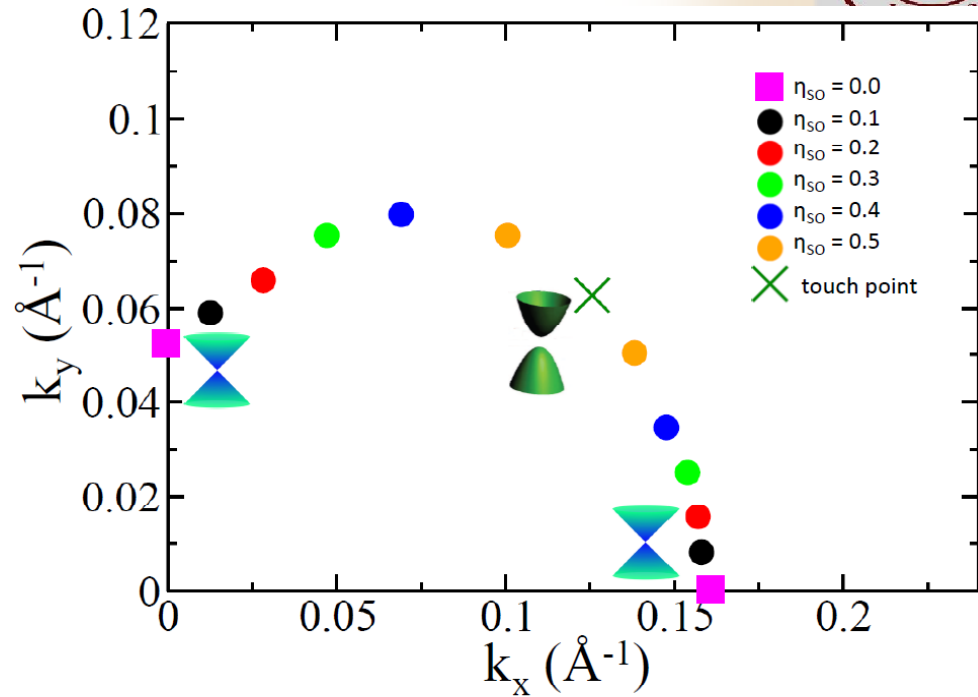
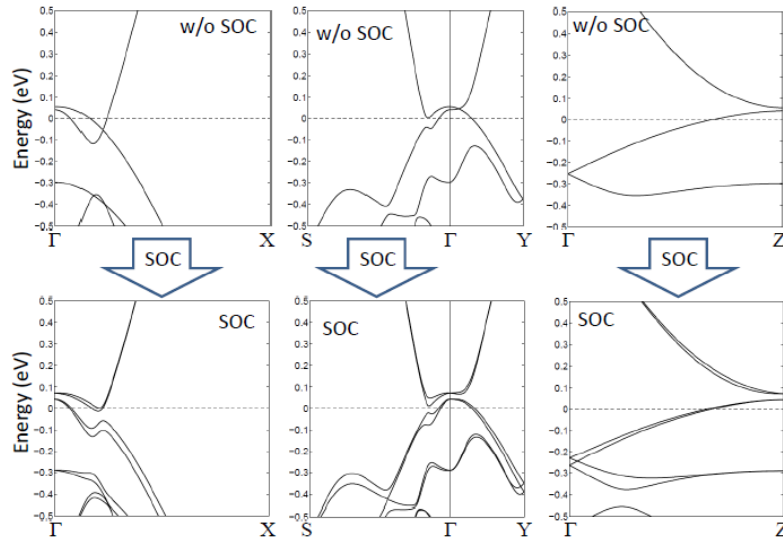
With SOC: small gap in all high-symmetry points and lines

Space group $P2_12_12_1$

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Weyl Nodes in Ag_2Se



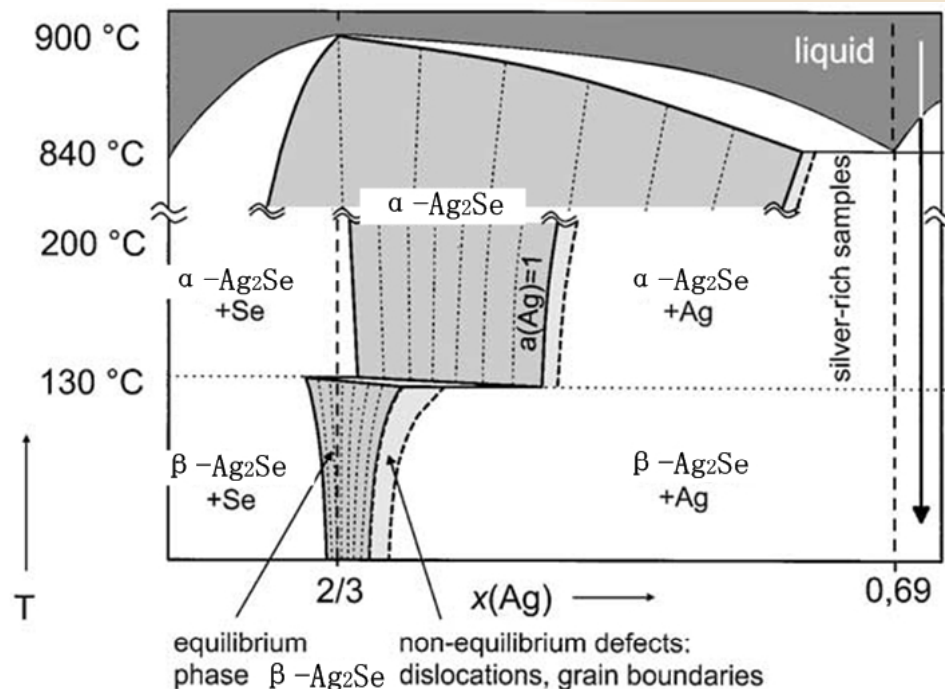
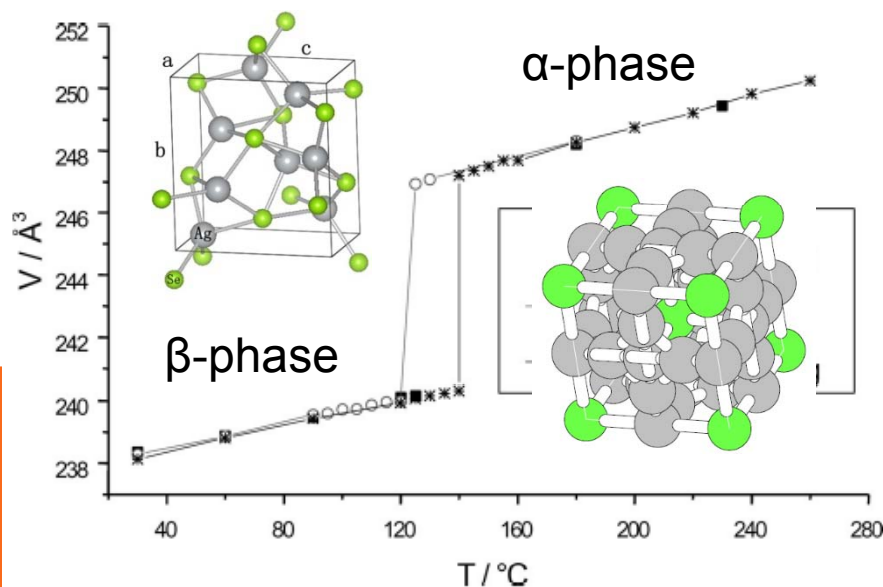
- 8 Weyl nodes born with small SOC
- Move in k -space with SOC increasing
- Annihilate at certain SOC

Difficulty for Growing Single Crystals

β -Ag₂Se



- High-T α -Phase: Super-Ionic, Cubic $Im-3m$, $a \sim 5 \text{ \AA}$
- Low-T β -Phase: Non-central-symmetric Orthorhombic $P2_12_12_1$
- β -Phase: $a \sim 4 \text{ \AA}$, b and $c \sim 7 \text{ \AA}$
- 1st Order Structural Phase Transition at $133 \text{ }^\circ\text{C}$ with Giant Volume Change



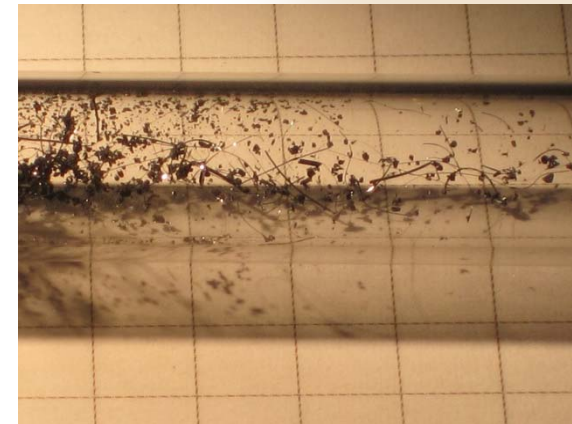
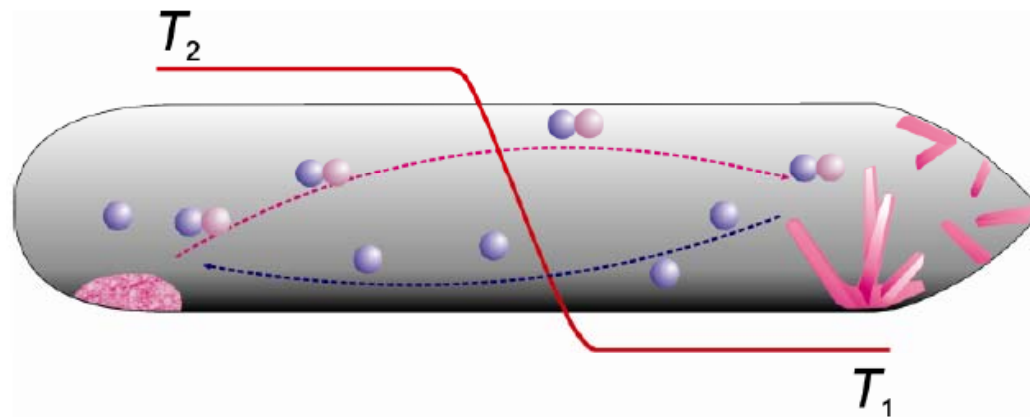
G. Beck and J. Janek *Physica B* 308-310: 1086 (2001)

Bias : We cannot grow single crystals of such thermally unstable compound.

H. Billetter and U. Ruschewitz *Z. Anorg. Allg. Chem.* 634: 241 (2008)

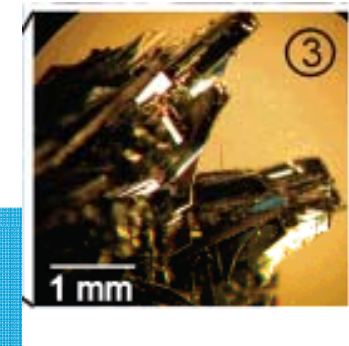


Single Crystal Growth: Vapor Transfer



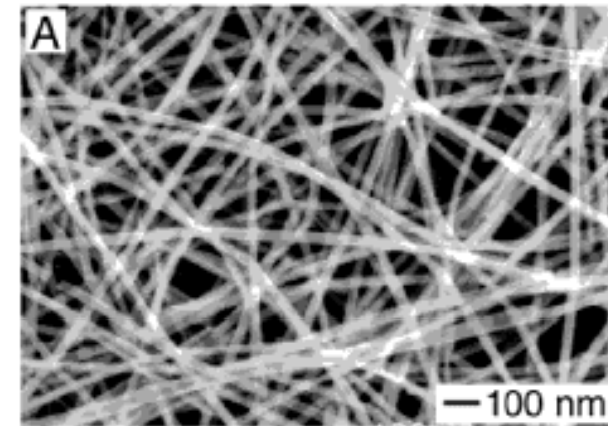
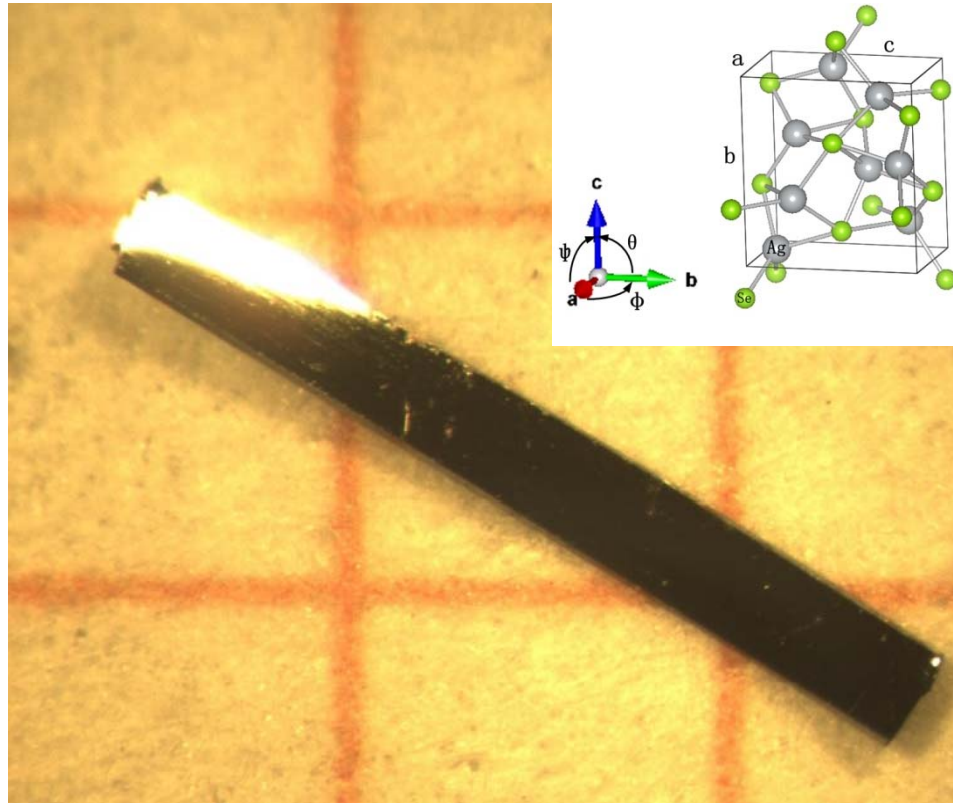
Ag₂Se in a quartz tube

- Forming crystals via gas phase
- Can be self-selected transfer or via agent
- Key point: agent, temperature, direction, “luck”
- Advantage: growing at low-T, clearer than flux-method, good for high-vapor-pressure elements



Black phosphorus via VT

Macro-size Single Crystals Synthesis of Ag_2Se

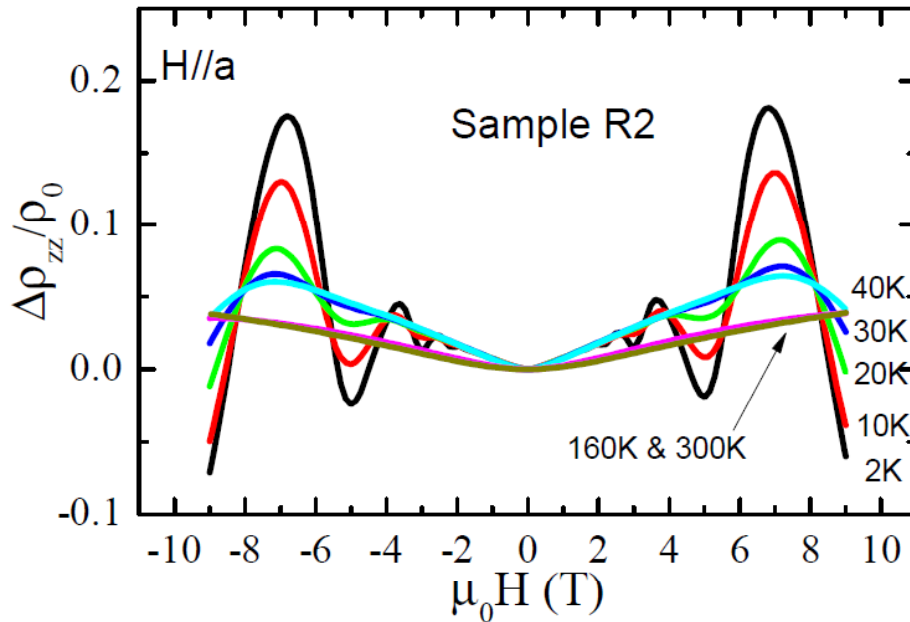


B. Gates et. al. *JACS* 123: 1150 (2001)

For comparison, only nano-size crystals were synthesized before.

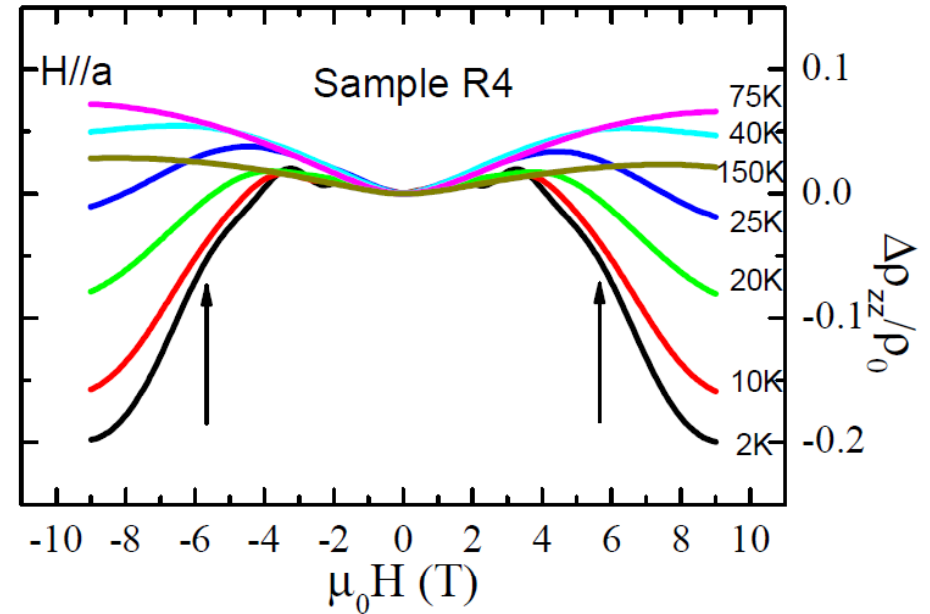
- Crystals Grown Via Vapor Transfer Method
- Ribbon-Like Shape
- Size: up to $10 \times 0.5 \times 0.5$ mm
- Confirmed by Single Crystal XRD

Longitudinal MR: Strong SdH Oscillations +Negative LMR



Large n

- Strong SdH Oscillations
- 1st Landau Level at 7T
- Negative LMR beyond 1st LL
- Oscillations and Negative LMR fade out when $T \uparrow$

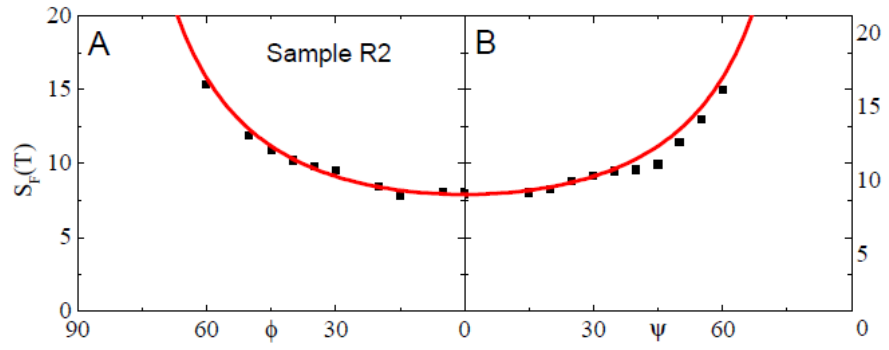
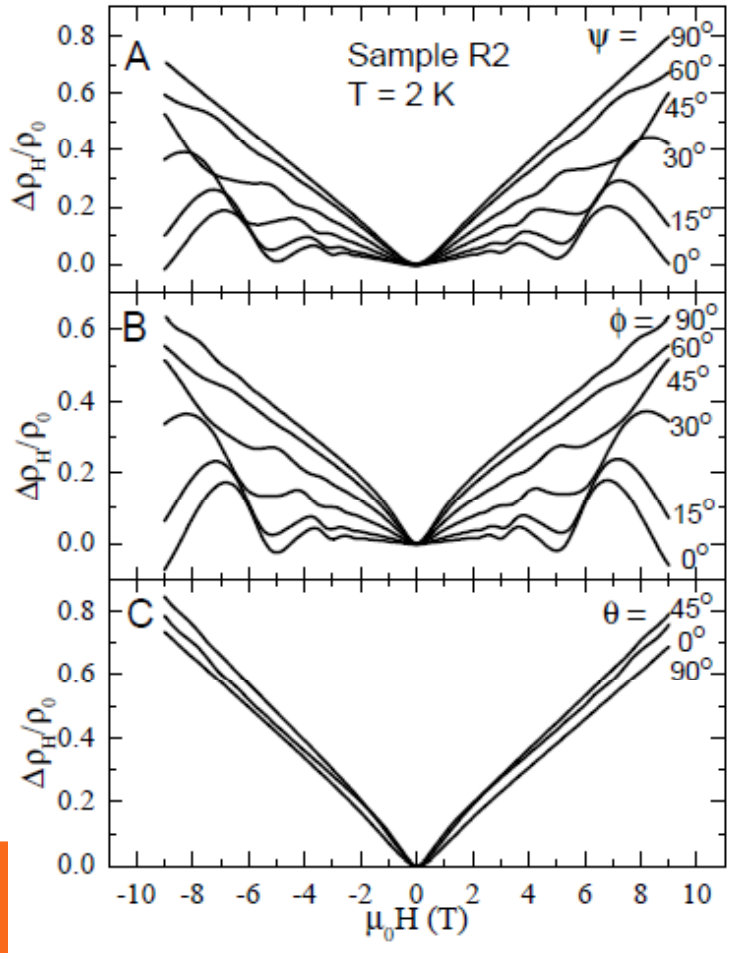


Small n

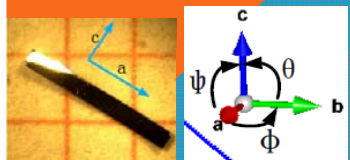
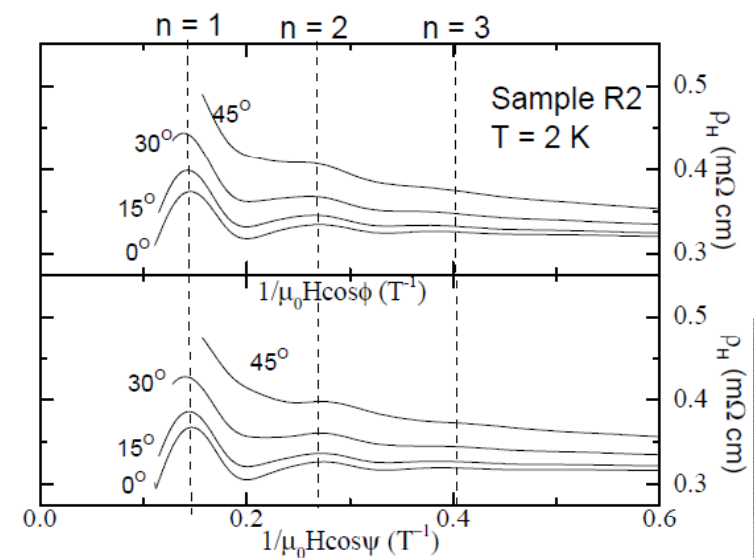
- Relatively Weak Oscillations
- 1st LL at 3T
- Clear Negative LMR beyond 1st LL
- Anomaly at 6 T

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Highly Anisotropic Fermi Surface



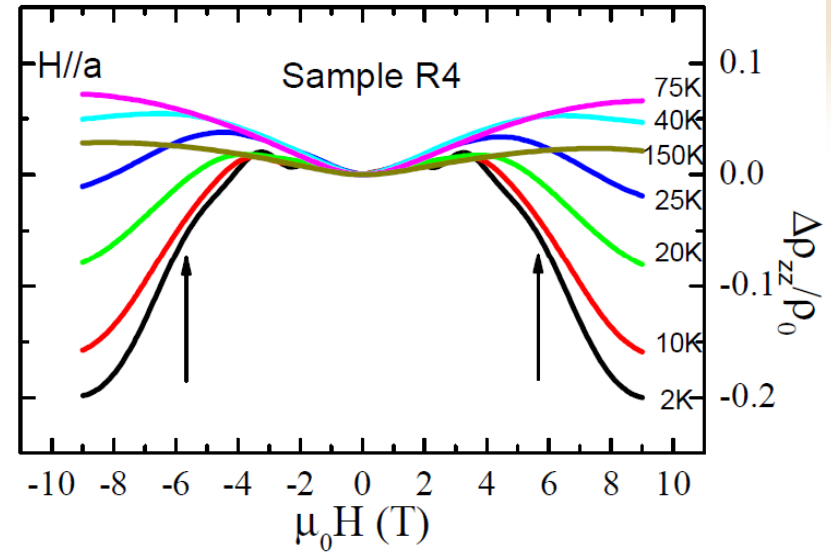
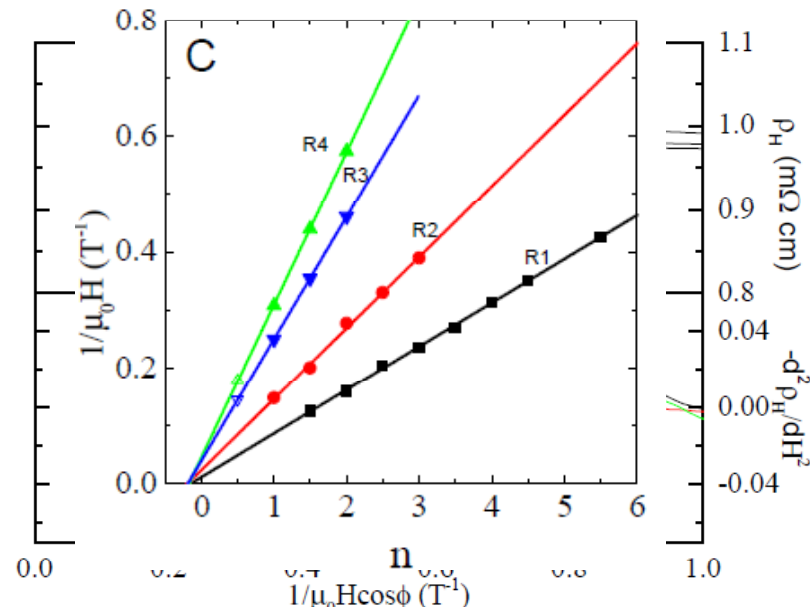
FS Follows $\frac{1}{\cos\phi}$ or $\frac{1}{\cos\psi}$



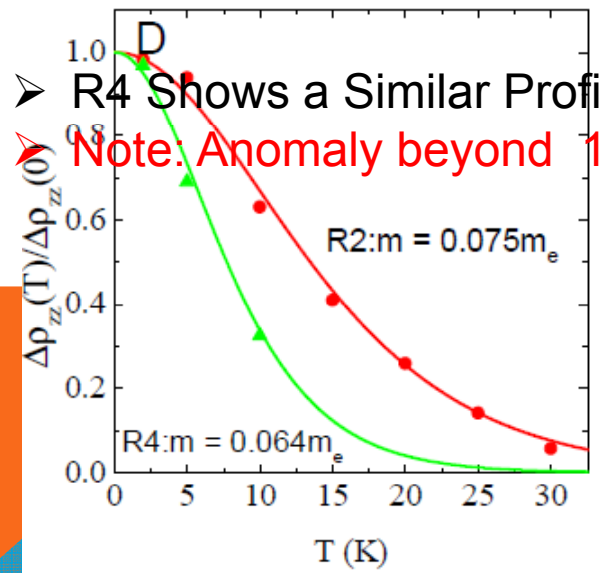
- MR from a to b/c are similar
- Invariant from b to c

FS : Cylinder or Highly Anisotropic Ellipsoid

Dirac Fermions with Tiny Surface



➤ R4 Shows a Similar Profile
 Note: Anomaly beyond 1st LL



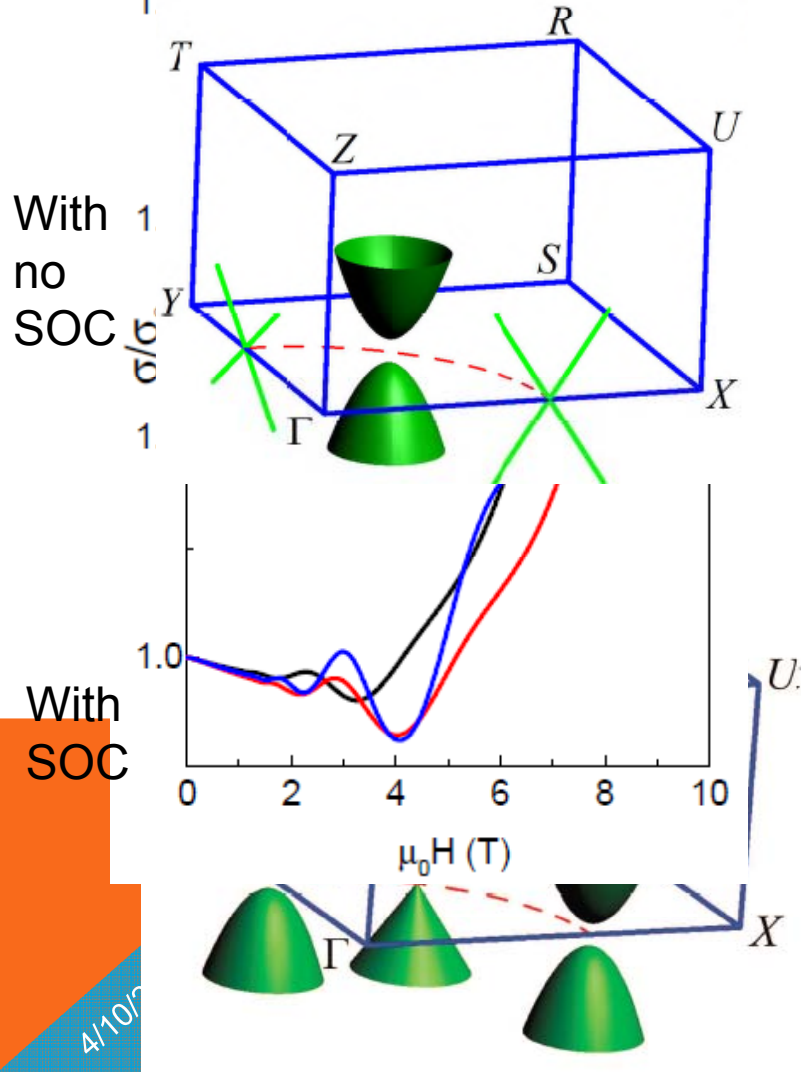
- FS: 13.3T~4T
- Onsager Phase $\gamma = -0.15 \pm 0.05$
- $m < 0.1m_e$
- $v_F \sim 2 \times 10^5 \text{ ms}^{-1}$
- Consistent with a Linear Dispersive Band near Fermi Level
- The Anomaly at $n = 1/2$



Negative LMR: Chiral Anomaly beyond QL

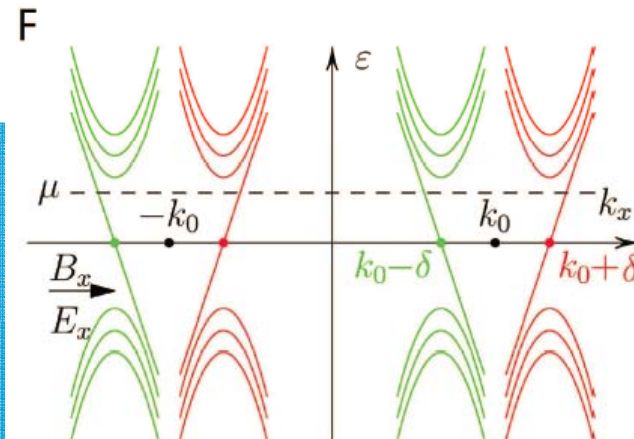
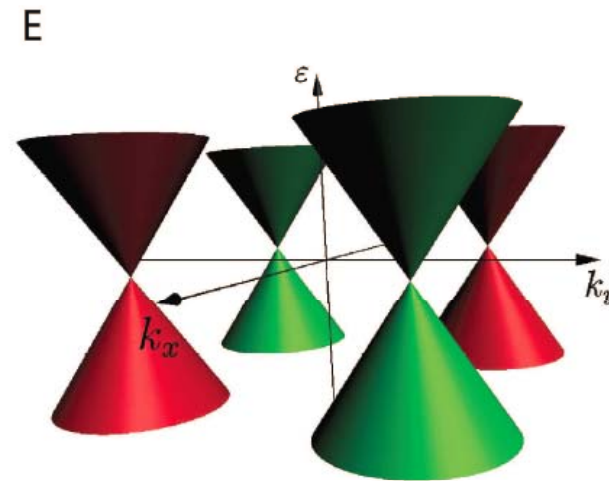
$P2_12_12_1$ Leads Symmetry

Protected Weyl Nodes due to SOC

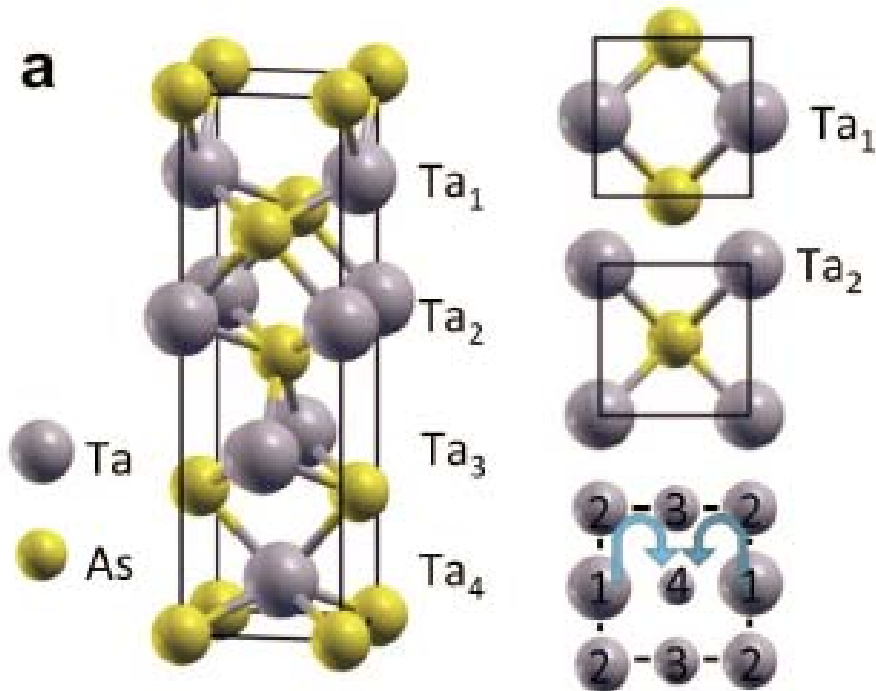


Charge Pumping Effect due to Chiral Anomaly beyond QL :

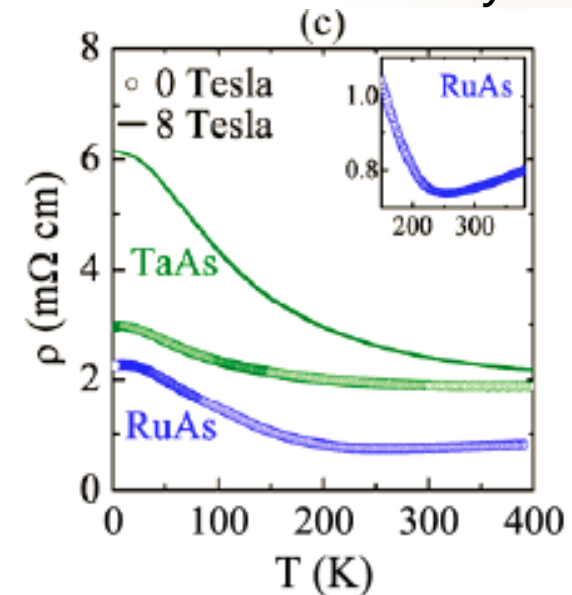
$$DOS \propto \vec{E} \cdot \vec{B} \propto \sigma$$



TaAs: a Semimetal with Non-central Symmetric Structure

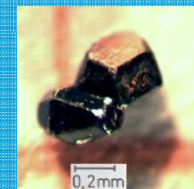


Electric properties of polycrystals are not noteworthy



B. Sapiro *et al.* 2012

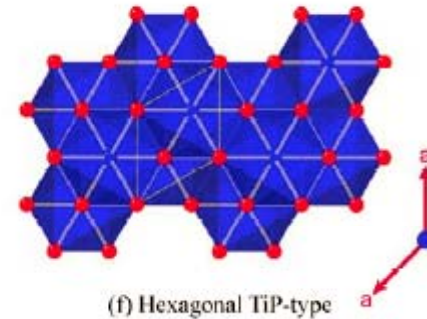
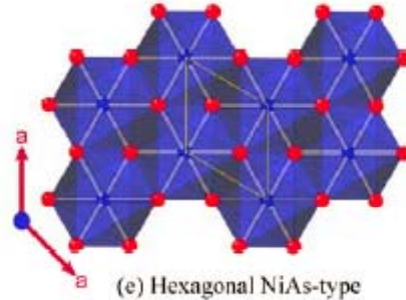
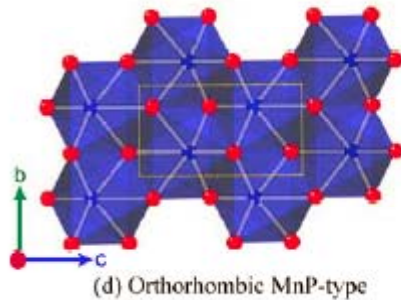
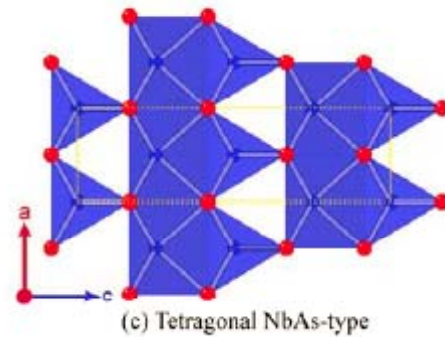
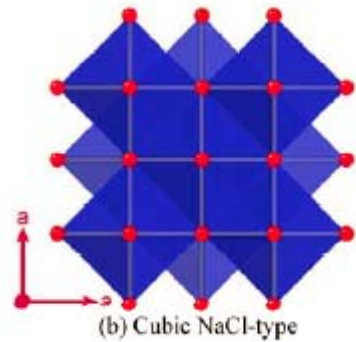
- One Ta atom in a 6-atom As prism
- Space group: $I4_1md$
- Previous study on polycrystals show semiconductor/semimetal behaviors
- Single crystals grown via vapor transfer method



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TaAs, NbAs, TaP and NbP: Relatively Structure in Pnictides



Periodic Table of the Elements

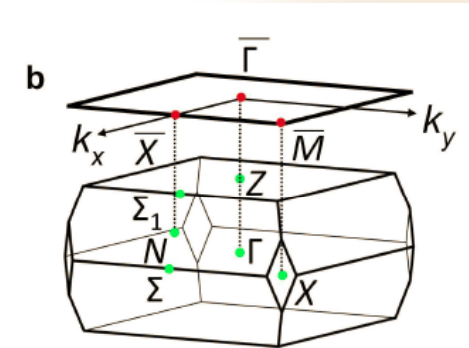
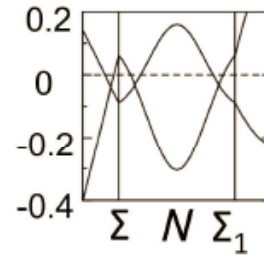
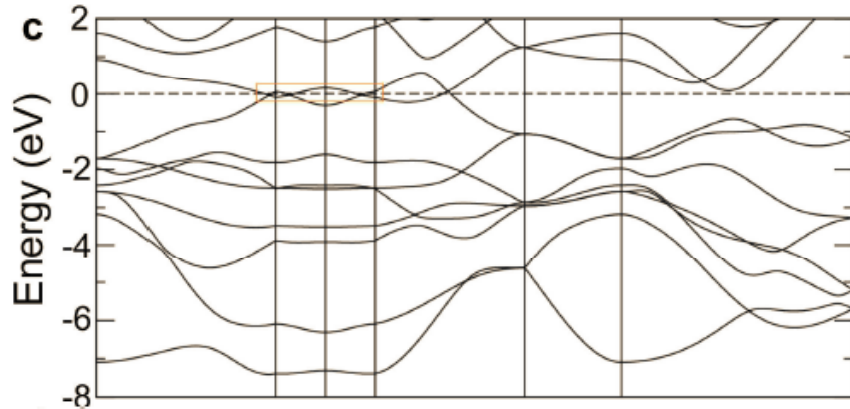
Representative (main group) elements																		Representative (main group) elements													
1 H 1.0079																	2 He 4.003														
3 Li 6.941	4 Be 9.012																	5 B 10.811	6 C 12.011	7 N 14.007	8 O 15.999	9 F 18.998	10 Ne 20.180								
Transition metals																															
11 Na 22.990	12 Mg 24.305																	13 Al 26.982	14 Si 28.086	15 P 30.974	16 S 32.065	17 Cl 35.453	18 Ar 39.948								
19 K 39.098	20 Ca 40.078	21 Sc 44.956	22 Ti 47.88	23 V 50.942	24 Cr 51.996	25 Mn 54.938	26 Fe 55.845	27 Co 58.933	28 Ni 58.69	29 Cu 63.546	30 Zn 65.39	31 Ga 69.723	32 Ge 72.61	33 As 74.922	34 Se 78.96	35 Br 79.904	36 Kr 83.8														
37 Rb 85.468	38 Sr 87.62	39 Y 88.908	40 Zr 91.224	41 Nb 92.906	42 Mo 95.94	43 Tc 98	44 Ru 101.07	45 Rh 102.906	46 Pd 106.42	47 Ag 107.868	48 Cd 112.411	49 In 114.82	50 Sn 118.71	51 Sb 121.76	52 Te 127.60	53 I 126.905	54 Xe 131.29														
55 Cs 132.905	56 Ba 137.327	57 La 138.906	58 Ce 140.115	59 Pr 140.908	60 Nd 144.24	61 Pm 144.9126	62 Sm 150.36	63 Eu 151.964	64 Gd 157.25	65 Tb 158.925	66 Dy 162.5	67 Ho 164.93	68 Er 167.26	69 Tm 168.934	70 Yb 173.04	71 Lu 174.967															
87 Fr 223	88 Ra 226.025	89 Ac 227.028	90 Th 232.038	91 Pa 231.036	92 U 238.029	93 Np 237.048	94 Pu 244	95 Am 243	96 Cm 247	97 Bk 247	98 Cf 251	99 Es 252	100 Fm 257	101 Md 258	102 No 259	103 Lr 262															
Rare earth elements																															
Lanthanides																															
Actinides																															

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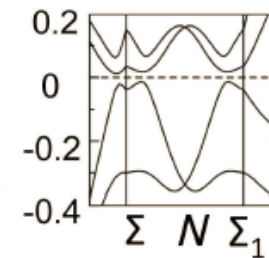
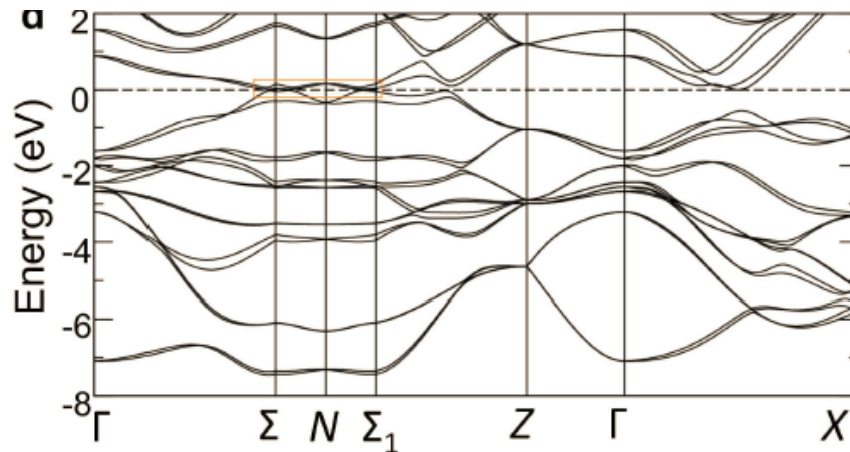
Band Structure of TaAs

With no SOC: band touching



The BZ of TaAs

With SOC: open a small gap in high symmetric lines

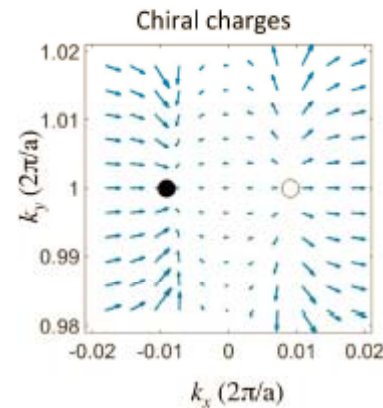
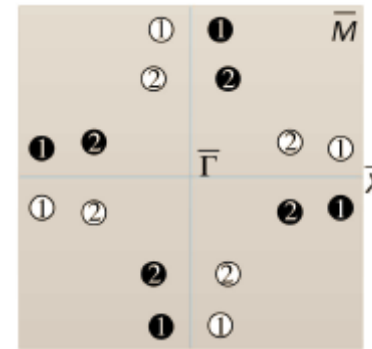
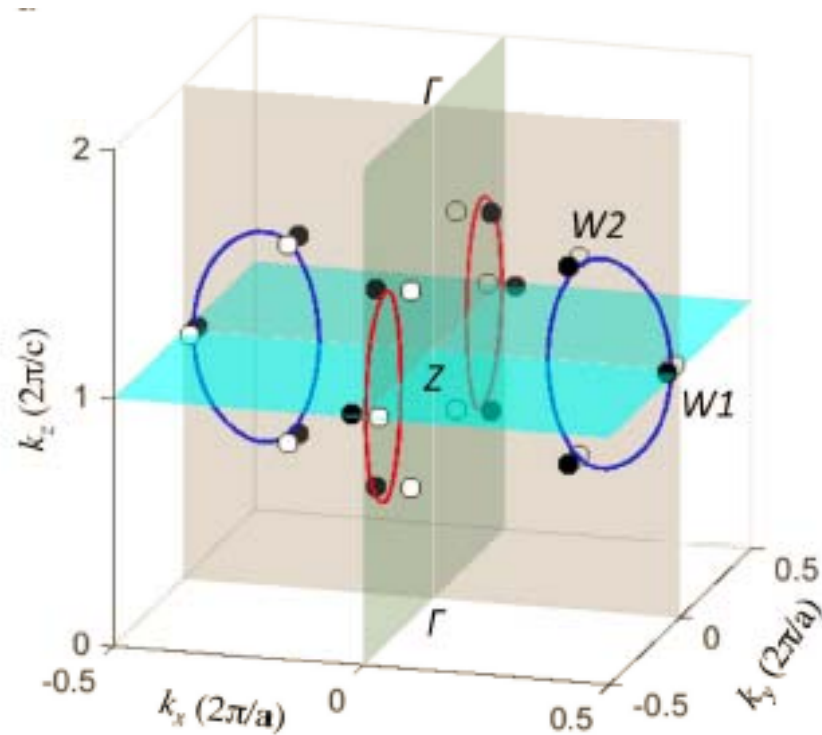


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S. Huang *etal.* Arxiv:1501.00755(2015)
Similar results in H. Weng *etal.* PRX 5 011029 (2015)



Weyl Nodes in TaAs (calculated)

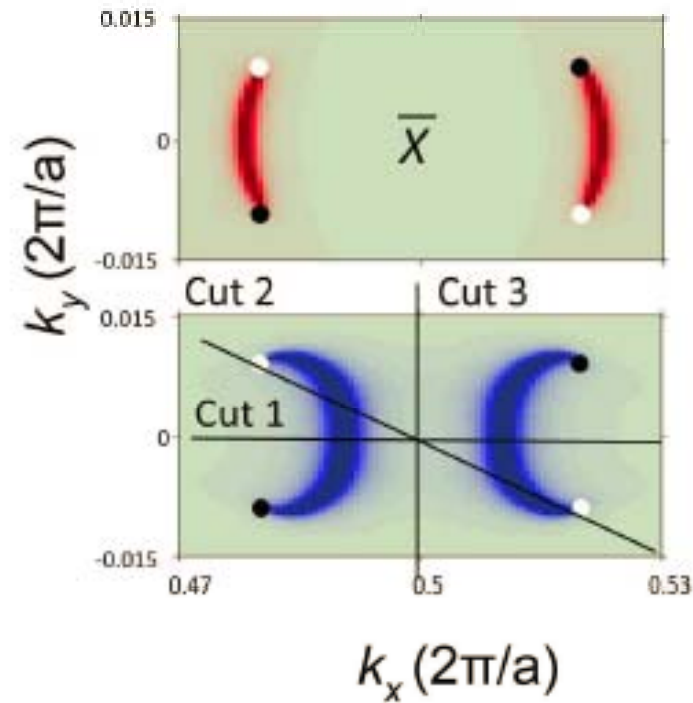
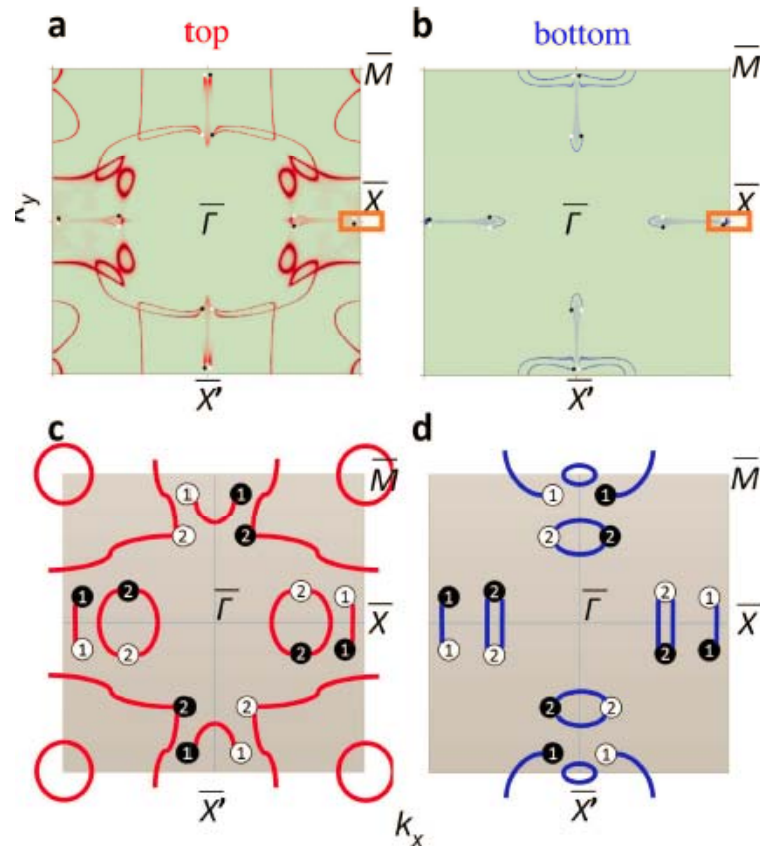


- Each line node vaporizes into six Weyl points
- 8 W1 on one plane, 16 W2 away from the plane
- Spin texture shows opposite chirality

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S. Huang *etal.* Arxiv:1501.00755(2015)
 Similar results in H. Weng *etal.* PRX 5 011029 (2015)

Surface Fermi Arc in TaAs



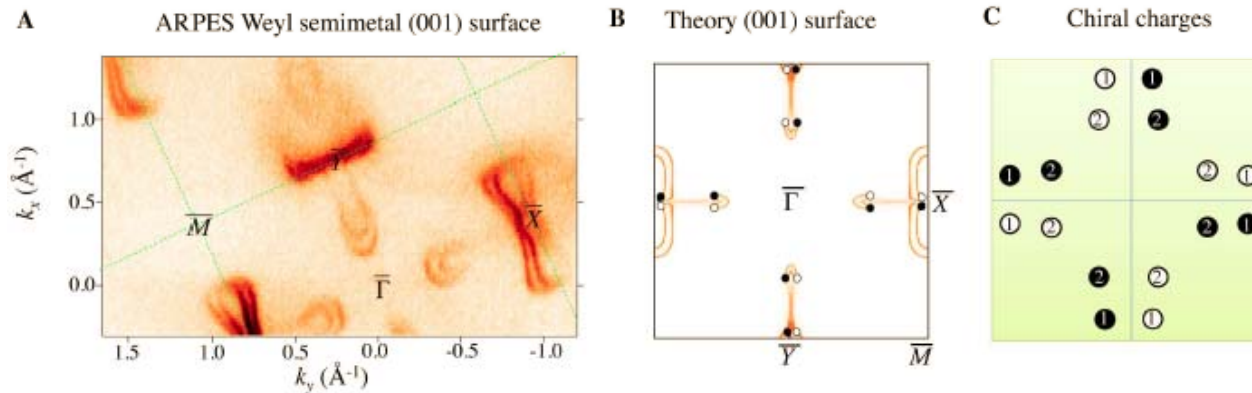
The (001) surface states on the top and bottom surface of TaAs

A close-up of the Fermi arc on the surface

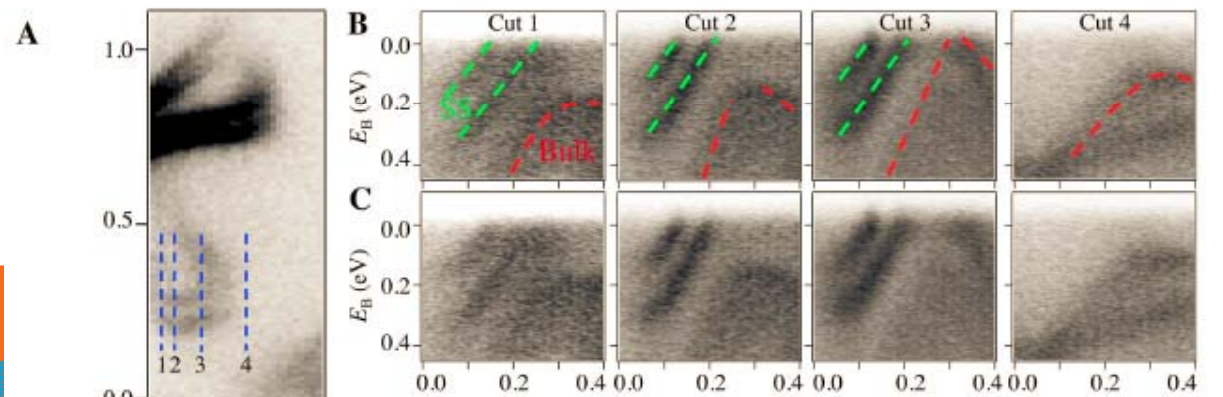
4/10/2015

S. Huang *et al.* Arxiv:1501.00755(2015)
 Similar results in H. Weng *et al.* PRX 5 011029 (2015)

ARPES: Observing the Surface Fermi Arc and Weyl Nodes in Experiment



Observed Fermi Arc surface states

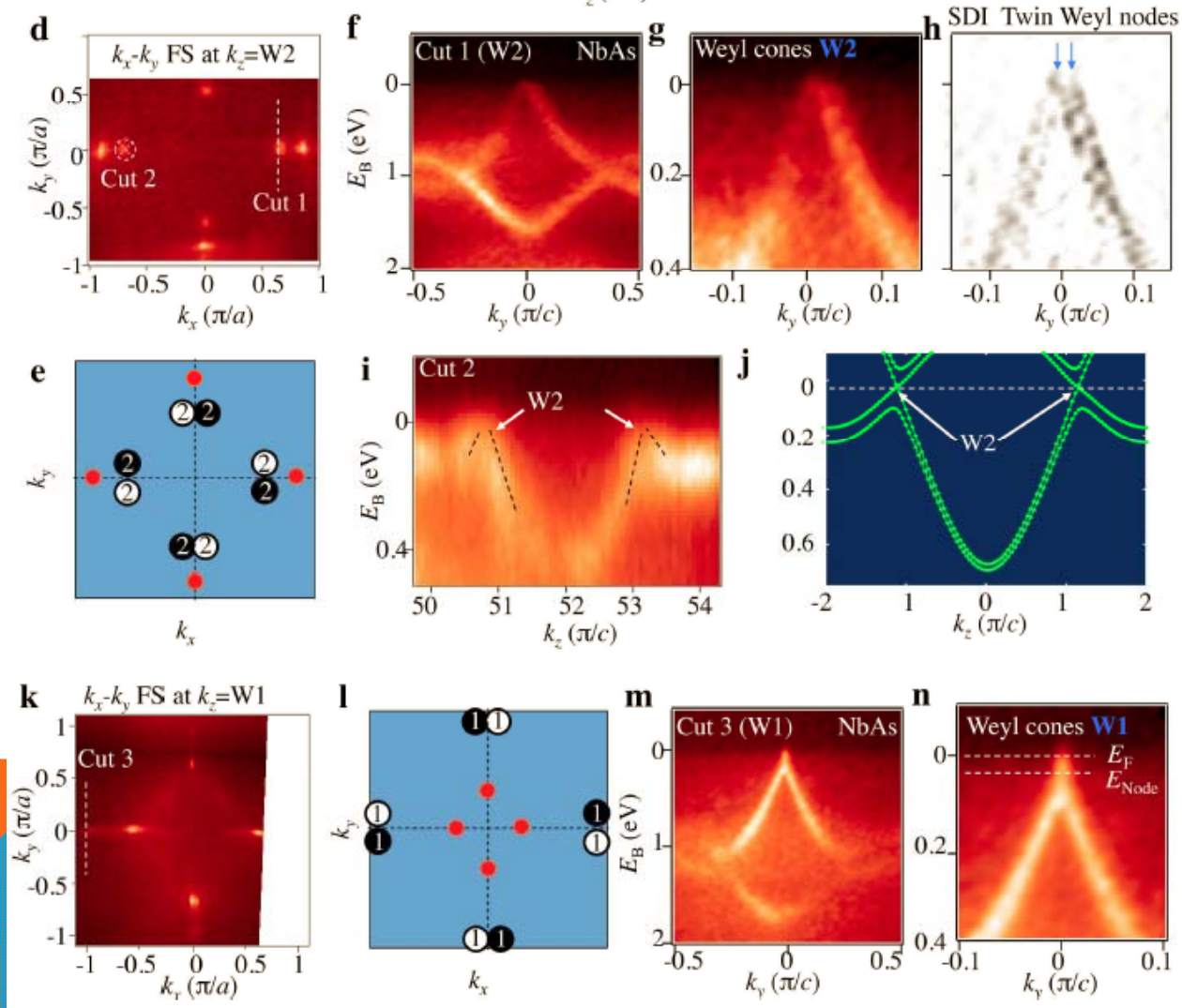


The Weyl cone rises up to meet the Fermi arcs

SY Xu *et al.* Arxiv:1501.03807(2015)
 Similar results in B.Q Lv *et al.* Arxiv: 1502.04684(2015)



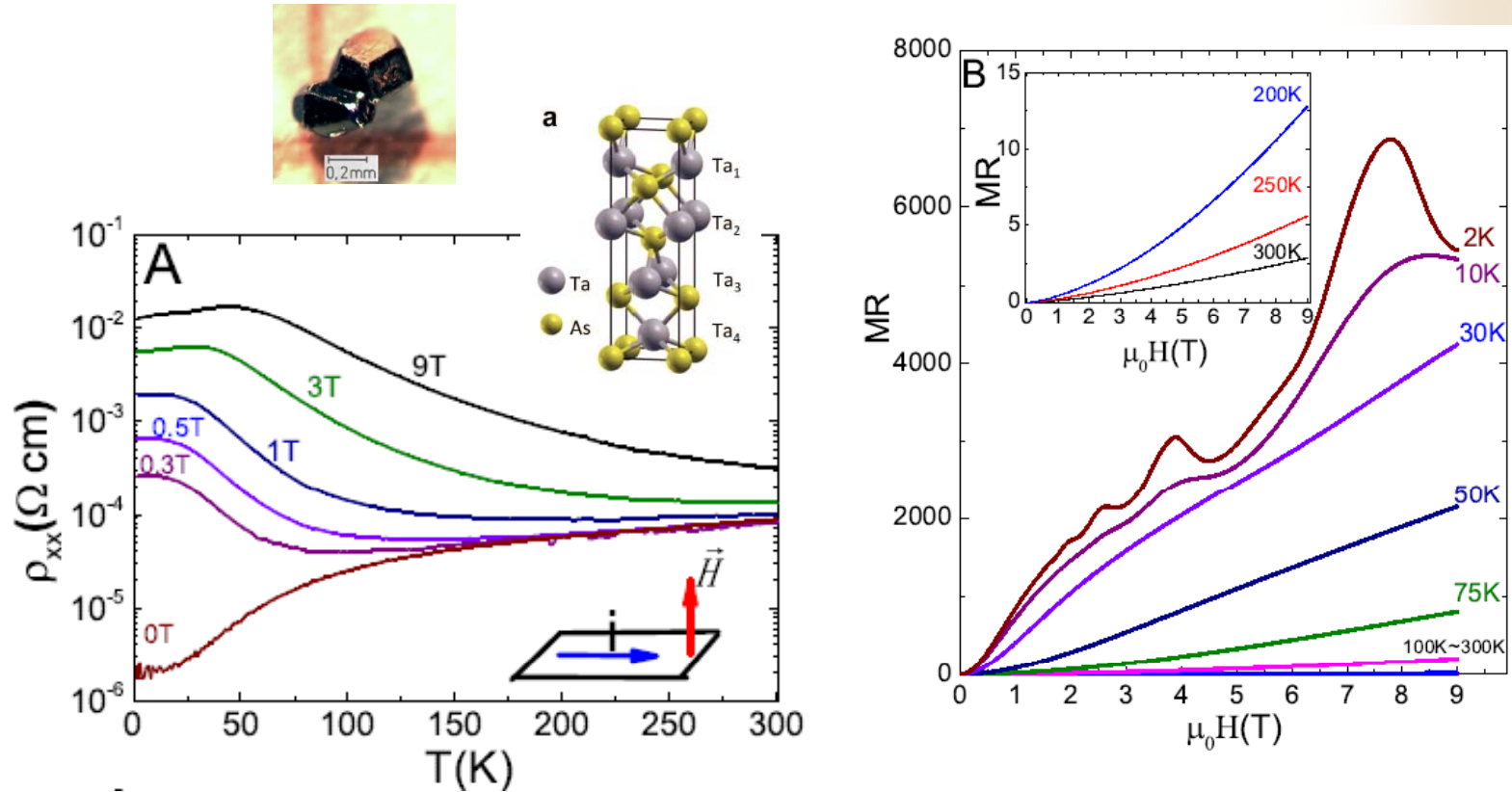
Observation of Clear Weyl Nodes and Fermi Arcs in NbAs



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More details in Su-yang Xu *et al.* arXiv:1504.01350

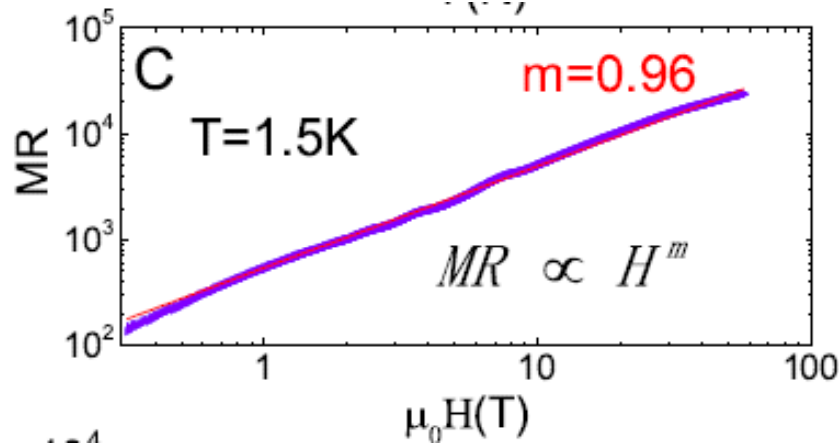
Transport Property: Large MR in TaAs



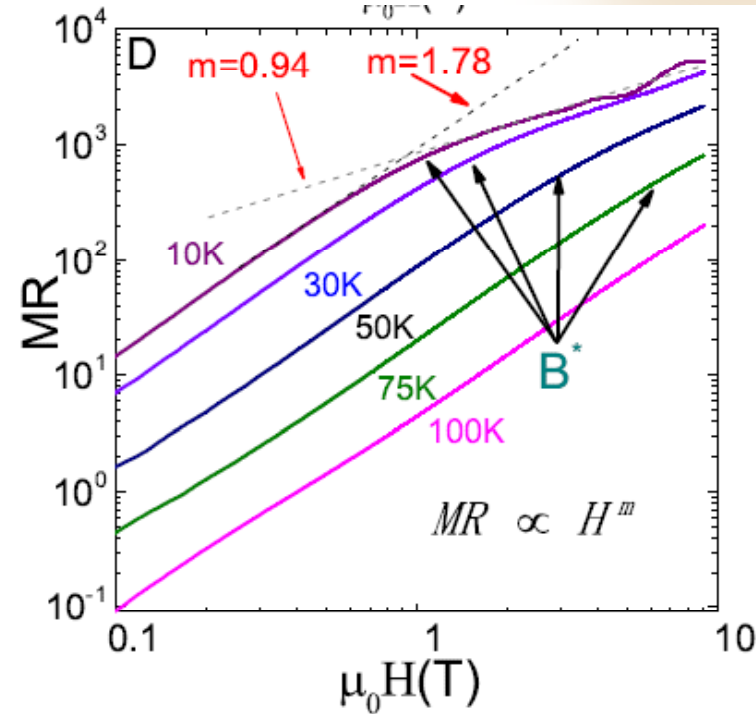
- A small field turns on TaAs to be an insulator
- In 9T, MR~5000 at 2K and ~3 at 300K ($MR = \rho_H/\rho_0 - 1$)
- Extremely strong SdH oscillations at low T



Linear MR with Strong SdH Oscillations at Low T



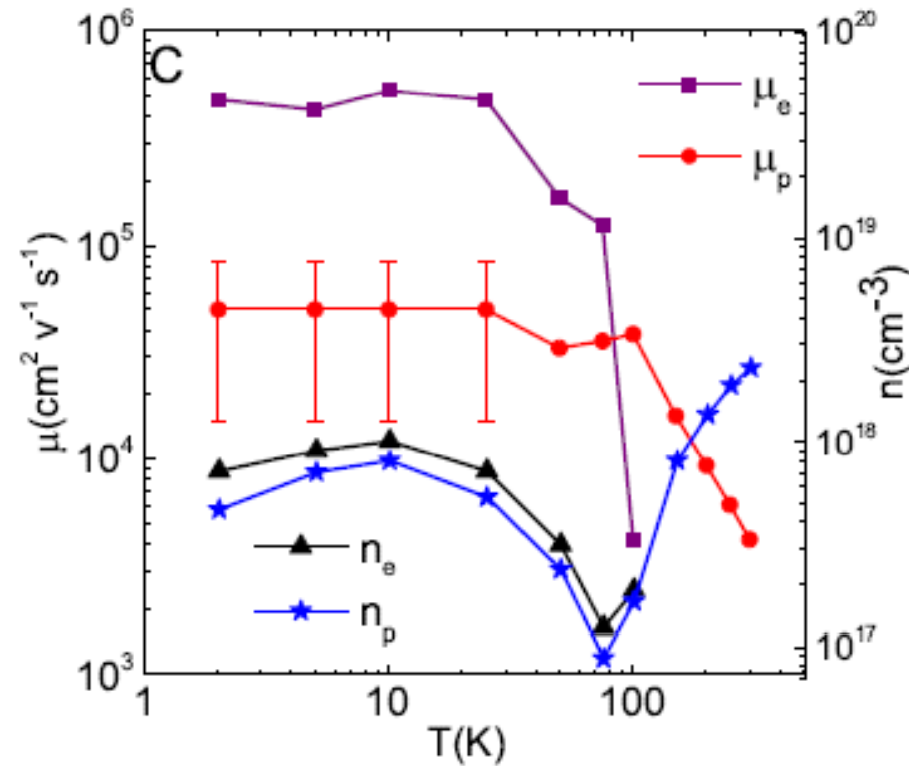
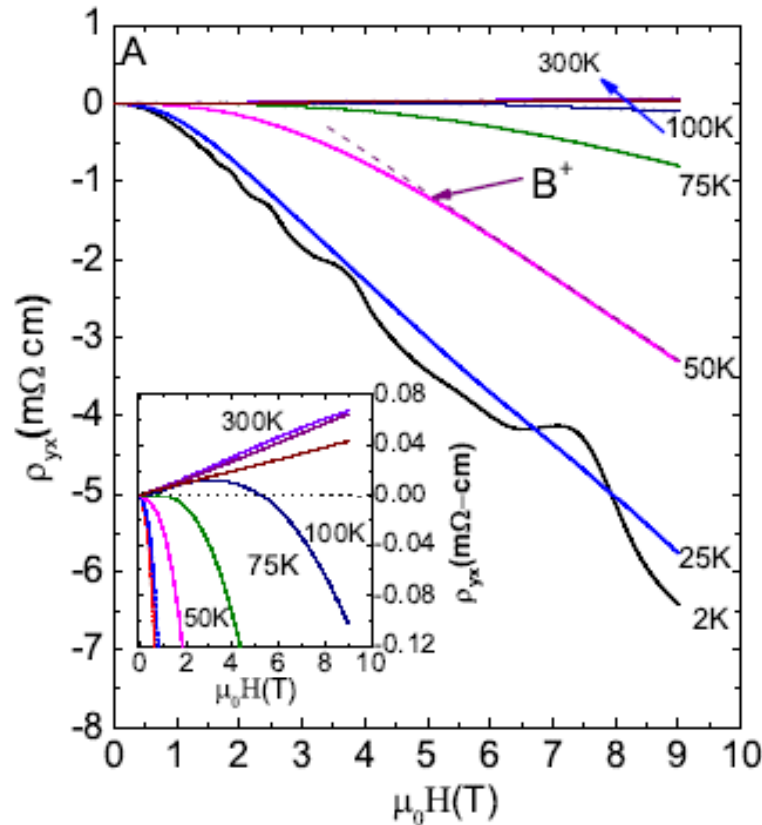
A linear MR beyond the QL up to 55T at low T



MR changes from parabolic to linear at certain B



Two-band Features in Hall Resistance

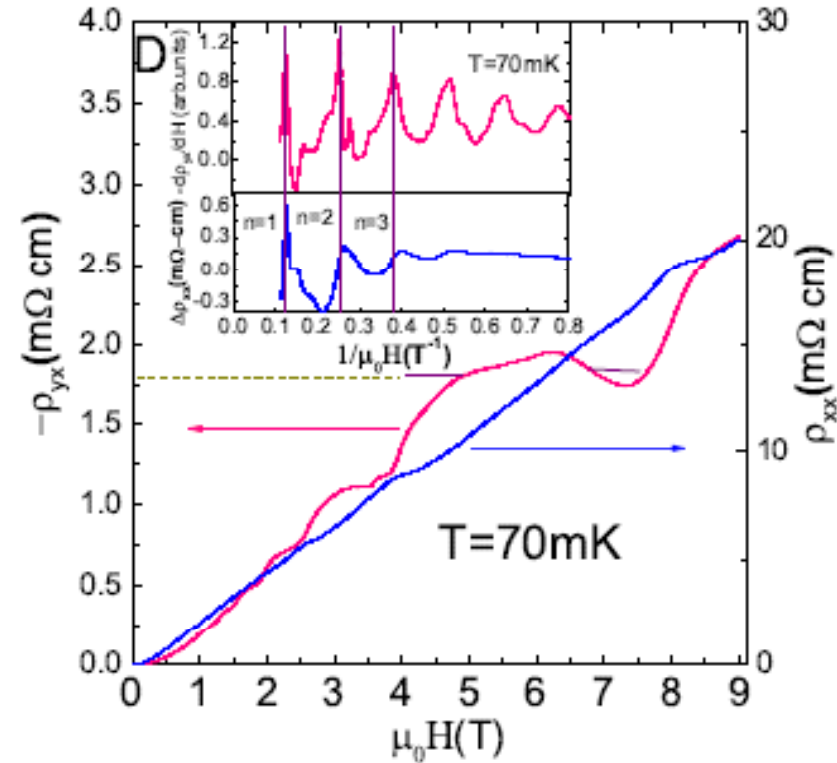
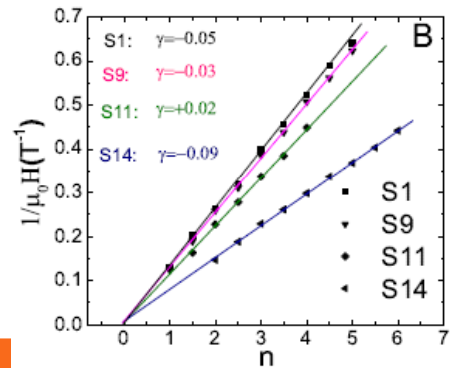
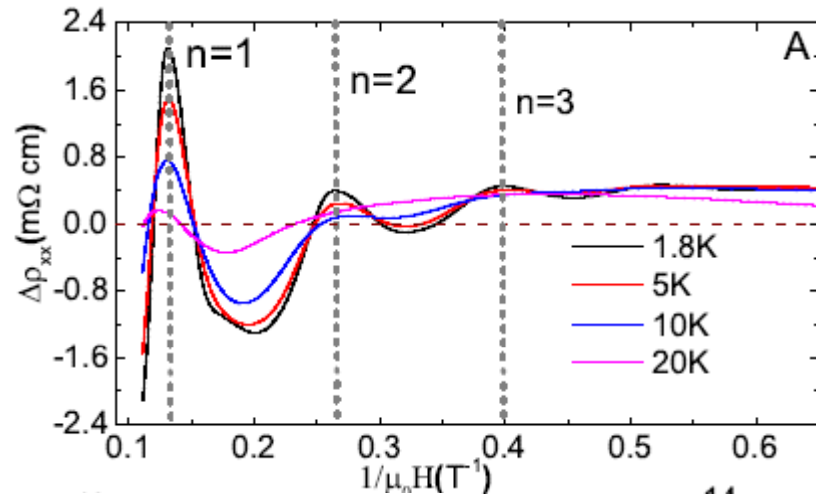


- Compensated semimetal with n and $p \sim 10^{17} \text{cm}^{-3}$
- **e channel with much higher mobility opens below 100K**
- Low T features are dominated by e channel

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Quantum Oscillations in ρ_{xx} and ρ_{xy}

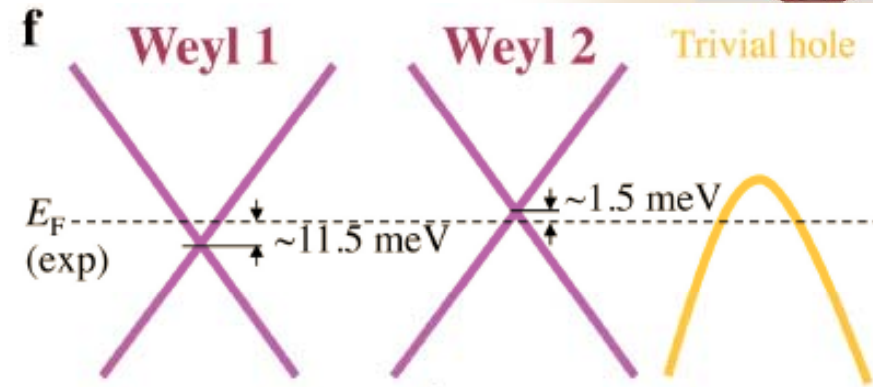
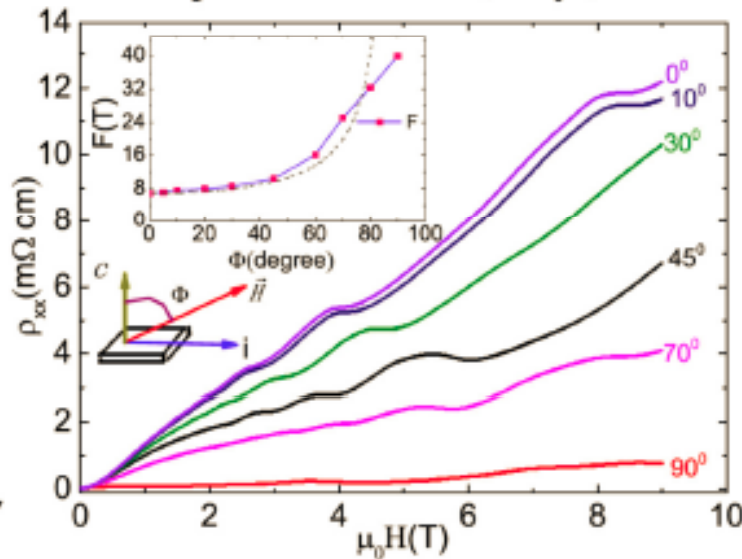


- Strong oscillations come from a tiny e pocket $\sim 8T$
- Non-trivial Berry phase

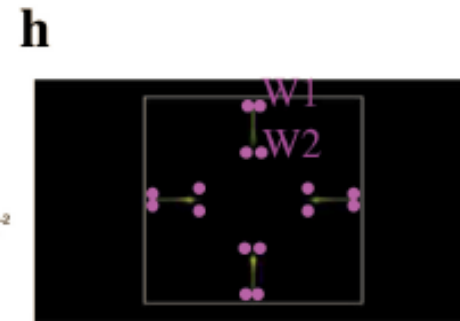
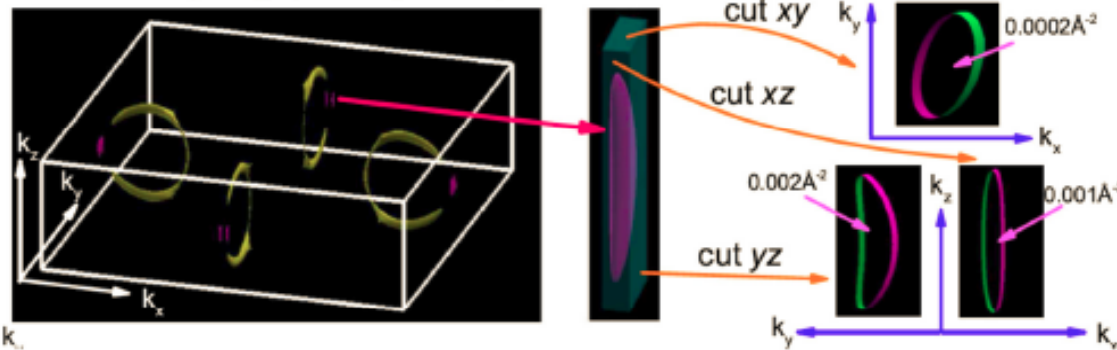
More interesting features at low-T

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Fermi Surface Topology from Transport



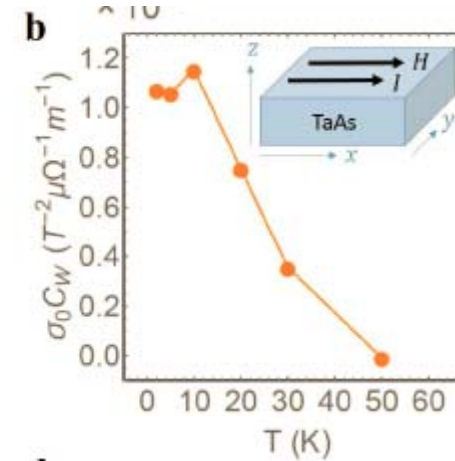
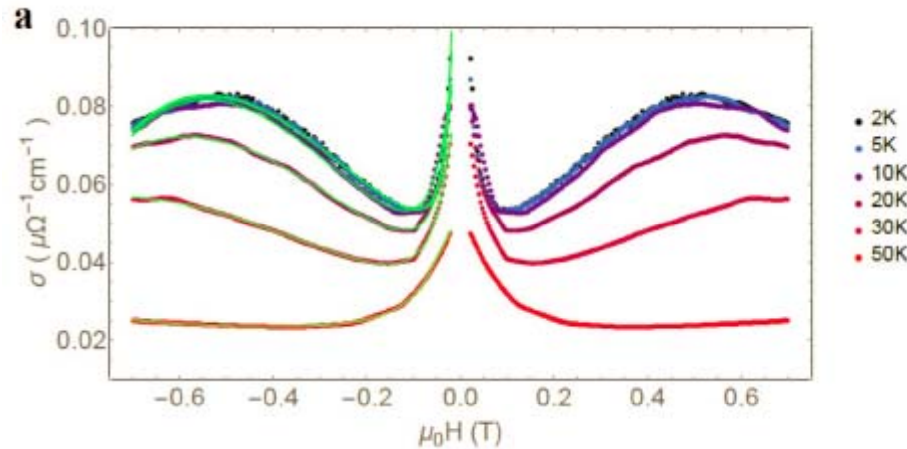
g Calc.: e (Weyl) + h (trivial)



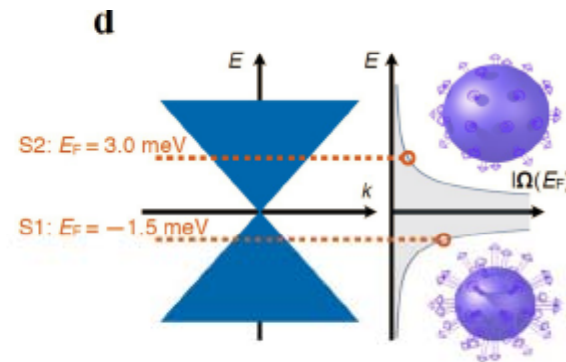
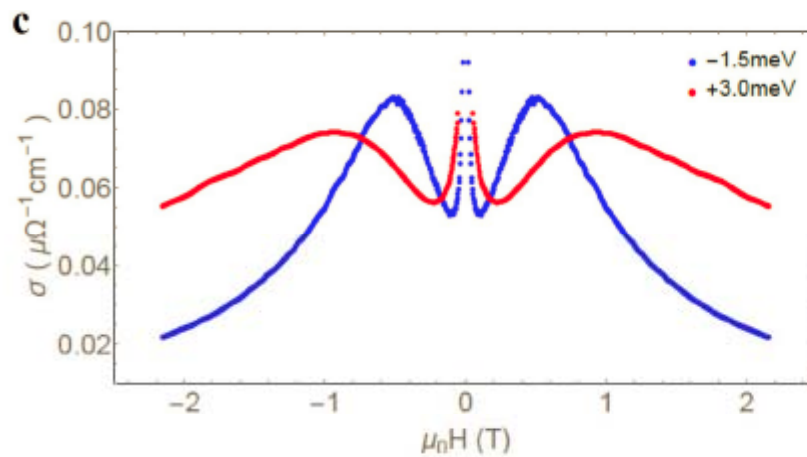
- An ellipsoid-like shape FS consistent with calculations
- Fermi level crosses two Weyl cones and one trivial hole band

Chiral Anomaly

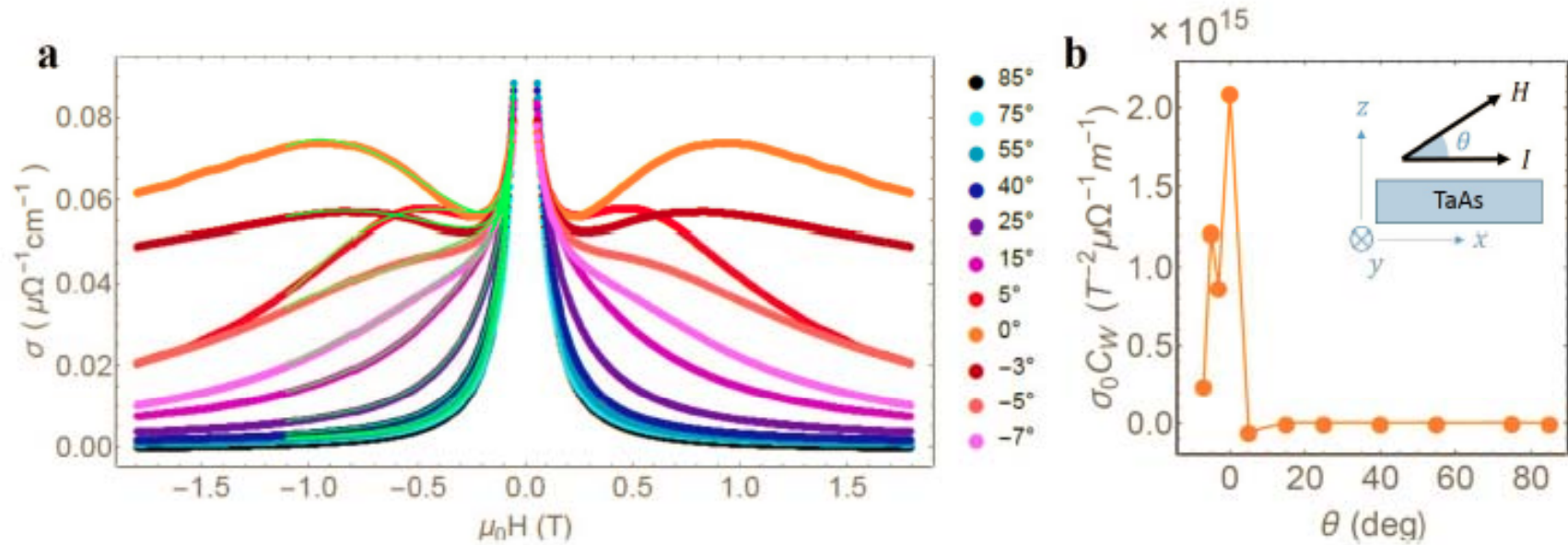
Low field σ is expressed as $\sigma(H) = (\sigma_0 + a\sqrt{H}) (1 + C_W H^2) + \frac{1}{\rho + AH^2} + \frac{1}{\rho' + A'H^2}$,



T dependent chiral anomaly contribution



Chiral Anomaly



➤ This Negative LMR (Positive conductance) is extremely sensitive to the angle deviation



SUMMARY

- We observed first Weyl semimetal TaAs in experiment
- Large MR and strong SdH oscillations
- Negative Longitudinal MR was understood as a Chiral anomaly in TaAs
- Need other candidates of Weyl semimetals and further studies in transport

PLAN

- Best candidates for studying the e-e interactions beyond QL
- Devices: the key point for observing the exotic chiral-related transport